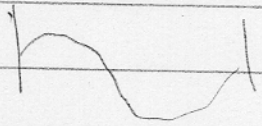


## Last Time

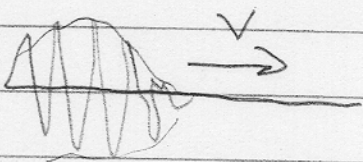
- ① Review - electrons are waves  $\lambda = h/p$  or  $p = \hbar k$
- ② • Standing waves:



only for discrete frequencies  
will an integral number  
of wavelengths fit in the  
confining region

Explains qualitatively why  
electron orbiting proton  
emits discrete light

- ③ • Traveling Waves



showed that a wave packet, e.g.

$$\psi = \cos(k_1 x - \omega_1 t) + \cos(k_2 x - \omega_2 t) + \cos(k_3 x - \omega_3 t)$$

moves with speed  $\omega_1 = p_1^2/2m$  ,  $\omega_2 = p_2^2/2m$

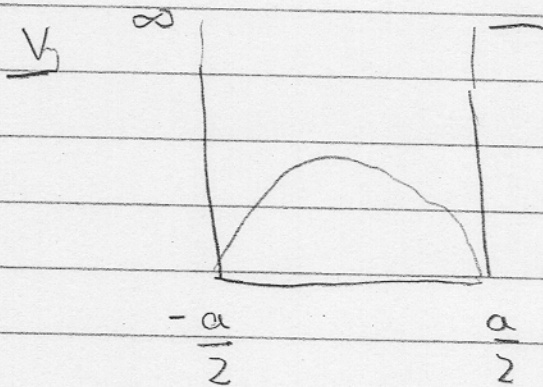
$$v = \frac{\hbar \bar{k}}{m}$$

$$\text{where } \bar{k} = \frac{k_1 + k_2 + k_3}{3}$$

good thing

④ Then we interpreted this wave probabilistically

Ex: Suppose you put an electron inside a box of length  $a$



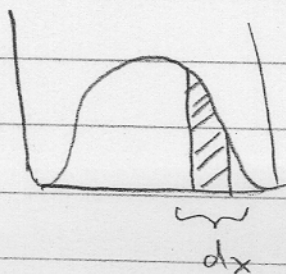
- kind of a model for an atom

- Electron wave fn must be inside box. No probability to be outside the box

We will show <sup>by solving</sup> that the wave fn is

$$\psi(x) = \begin{cases} C \cos(\pi x/a) & |x| < a/2 \\ 0 & |x| > a/2 \end{cases}$$

What does it mean:



$P(x) = |\psi(x)|^2$  is the probability per unit length  
↑  
probability density

Or

$$dP = |\psi(x)|^2 dx = \overset{\checkmark}{P(x)} dx$$

↑

probability of finding the electron in a bin centered at  $x$   
↳ of width  $dx$

Then

$$1 = \sum_{\text{bins}} dP = \int_{-\infty}^{\infty} dx |\psi(x)|^2 = \text{electron must be somewhere}$$

$$\overline{x} = \sum_x x dP = \int_{-\infty}^{\infty} dx x |\psi(x)|^2$$

↑ the average position over the  $x$  weighted by the chance that you will find the electron at  $x$ ,  $|\psi(x)|^2 dx$

$$\overline{x^2} = \int_{-\infty}^{\infty} dx x^2 |\psi(x)|^2$$

$$\overline{3} = \int_{-\infty}^{\infty} dx 3 |\psi|^2 = 3 \int_{-\infty}^{\infty} |\psi|^2 = 3 \quad \text{so } \overline{x} = \overline{x}$$

Then we can use this:

$$1 = \int_{-\infty}^{\infty} dx |\psi(x)|^2 = \int_{-a/2}^{a/2} dx C^2 \cos^2\left(\frac{\pi x}{a}\right)$$

$$= C^2 a \int_{-a/2}^{a/2} \frac{dx}{a} \cos^2\left(\frac{\pi x}{a}\right)$$

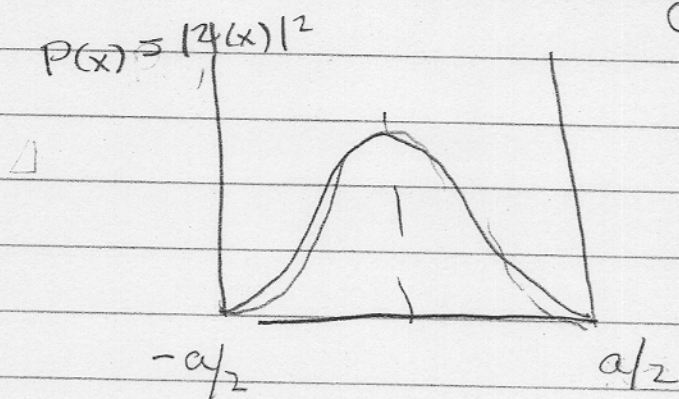
$$1 = C^2 a \int_{-1/2}^{1/2} du \cos^2(\pi u)$$

$$1 = C^2 a \frac{1}{2}$$

$$\sqrt{\frac{2}{a}} = C$$

I explained how to do this integral in recitation

Now what's the average position:

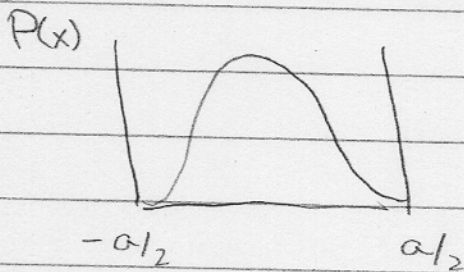


• What's the most likely position?

$x=0$  found by maximizing  $P(x)$

What is the average position?

$$\bar{x} = \int_{-\infty}^{\infty} dx \ x \ |\psi(x)|^2 = 0$$



we will show this  
- it is "intuitively obvious"

$$\bar{x} = \int_{-\infty}^{\infty} dx \ x \ \underbrace{\frac{2}{a}}_{\text{odd}} \underbrace{\cos^2\left(\frac{\pi x}{a}\right)}_{\text{even}} = 0$$

Now what is the spread or standard deviation in the average?

Want to measure

$$x - \bar{x}$$

$$\text{But } \overline{x - \bar{x}} = \bar{x} - \bar{\bar{x}} = \bar{x} - \bar{x} = 0$$

Could take

$$|x - \bar{x}|$$

but integrals are hard usually  
So take

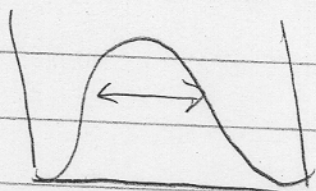
$$\begin{aligned}\sigma^2 \equiv \overline{\Delta x^2} &= \overline{(x - \bar{x})^2} = \overline{(x - \bar{x})(x - \bar{x})} \\ &= \overline{x^2 - 2x\bar{x} + \bar{x}^2} \\ &= \overline{x^2} - 2\bar{x}\bar{x} + \bar{x}^2\end{aligned}$$

$$\sigma^2 = \overline{x^2} - (\bar{x})^2$$

Then the "root mean square" deviation

$$(\Delta x)^2 \equiv \left[ \overline{(x - \bar{x})^2} \right]^{1/2}$$

Example: Compute the standard deviation  
for the particle in the box



First estimate  $\Delta x \sim \frac{a}{2}$

## Solution

$$\Delta X^2 \equiv \overline{X^2} - \overline{X}^2$$

$$(\Delta X)^2 = \int_{-\infty}^{\infty} dx \, x^2 |\psi(x)|^2$$

$$(\Delta X)^2 = \int_{-a/2}^{a/2} dx \, x^2 \frac{2}{a} \cos^2\left(\frac{\pi x}{a}\right)$$

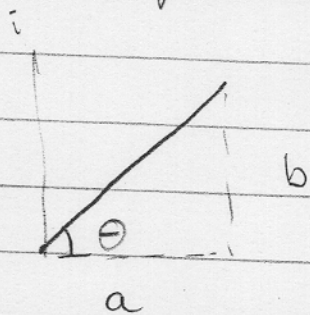
$$(\Delta X)^2 = 2a^2 \int_{-a/2}^{a/2} \frac{dx}{a} \left(\frac{x}{a}\right)^2 \cos^2\left(\frac{\pi x}{a}\right)$$

$$(\Delta X)^2 = 2a^2 \int_{-1/2}^{1/2} du \, u^2 \cos^2(\pi u) \quad \left. \begin{array}{l} \\ \\ \end{array} \right\} u = \frac{x}{a}$$

$$(\Delta X)^2 = 2a^2 \left( \frac{-8 + \pi^2}{2\pi^3} \right) = a^2 (0.06)$$

$$\Delta X = \sqrt{a^2 (0.06)} \approx a (0.25)$$

## Reviewed Complex #'s



$$z = a + bi$$

$$z = r \cos \theta + i r \sin \theta$$

$$r = \sqrt{a^2 + b^2}$$

$$\theta = \tan^{-1} \left( \frac{b}{a} \right)$$

Then

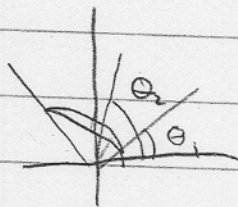
$$e^{i\theta} = \cos \theta + i \sin \theta$$

So

$$z = r e^{i\theta}$$

## Multiply Complex #'s

$$z_1 z_2 = r_1 e^{i\theta_1} r_2 e^{i\theta_2} = r_1 r_2 e^{i(\theta_1 + \theta_2)}$$





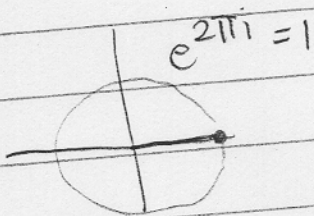
$$\text{Also } \sqrt{z} = \sqrt{r e^{i\theta}} = \sqrt{r} e^{i\theta/2}$$

So Example

$$\sqrt{1} = \pm 1$$



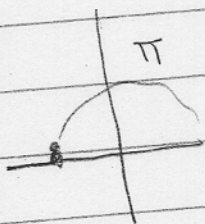
$$1 = e^0 \quad \text{and} \quad 1 = e^{2\pi i}$$



So

$$\sqrt{1} = e^{0/2} \quad \text{and} \quad \sqrt{1} = e^{(2\pi i)/2}$$

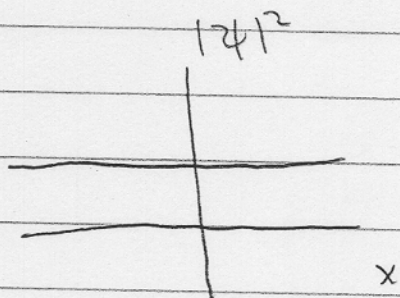
$$\sqrt{1} = e^{i\pi} = -1$$



Then we talked about the purest wave:

$$\psi = c e^{ikx} \quad \leftarrow \text{wavelength is precisely known}$$

$$\psi^* \psi = c^2 e^{-ikx} e^{ikx} = c^2$$



But then the electron is equally likely to be anywhere

## Superposition of two waves (Complex)

$$\psi_1 = e^{ik_1x} + e^{ik_2x} = e^{i15x} + e^{i16x}$$

$$\psi = (\cos(k_1x) + i\sin(k_1x)) + (\cos(k_2x) + i\sin(k_2x))$$

$$\psi = [\cos(k_1x) + \cos(k_2x)] + i[\sin(k_1x) + \sin(k_2x)]$$

$\text{Re}\psi$  = super position of two cos's

$\text{Im}\psi$  = super position of two sines

$$\bar{k} = \frac{k_1 + k_2}{2} = 15.5 \quad \Delta k = k_2 - k_1 = 16 - 15 = 1$$

$$k_1 = \bar{k} - \frac{\Delta k}{2} \quad k_2 = \bar{k} + \frac{\Delta k}{2}$$

So

$$\psi = e^{i(\bar{k} - \Delta k/2)x} + e^{i(\bar{k} + \Delta k/2)x} = e^{i\bar{k}x} (e^{i\frac{\Delta k}{2}x} + e^{-i\frac{\Delta k}{2}x})$$

$$\psi = \underbrace{e^{i\bar{k}x}}_{\text{Carrier wave}} \underbrace{2\cos\frac{\Delta k}{2}x}_{\text{envelope}} = 2(\cos\bar{k}x + i\sin(\bar{k}x)) \cos\frac{\Delta k}{2}x$$

Carrier wave      envelope

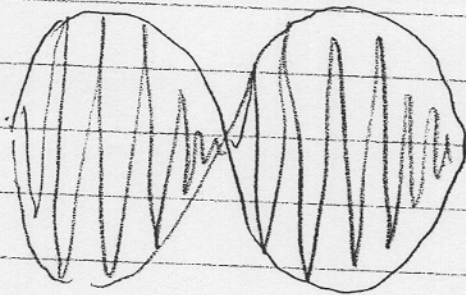
$$\text{Im}\psi = \underbrace{\sin(\bar{k}x)}_{\text{carrier wave}} \underbrace{2\cos\frac{\Delta k}{2}x}_{\text{envelope}}$$

## Summary

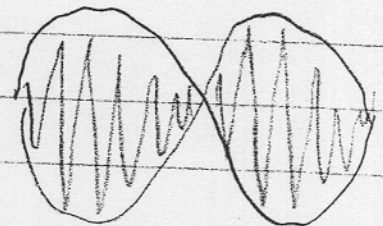
$$\psi_1 = e^{+ik_1x} + e^{iK_2x} = e^{i\bar{k}x} \overbrace{2\cos(\Delta k x)}^{A(x) = \text{envelope funct}} = A(x) e^{i\bar{k}x}$$

## Picture

Re  $\psi$



Im  $\psi$



$$|\psi|^2 = (A^*(x) e^{-i\bar{k}x}) (A(x) e^{+i\bar{k}x}) = A^*(x) A(x) e^0$$

①  $|\psi|^2 = |A(x)|^2$

↑ The probability is determined by the envelope

② But  $\psi = e^{i\bar{k}x} A(x)$

So the momentum / wavelength / wave number is determined by the rapidly fluctuating phase

$$\bar{p} = \hbar \bar{k}$$

Computing the average momentum (see handout)



$$\bar{p} = \sum_{\text{bins}} \hbar \bar{k}_{\text{bin}} P_{\text{bin}}$$

$\uparrow$  momentum in bin       $\swarrow$  prob to be in bin

$$\bar{p} = \sum_{\text{bins}} \hbar \bar{k} \psi^*(x) \psi(x) dx$$

$$= \sum_{\text{bin}} \underbrace{A^* e^{-ik \cdot x}}_{\psi^*} p \underbrace{e^{+ik \cdot x} A}_{\psi} dx$$

Watch:

$$p = i\hbar \frac{d}{dx} e^{+ik \cdot x} = -i\hbar e^{+ik \cdot x} + i\hbar k e^{+ik \cdot x} = \hbar k e^{+ik \cdot x} = p e^{+ik \cdot x}$$

$$\boxed{-i\hbar \frac{d}{dx} e^{+ik \cdot x} = p e^{+ik \cdot x}}$$

$$\bar{S}_p = \sum \psi^* p \psi$$

$$\bar{p} = \sum_{\text{bin}} A^* e^{-ik \cdot x} \left( -i\hbar \frac{d}{dx} \right) e^{+ikx} A$$

$\underbrace{\hspace{10em}}_{\psi^*} \qquad \qquad \qquad \underbrace{\hspace{10em}}_{\psi}$

$$\bar{p} = \int_{-\infty}^{\infty} dx \psi^* \left( -i\hbar \frac{d}{dx} \right) \psi$$

Magic formula

$$-i\hbar \frac{d}{dx} e^{+ikx} = p e^{+ikx}$$

$$-i\hbar \frac{d}{dx} \left( -i\hbar \frac{d}{dx} e^{+ikx} \right) = -i\hbar \frac{d}{dx} p e^{+ikx} = p^2 e^{+ikx}$$

So

$$\bar{p}^2 = \sum_b \psi^* p^2 \psi$$

$$= \sum_b \psi^* \left( -i\hbar \frac{d}{dx} \right)^2 \psi$$

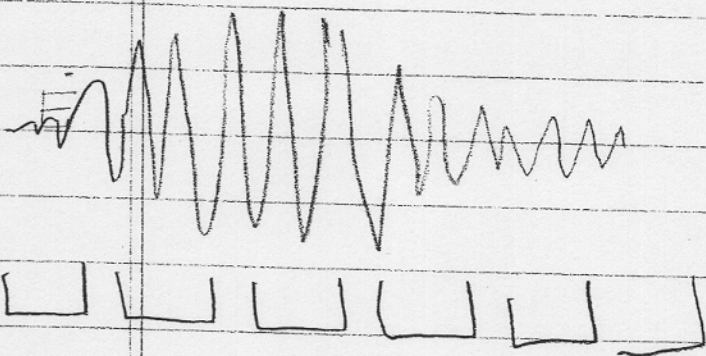
$$\bar{p}^2 = \sum_b \psi^* \left( -\hbar^2 \frac{d^2}{dx^2} \right) \psi$$

$$KE = \frac{p^2}{2m} \quad \text{so, ...}$$

$$\overline{KE} = \int_{-\infty}^{\infty} dx \psi^* \left( -\frac{\hbar^2}{2m} \frac{d^2}{dx^2} \right) \psi$$

$$KE \sim \frac{\hbar^2}{2mL^2}$$

So Far we've been discussing space what about time? and energy



$$\overline{E} = \sum_{\text{bins}} E_{\text{bin}} \cdot P_{\text{bin}}$$

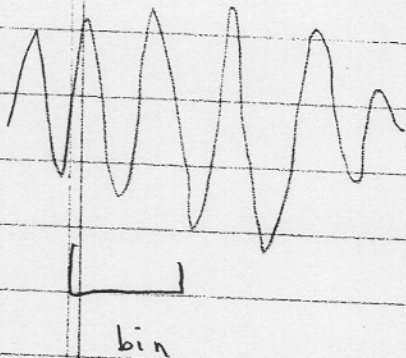
$$\overline{E} = \sum_{\text{bins}} E_{\text{bin}} |A_{\text{bin}}|^2 dx = \sum_{\text{bins}} \underbrace{A^* e^{+i\omega t + ikx}}_{\psi^*} E_b \underbrace{A e^{-i\omega t - ikx}}_{\psi}$$

Watch

$$\hbar i \hbar \frac{d}{dt} e^{-i\omega t} = \hbar i \hbar (-i\omega) e^{-i\omega t} = \hbar \omega e^{-i\omega t} = E e^{-i\omega t}$$

$$\overline{E} = \int_{-\infty}^{\infty} dx \psi^* \hbar i \hbar \frac{\partial}{\partial t} \psi$$

# The Schrödinger Equation



In a single bin we have two expressions for the energy

$$\psi^* \left[ \frac{p^2}{2m} + V(x) \right] \psi(x) = \psi^* [E] \psi$$

$$\psi^* \left[ \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \psi(x) = \psi^* \left[ i\hbar \frac{\partial}{\partial t} \right] \psi$$

So Find:

$$\left[ \frac{-\hbar^2}{2m} \frac{d^2}{dx^2} + V(x) \right] \psi(x) = i\hbar \frac{\partial \psi}{\partial t}$$