

Last Time

- ① Wrote down the Schrödinger Equation in 3D


$$\left[\frac{-\hbar^2}{2m} \nabla^2 + V(r) \right] \psi = E \psi$$

$$\nabla^2 = \frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2} = \text{"Laplacian"}$$

$$= \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{1}{r^2} \left[\frac{1}{\sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} + \frac{1}{\sin^2 \theta} \frac{\partial^2}{\partial \phi^2} \right]$$

Try Separation of variables

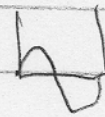
$$\psi_{nlm} = R_{nl}(r) \Theta_{lm}(\theta) \Phi_m(\phi)$$

① $n-1$ = total number of excitations angular or radial  0 excitations

$$= 0, 1, 2, 3, \dots$$

② l = total number of angular excitations

$$= 0, 1, \dots, n-1$$



1 excitation

angular momentum

$$\leftarrow L^2 = l(l+1)\hbar^2 = L_x^2 + L_y^2 + L_z^2$$

③ $(n-1) - l = \text{total number}$
radial excitations

④ $m = 0, \pm 1, \pm 2, \dots, l$

$L_z = \text{angular momentum around } z\text{-axis} = m\hbar$

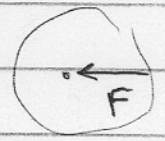
$|m| = \text{total number of azimuthal}$
excitations

$+$ = counter clockwise

$-$ = clockwise

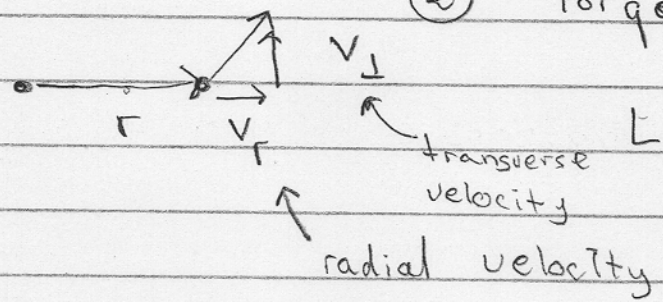
Lets Study Classical Orbits:

① Energy is Constant!



$$KE + PE = E$$

② Torque $\vec{F} \times \vec{r} = 0$



$$L = m v_{\perp} r = \text{constant}$$

So

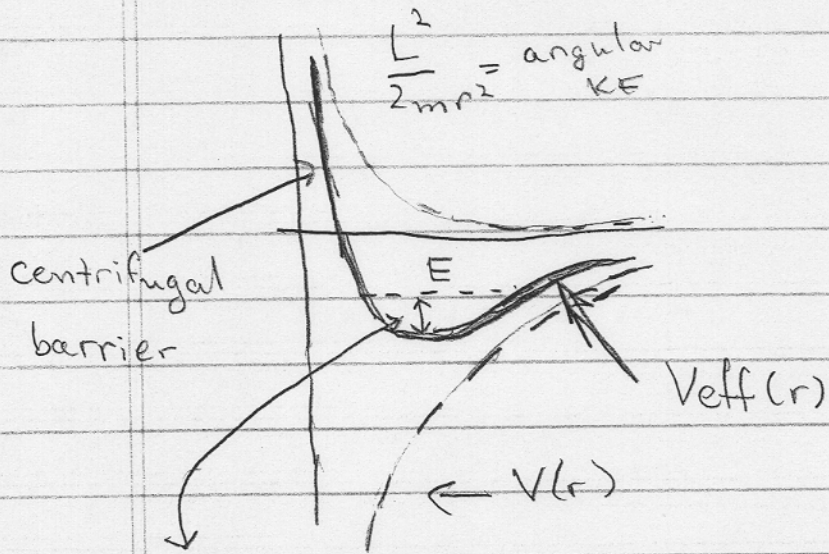
$$\begin{aligned} KE &= \frac{1}{2} m v_r^2 + \frac{1}{2} m v_{\perp}^2 \\ &= \frac{1}{2} m v_r^2 + \frac{1}{2} \frac{(m v_{\perp} r)^2}{m r^2} \\ &= \frac{1}{2} m v_r^2 + \frac{L^2}{2 m r^2} \end{aligned}$$

Thus we see that

$$\frac{1}{2} m v_r^2 + \underbrace{\frac{L^2}{2 m r^2} + V(r)} = E$$

$$\frac{1}{2} m v_r^2 + \underbrace{\frac{L^2}{2 m r^2} + V(r)}_{\equiv V_{\text{eff}}(r)} = E \quad \text{effective potential}$$

$$V(r) = -\frac{e^2}{4\pi\epsilon_0 r}$$



$$\text{radial KE} = \frac{1}{2} m v_r^2 = E - V_{\text{eff}}(r)$$

Note the minimum happens

$$\frac{\partial V_{\text{eff}}}{\partial r} = \frac{L^2}{mr^3} + \frac{\partial V}{\partial r} = 0$$

$$\frac{L^2}{mr^3} = \frac{\partial V}{\partial r}$$



$$\frac{m v_{\perp} r^2}{r^2} = |\vec{L}|$$

$$m \frac{v^2}{r} = |\vec{F}|$$

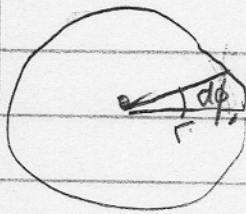
Condition for circular orbit

Now return to quantum mechanics

$$(KE + PE)\psi = E\psi$$

KE = radial KE + angular KE \leftarrow how to define

①



$$L_z = r p_{\phi}$$

$$= r \left(-i\hbar \frac{d}{dx} \right)$$

$$= \cancel{r} \left(-i\hbar \frac{d}{\cancel{r} d\phi} \right)$$

$$L_z = -i\hbar \frac{\partial}{\partial \phi}$$

\leftarrow angular momentum is an angular derivatives

②

$$\vec{L} = \begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ x & y & z \\ p_x & p_y & p_z \end{vmatrix} = \vec{r} \times \vec{p}$$

$$p_x = -i\hbar \frac{\partial}{\partial x}$$

After a considerable amount of algebra
(See Appendix M)

Find

$$L^2 = L_x^2 + L_y^2 + L_z^2$$

$$L^2 = -\hbar^2 \left[\frac{1}{\sin\theta} \frac{\partial}{\partial\theta} \sin\theta \frac{\partial}{\partial\theta} + \frac{1}{\sin^2\theta} \frac{\partial^2}{\partial\phi^2} \right]$$

Cool: '

$$KE = \frac{-\hbar^2}{2m} \nabla^2$$

$$= \underbrace{\frac{-\hbar^2}{2m} \left(\frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} \right)}_{\text{radial KE}} + \underbrace{\frac{L^2}{2mr^2}}_{\text{angular KE}}$$

radial KE

angular KE

Then the Schrödinger Equation becomes

$$\frac{1}{\psi} (KE + PE) \psi = E \psi$$

$$(**) \frac{1}{\psi} \left[\frac{-\hbar^2}{2m} \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{L^2}{2mr^2} + V(r) \right] \psi = E$$

Now try separation of Variables

$$* \quad \psi = R \underbrace{\Theta \Phi}_{\equiv Y(\theta, \varphi)}$$

known as a Spherical harmonic

Substituting (*) into the Schrödinger Equation (**) and going through the usual separation of variables (Hmunk)

$$\rightarrow \left[\left(-\frac{\hbar^2}{2m} \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} \right) + \underbrace{E_{\Omega}}_{\text{angular KE}} + V(r) \right] R = E R$$

radial Schrödinger radial wave fn

$$\underbrace{\left(\frac{\hbar^2}{2mr^2} \right)}_{\text{angular KE}} \underbrace{Y(\theta, \varphi)}_{\text{angular wave-fcn}} = E_{\Omega} Y(\theta, \varphi)$$

lives on sphere!

Very much like particle in box
 E_{Ω} (the angular KE) must be adjusted
so that the wave fn $Y(\theta, \varphi)$ fits
an integral number of wavelengths on the sphere

Find E_{Ω} (angular KE) is discrete

$$E_{\Omega} = \frac{l(l+1)\hbar^2}{2mr^2} \quad \text{compare} \quad \text{PIB} = \frac{\hbar^2 \pi^2}{2mL^2} n^2$$

$$= \frac{\overline{L^2}}{2mr} \quad \overline{L^2} = l(l+1)\hbar^2$$

where $l = 0, 1, \dots, n-1$

So we have

$$\frac{-\hbar^2}{2m} \left[\frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \underbrace{\frac{l(l+1)\hbar^2}{2mr^2} + V(r)}_{V_{\text{eff}}^l(r)} \right] R_{nl} = E_{nl} R_{nl}$$

Finally motivated by $\mathcal{P}(r) = 4\pi r^2 |R|^2$ define

$u_{nl} \equiv \sqrt{4\pi r} R_{nl}$ so $\mathcal{P}(r) = |u_{nl}|^2$, find

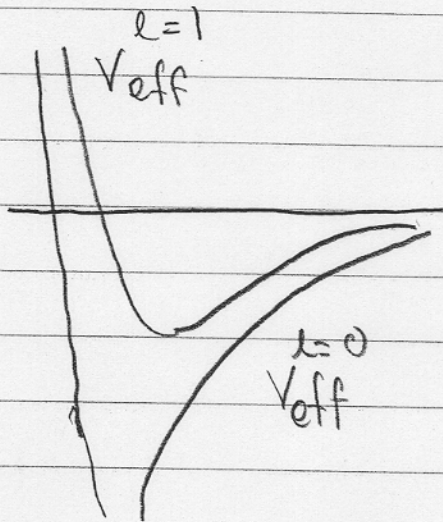
that (Hmwrk)

$$\left[-\frac{\hbar^2}{2m} \frac{\partial^2}{\partial r^2} + V_{\text{eff}}^l(r) \right] u_{nl}(r) = E_{nl} u_{nl}$$

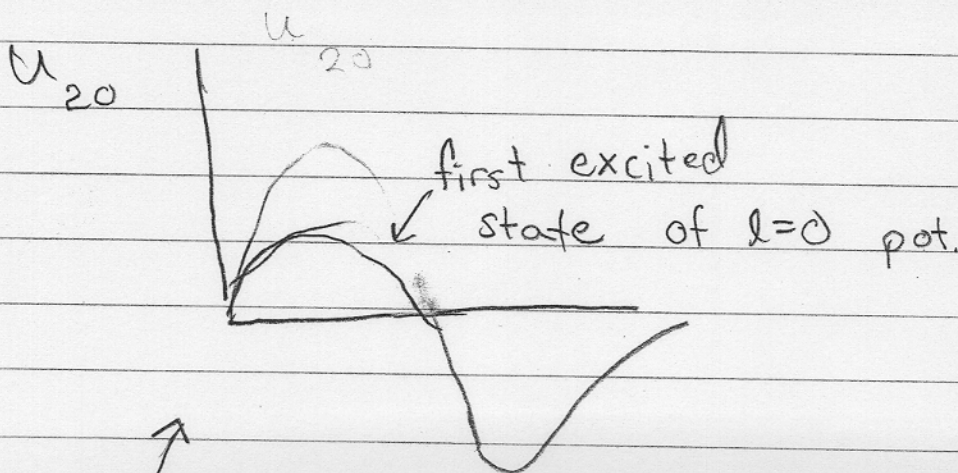
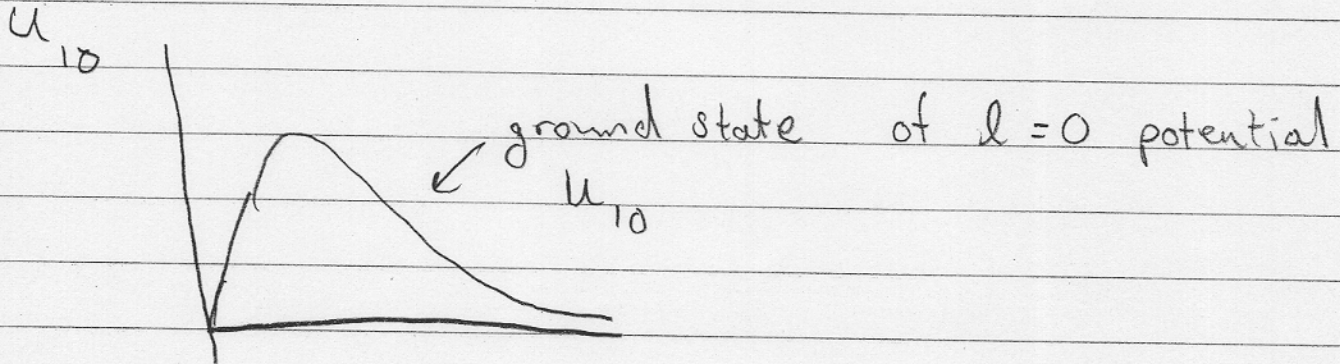
↑ Effective one 1d Schrödinger Equation

Picture

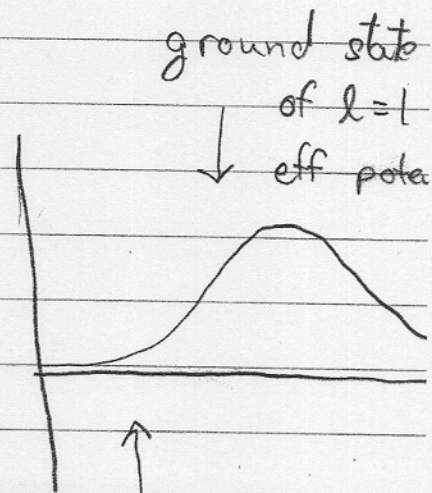
$$V_{\text{eff}} = \frac{l(l+1)\hbar^2}{2mr^2} + V(r)$$



Now (see slides)

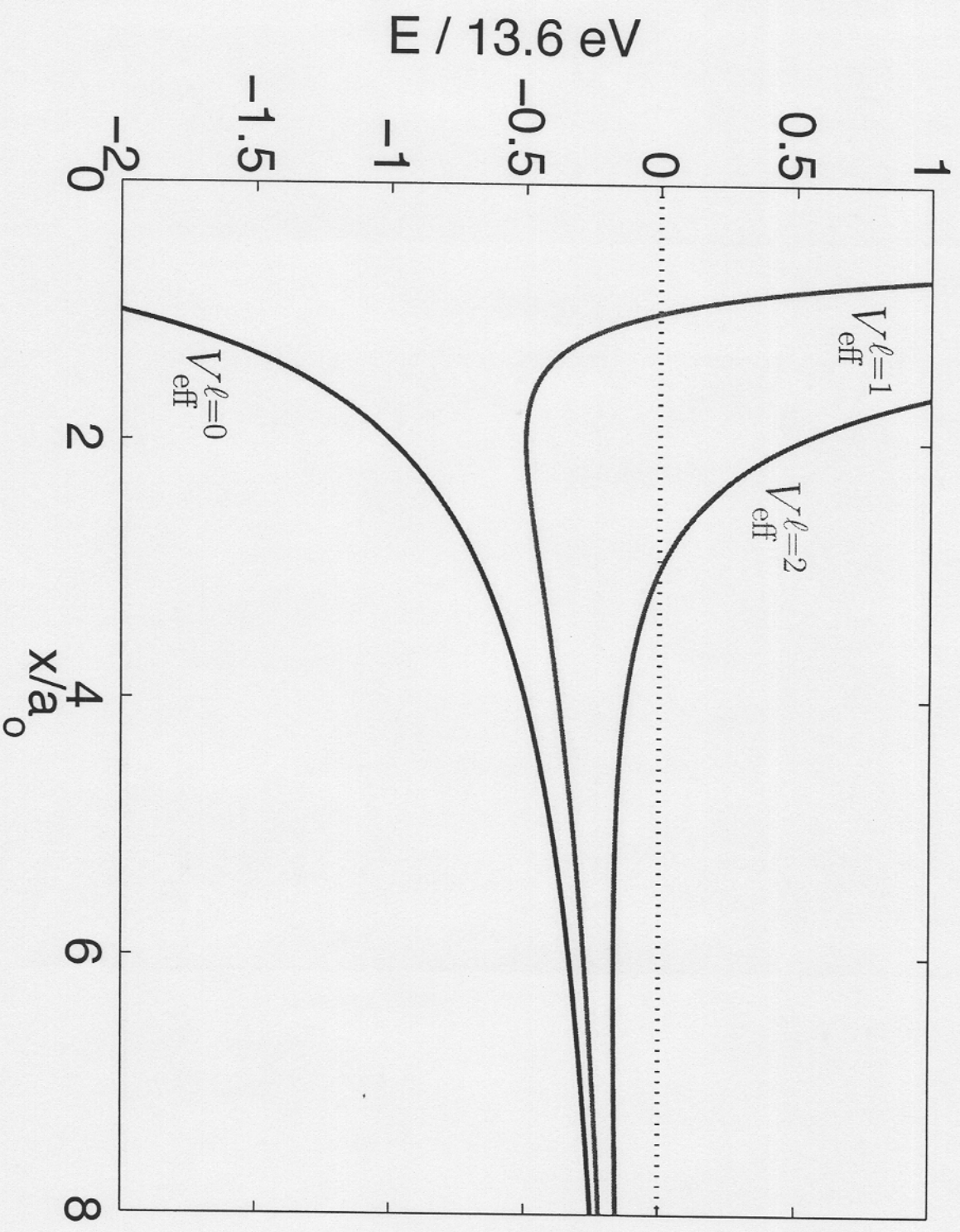


one radial excitation, no angular

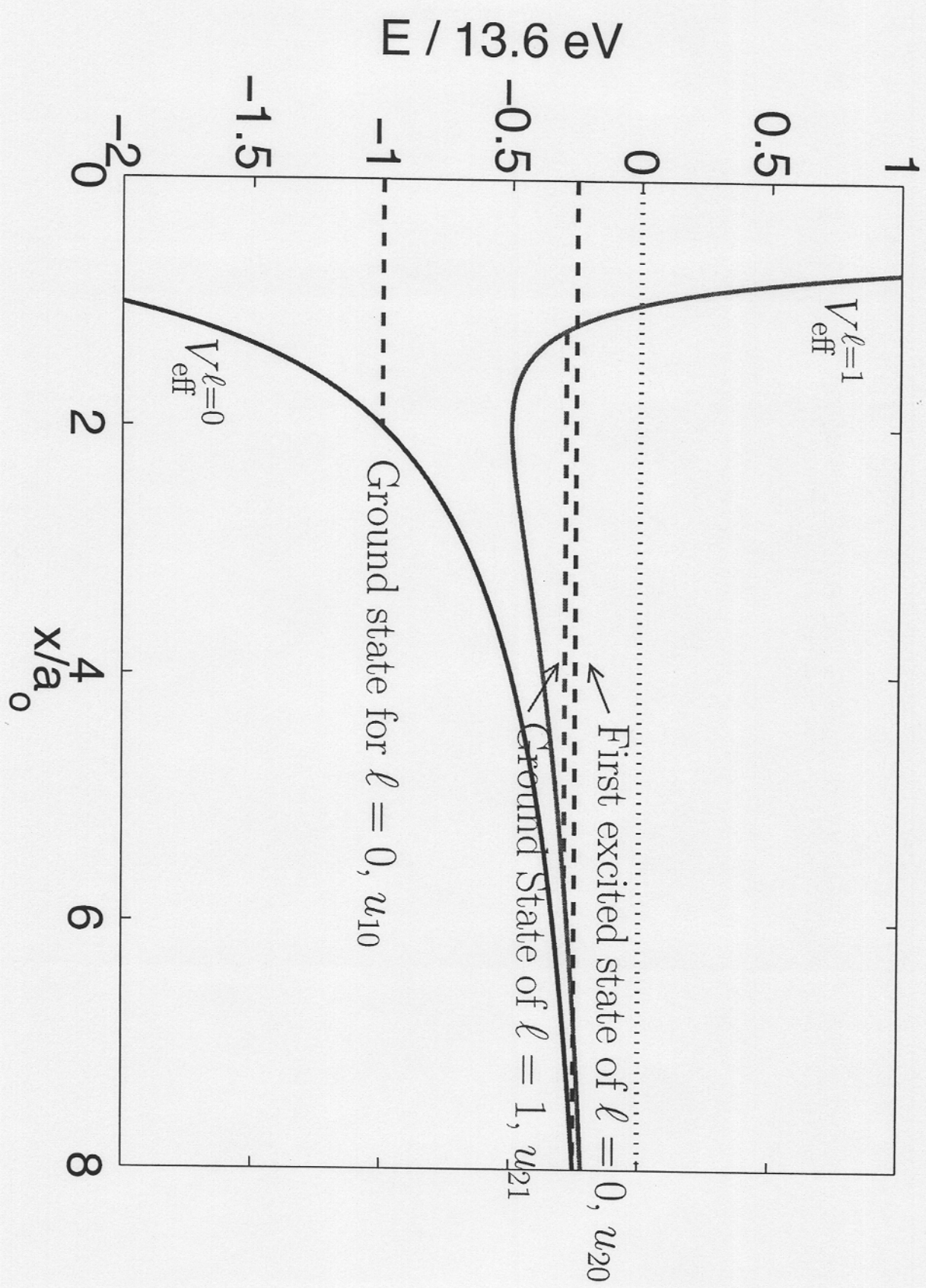


one angular ex but no radial

The effective potential for various l



The ($n = 1, \ell = 0$) and the ($n = 2, \ell = 0$) and ($n = 2, \ell = 1$) states



The ($n = 1, \ell = 0$), the ($n = 2, \ell = 1$) and the ($n = 3, \ell = 2$) states

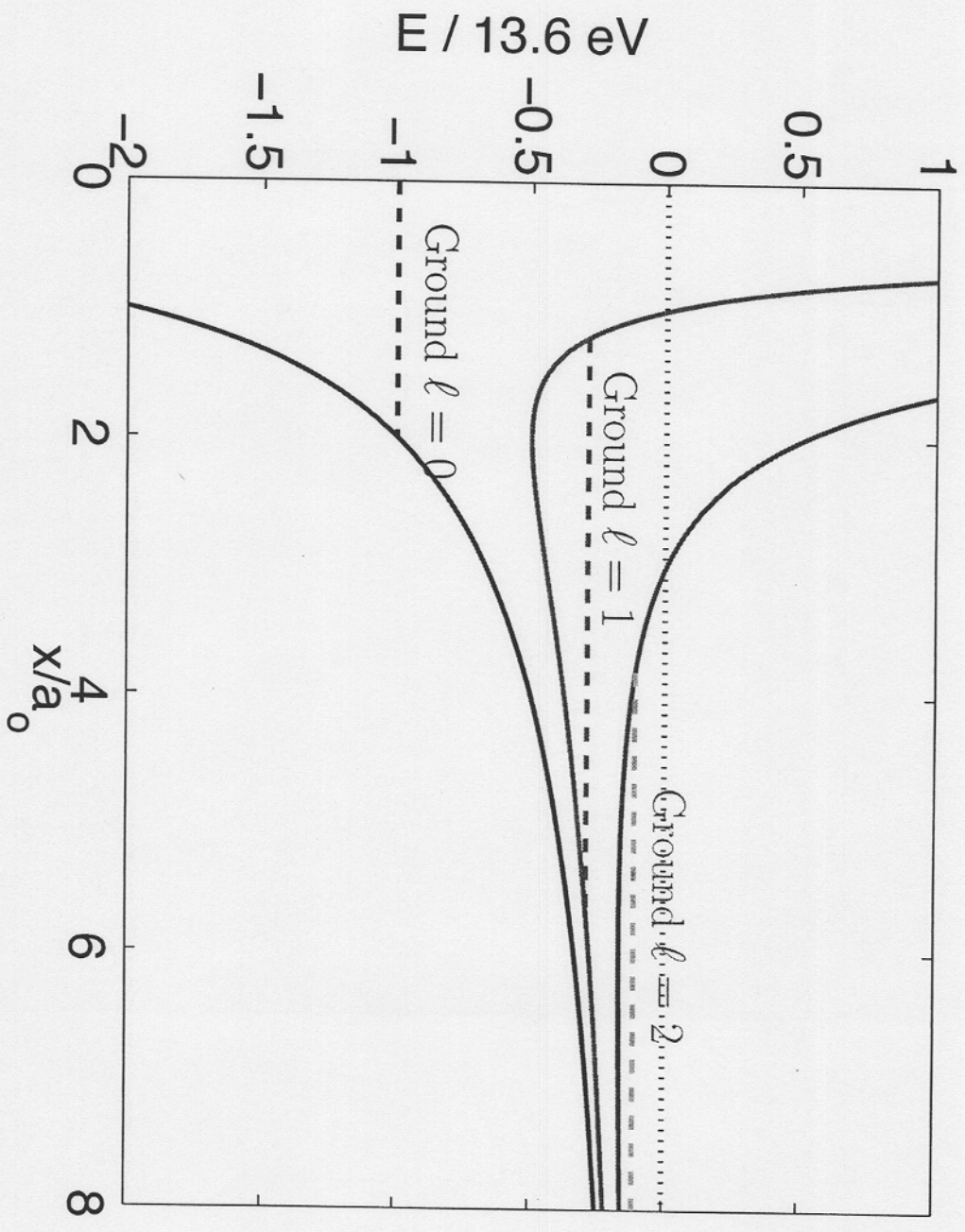
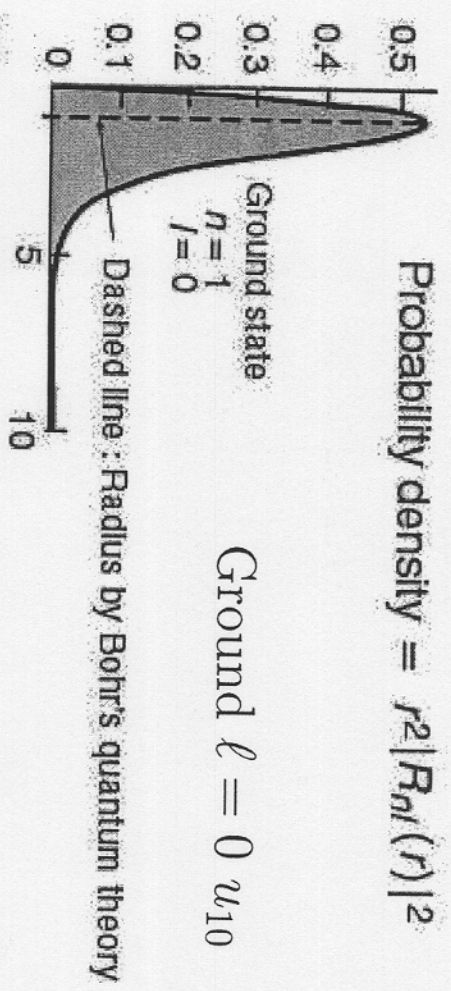
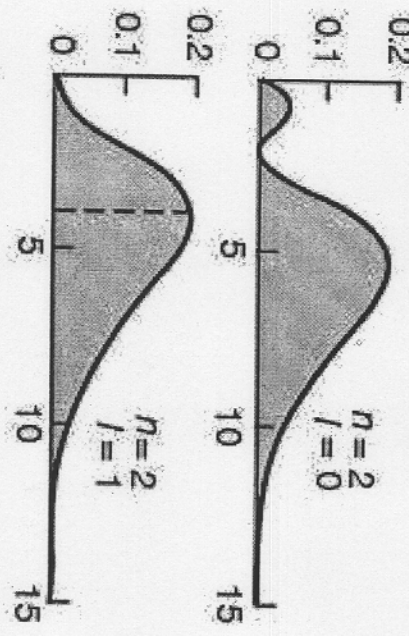


Fig. (C)

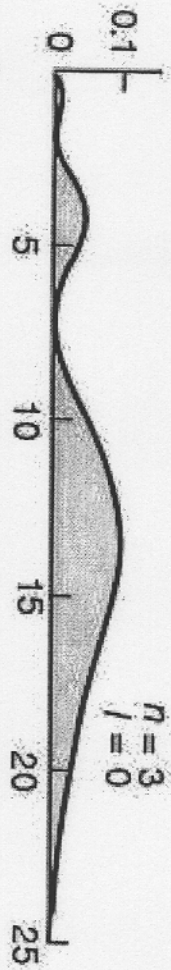
Probability density = $r^2 |R_{nl}(r)|^2$



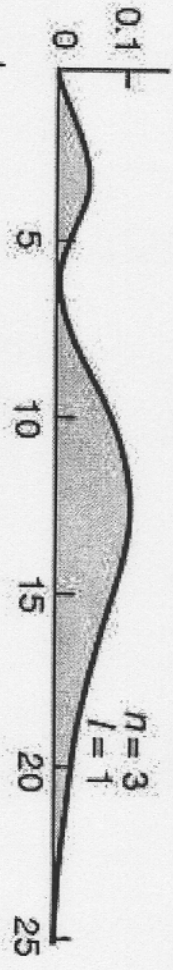
First exc for $l = 0$ u_{20}



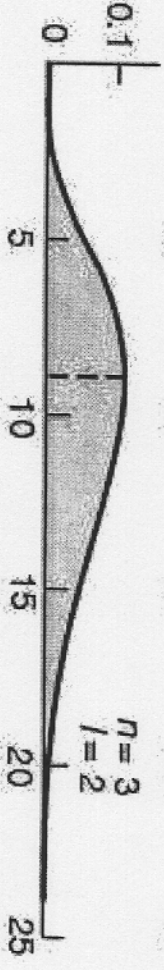
Ground for $l = 1$ u_{21}



Second $l = 0$ u_{30}



First exc. for $l = 1$ u_{31}



Ground $l = 2$ u_{32}

r/a_0 (a_0 : Bohr radius)