

Last Time

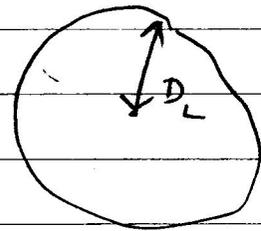
- Discussed how observations of supernova at various red shifts reveal the acceleration history of the universe

$$L = \frac{\text{Energy}}{\text{Time}} = \text{Abs. Luminosity}$$

$$l = \frac{\text{Energy}}{\text{Area (time)}} = \text{Apparant Luminos}$$

$$l = \frac{L}{4\pi D_L^2}$$

$$\Rightarrow D_L = \left(\frac{L}{4\pi l} \right)^{1/2}$$



= Luminosity distance

\Rightarrow Red Shift

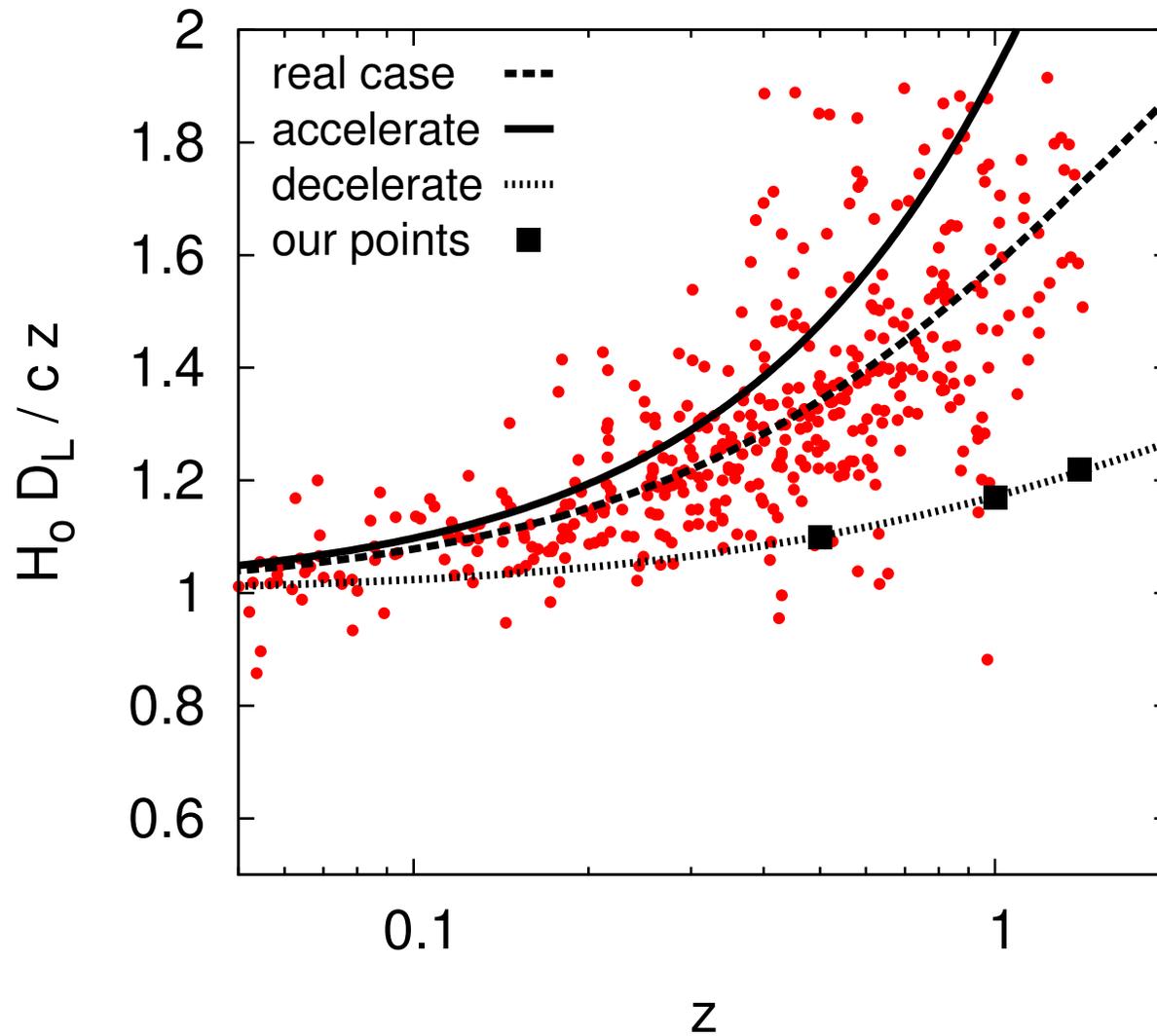
$$\lambda_{\text{obs}} = (1 + z) \lambda_{\text{emit}}$$

Then D_L vs. z , allows one to

determine $a(t)$ vs. t . D_L vs z . is like dist vs time.

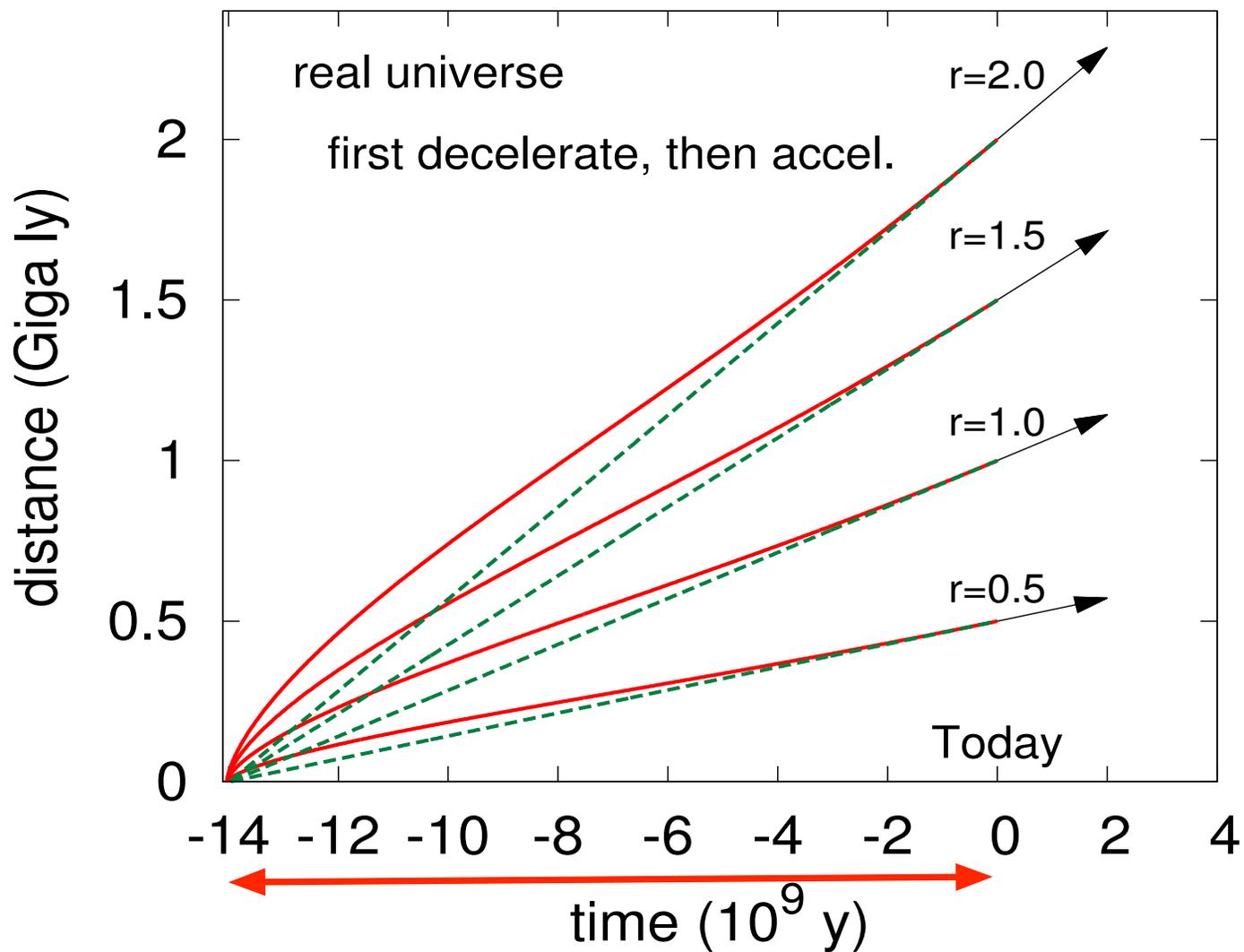
\uparrow
scale factor = % smaller the universe was at time = t relative to today

What cosmology does the data prefer?



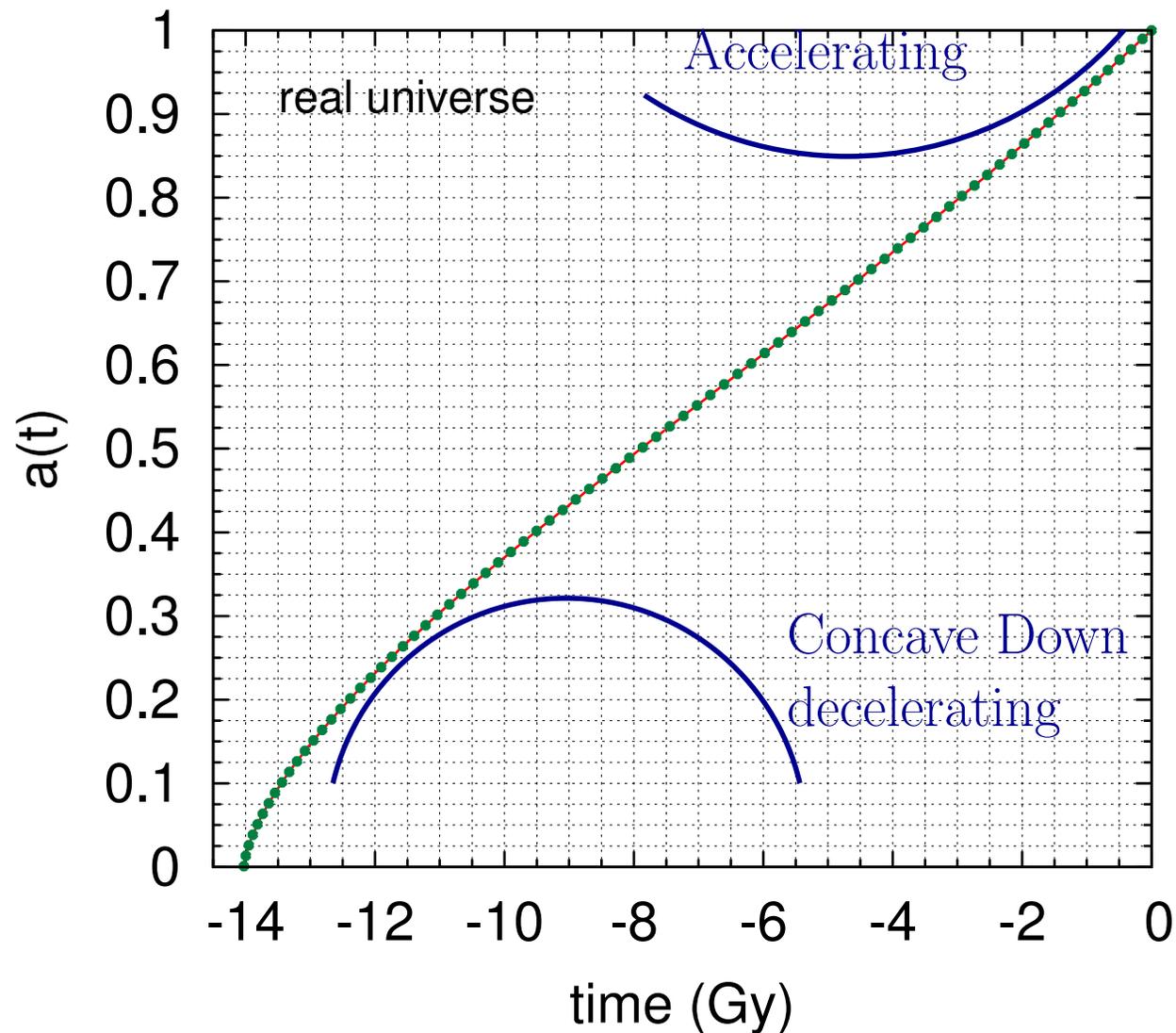
The data want first a period of deceleration then a period of acceleration!

Preferred acceleration history:



Lifetime of Universe is (for this cosmology)
14 Billion Years

The expansion of the universe is summarized by: $a(t)$ =fraction of universe-size today

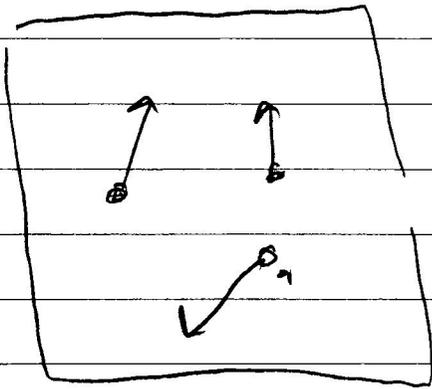


The data want first a period of deceleration, then a period of acceleration

8 Gy after big bang or 6 Gy ago!

Review of Temperature

Then we started talking about Black Body radiation and temperature



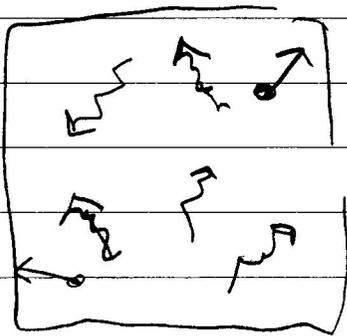
$$\frac{E}{N} \sim k_B T$$

↑
energy per particle

Where k_B is a constant:

$$k_B = \frac{1/40 \text{ eV}}{300^\circ \text{K}} = \frac{\text{energy}}{^\circ \text{K}}$$

Then why do things glow when they get hot:



Atoms bounce around giving some of their energy to ^{creatings} the light particles -- photons

- ① The higher the temperature the higher the frequency, the shorter the wavelength.
- ② Also see more photons the higher the temperature (since the wavelength is shorter more fit in box)

typical Temperature of Black Body

The energy of a photon is:

$$E_{\text{typ}} = h f_{\text{typ}} \sim k_B T$$

In general don't just find photons at one frequency but find photons at a range of frequencies

$$h f_{\text{mp}} = E_{\text{mp}} = 2.82 k_B T$$

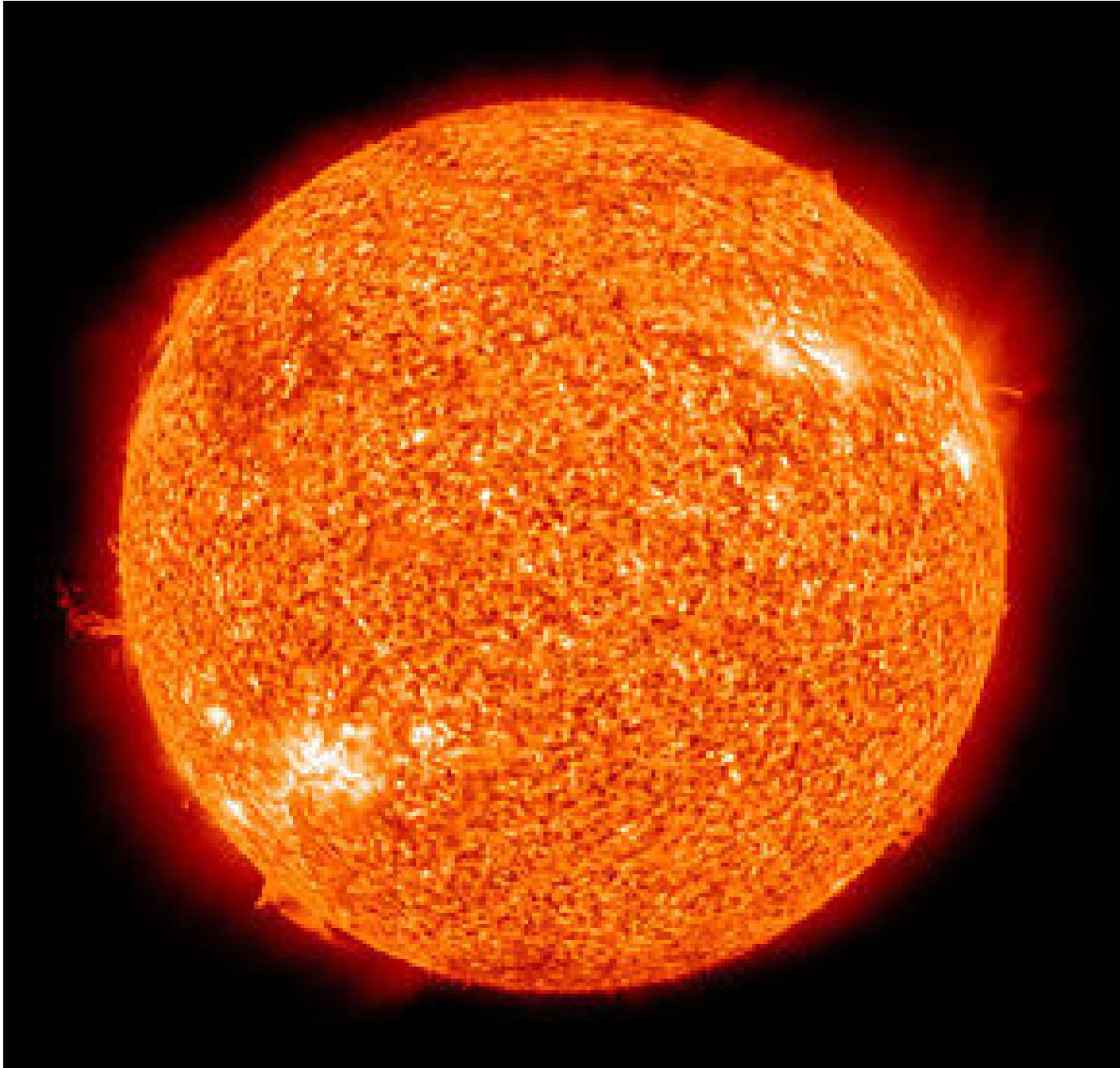
Energy and frequency where it is most likely to find a photon

* Do Problem

The hot walls in equilibrium with the radiation in an oven



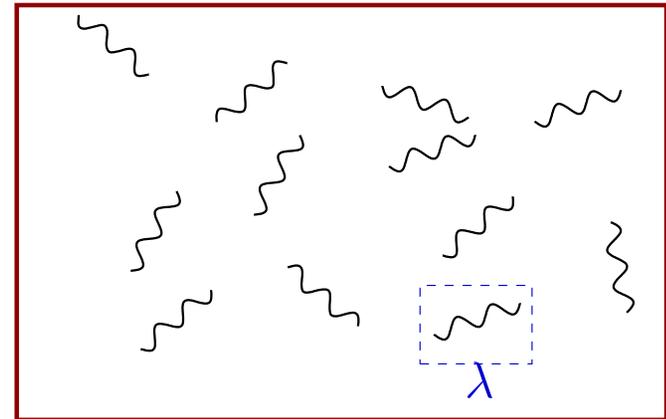
The light from the sun which is in equilibrium with the free electrons and protons of the sun



Real Life – Black Body



Theoretical Physicists Black Body



Photon Gas in Box

Problem on BB

- Estimate the energy and velocity of an electron in a 5000°K plasma like the sun:

Hint: Useful formulas $E \approx \frac{1}{2}mv^2$

Ignore the electrons potential energy

- Estimate the typical wavelength of a photon in a 5000°K plasma

Useful stuff you should know:

$$E = \frac{hc}{\lambda}$$

$$m_e = 0.5 \frac{\text{MeV}}{c^2} = \frac{0.5 \times 10^6 \text{ eV}}{(3 \times 10^8 \text{ m/s})^2}$$

$$k_B = \frac{1/40 \text{ eV}}{300^\circ\text{K}}$$

a unit of mass (mega electron volts) / (speed of light)²

$$hc = 1240 \text{ eV nm}$$

Answers:

(1a) $E \sim 0.41 \text{ eV}$ (1b) $v \sim \frac{1}{1000} c$ (2a) $E_{\text{phot}} \sim 1.15 \text{ eV}$ (2b) $\lambda \sim 1000$
↑
infrared

Solution

① Using

$$E_e \sim k_B T$$

$$\sim \frac{1/40 \text{ eV}}{300^\circ \text{K}} \times 5000^\circ \text{K}$$

a) $\sim 0.41 \text{ eV}$

Now

$$E = \frac{1}{2} m v^2 \Rightarrow v = \sqrt{\frac{2E}{m}} = \left(\frac{2 \times 0.41 \text{ eV}}{0.5 \times 10^{-6} \text{ eV}} \right)^{1/2} \frac{1}{(3 \times 10^8 \text{ m/s})^2}$$

Oops, I messed up here plugging the numbers into the calculator, replace 2.7 with 3.8

b) $= 2.7 \times 10^5 \text{ m/s}$

$$\sim \frac{1}{1000} \text{ of Light speed}$$

② For a photon in a 5000°K black body we have

$$E \approx 2.8 k_B T \approx 1.15 \text{ eV}$$

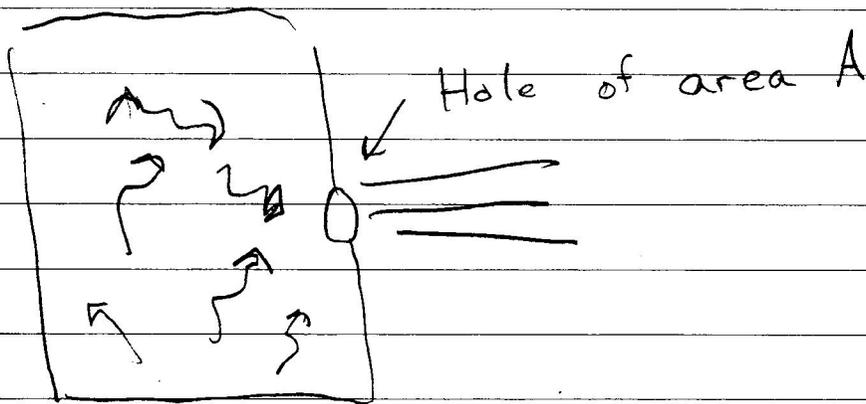
So

$$E = \frac{hc}{\lambda} \Rightarrow \lambda = \frac{hc}{E} = \frac{1240 \text{ eV nm}}{1.15 \text{ eV}} = 1078 \text{ nm}$$

↑
infrared

Radiation from Black Body

Now how much light energy does the box emit:



$$\frac{\Delta E}{A \Delta t} = \sigma T^4 \quad \text{or} \quad \frac{\Delta E}{\Delta t} = \sigma T^4 (\text{Area})$$

↑
the energy per area per time emitted
by the photon gas

$$\sigma = 5.6 \times 10^{-8} \frac{\text{W}}{\text{m}^2 \text{K}^4}$$

$$\sigma = \frac{2\pi^5}{15} c \frac{k_B^4}{(hc)^3} \leftarrow \text{an exact number}$$

Problem

- Estimate the energy per time
(The absolute luminosity) emitted by
the sun. Assume the sun is a 5800°K
black body
- Estimate the apparent luminosity of the sun
just outside the earth's atmosphere.
(If you have time)

Answers:

① $3.9 \times 10^{26} \text{ W}$ ② $l = 1380 \frac{\text{W}}{\text{m}^2}$

↑ agrees with observation

↑
 $\sim 1 \text{ kW per meter}^2$

Useful

$R_{\odot} \approx 700,000 \text{ km}$

Solution

$$(1) \frac{\Delta E}{\Delta t} = \sigma T^4 \times (\text{Area of Sun})$$

$$= 5.6 \times 10^{-8} \frac{\text{W}}{\text{m}^2} (5800^\circ\text{K})^4 \times (4\pi (7 \times 10^8 \text{m})^2)$$

$$\frac{\Delta E}{\Delta t} = 3.9 \times 10^{26} \text{ W}$$

(2) To find the energy per area on earth (outside atmosphere)

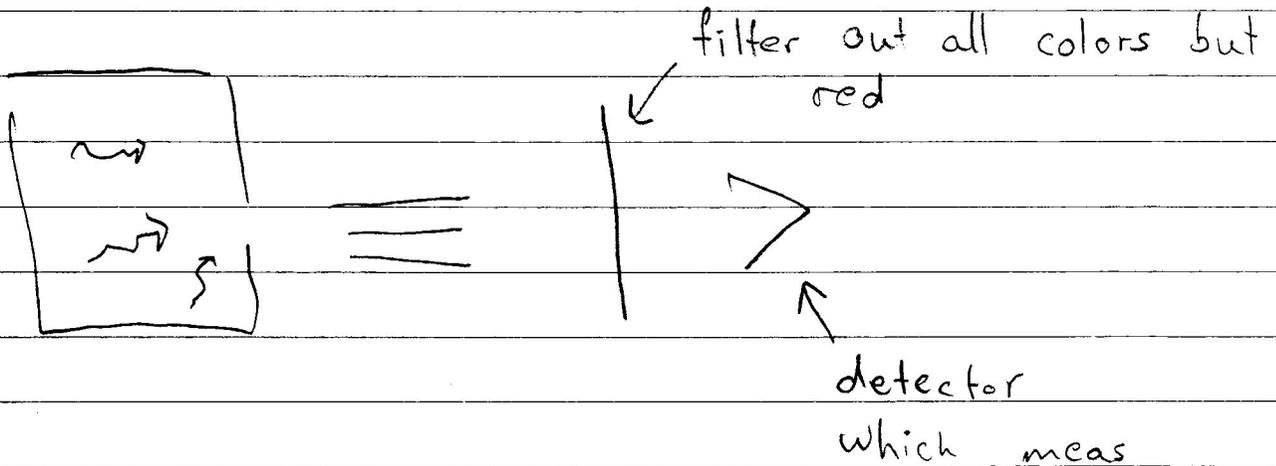
$$I = \frac{L}{4\pi R^2} = \frac{3.9 \times 10^{26} \text{ W}}{4\pi (1.5 \times 10^8 \text{ km})^2}$$

↑
distance to sun

$$I = 1380 \frac{\text{W}}{\text{m}^2}$$

Frequency dependence of BB radiation

Then we started talking about the color of the light



• In reality a filter will filter in a certain range, for example:

$$620\text{nm} \lesssim \lambda \lesssim 650$$

$$2.0\text{ eV} \lesssim hf \lesssim 2.1\text{ eV}$$

$$\Delta f \approx 0.1\text{ eV} \quad \text{for this case}$$

Then

$$\frac{\Delta E}{A \Delta t} = I_f(\tau) \Delta f$$

power per area

power per area per frequency

frequency of detector

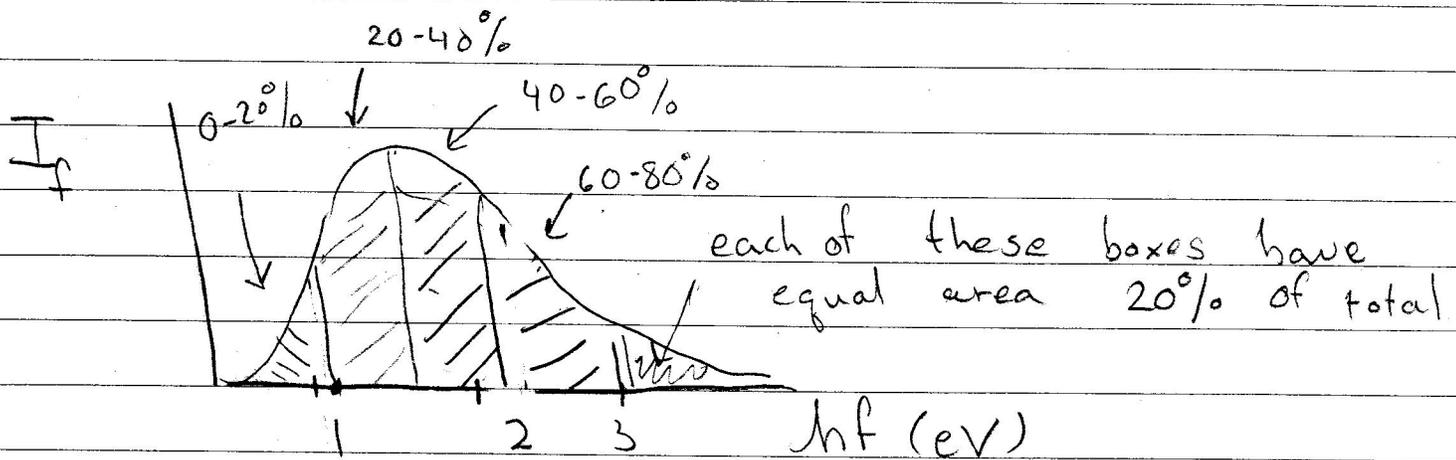
Frequency Dependence of Black Body Spectrum (continued)

Last time we stated w/out proof :

$$I_f \propto \frac{f^3}{e^{hf/k_B T} - 1}$$

• See slide :

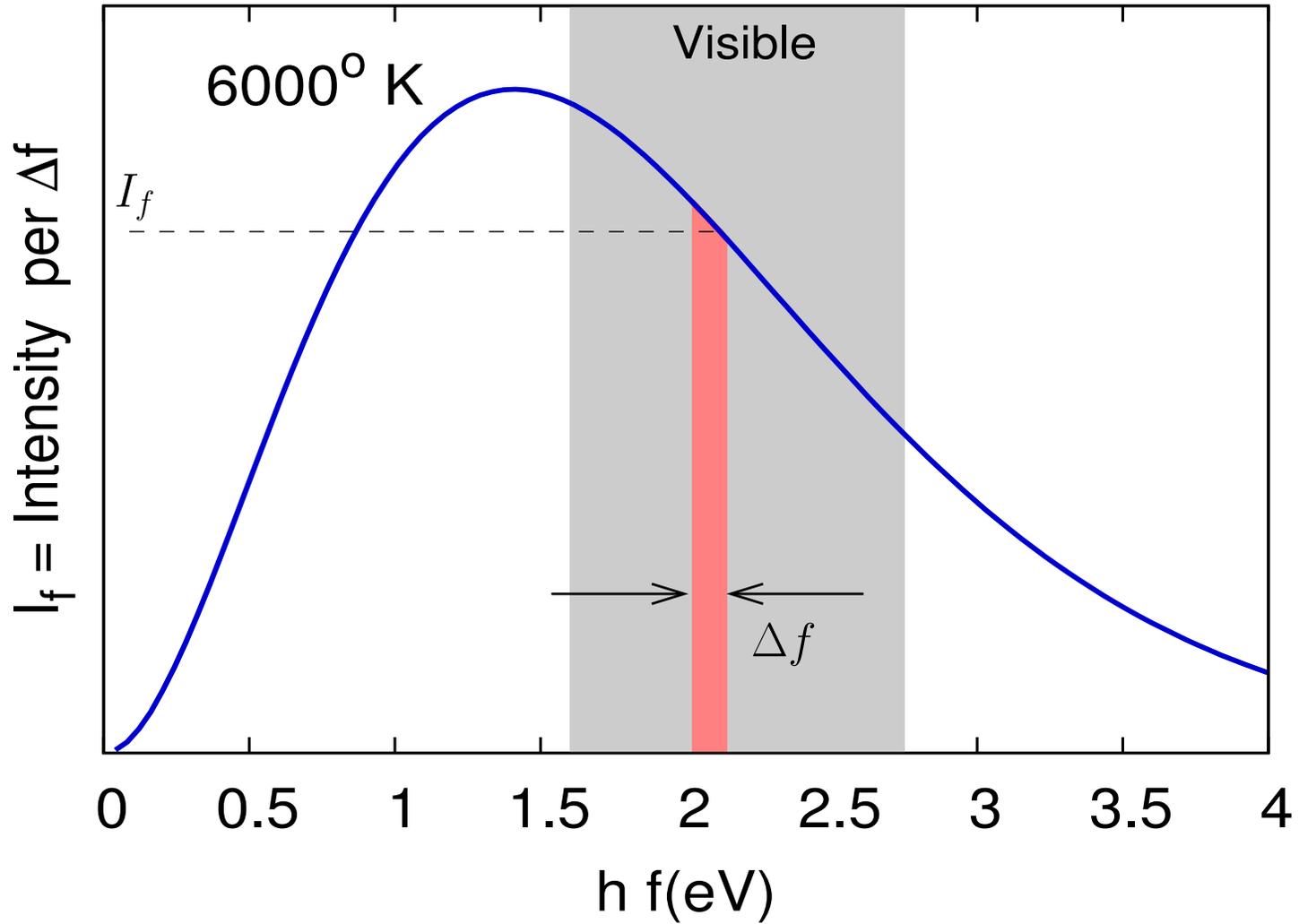
• See slide



• Thus 20% of the energy is carried by photons with energy less than 1 eV

• And 20% of the energy is carried by photons with $E > 2.7$ eV.

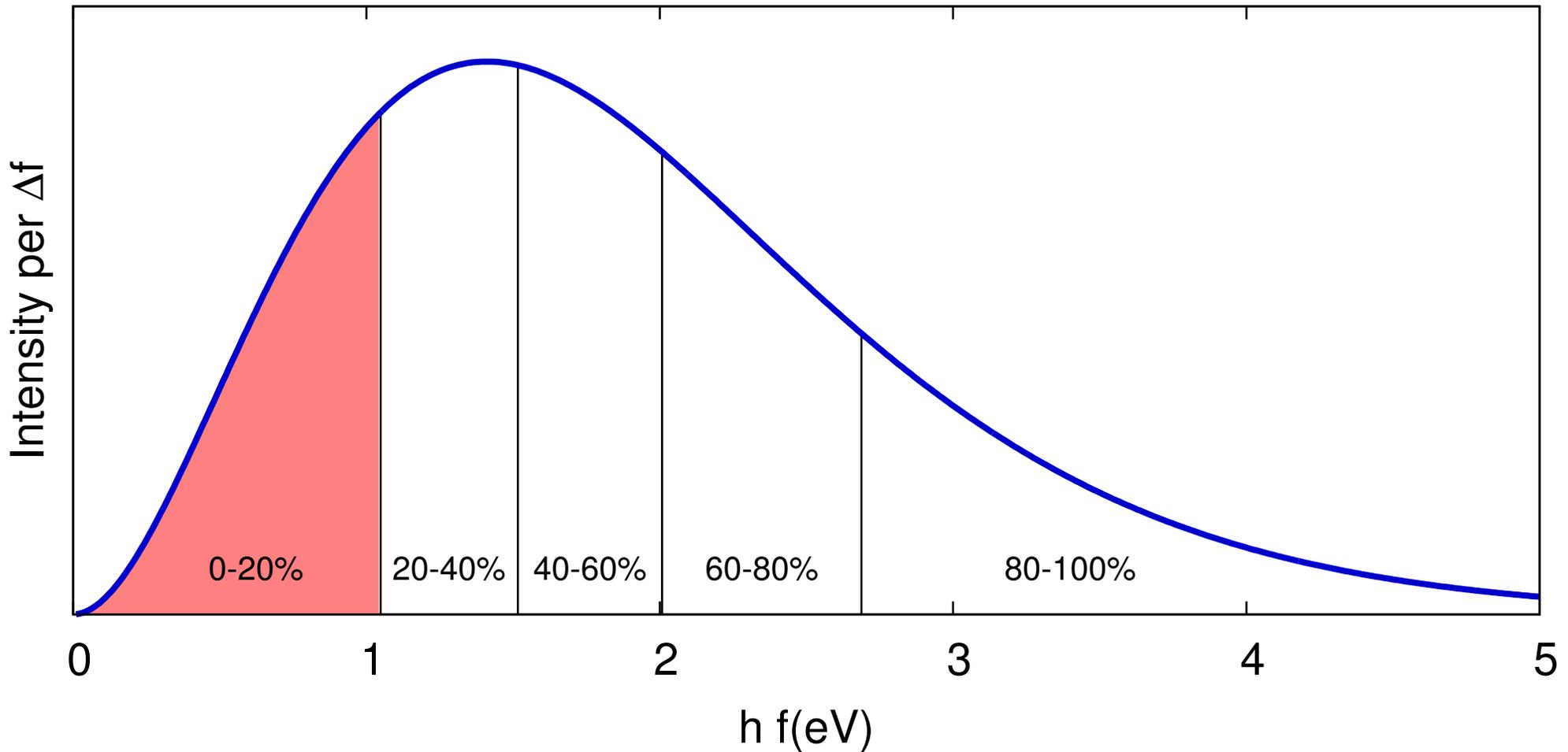
Black Body Radiation



$$\frac{\Delta E}{A \Delta t} = I_f \times \Delta f$$

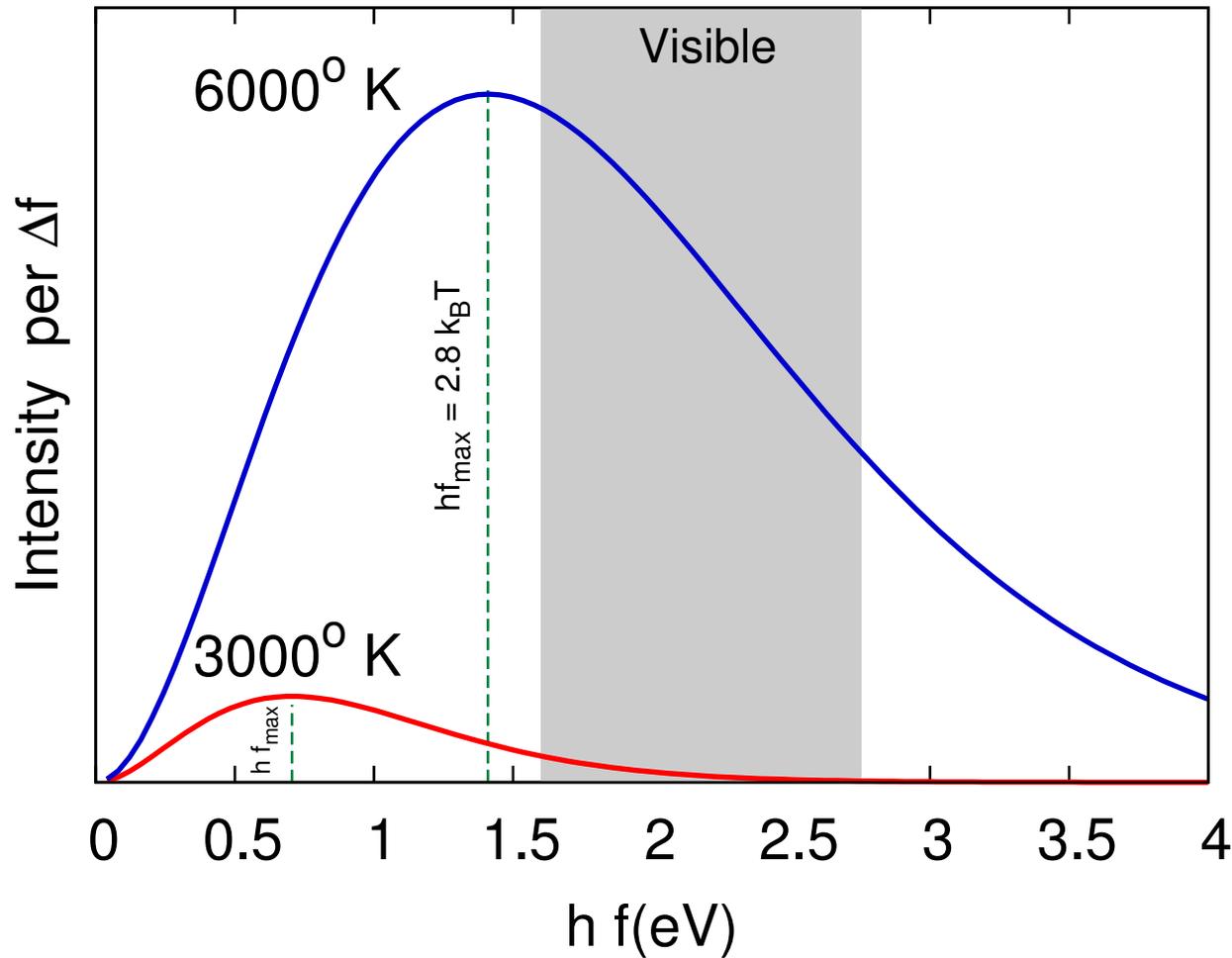
The black black body spectrum records how much light is carried by a given frequency

Black Body Radiation at 6000^o K



Each of these frequency ranges have equal area under the curve and carry 20% of the energy

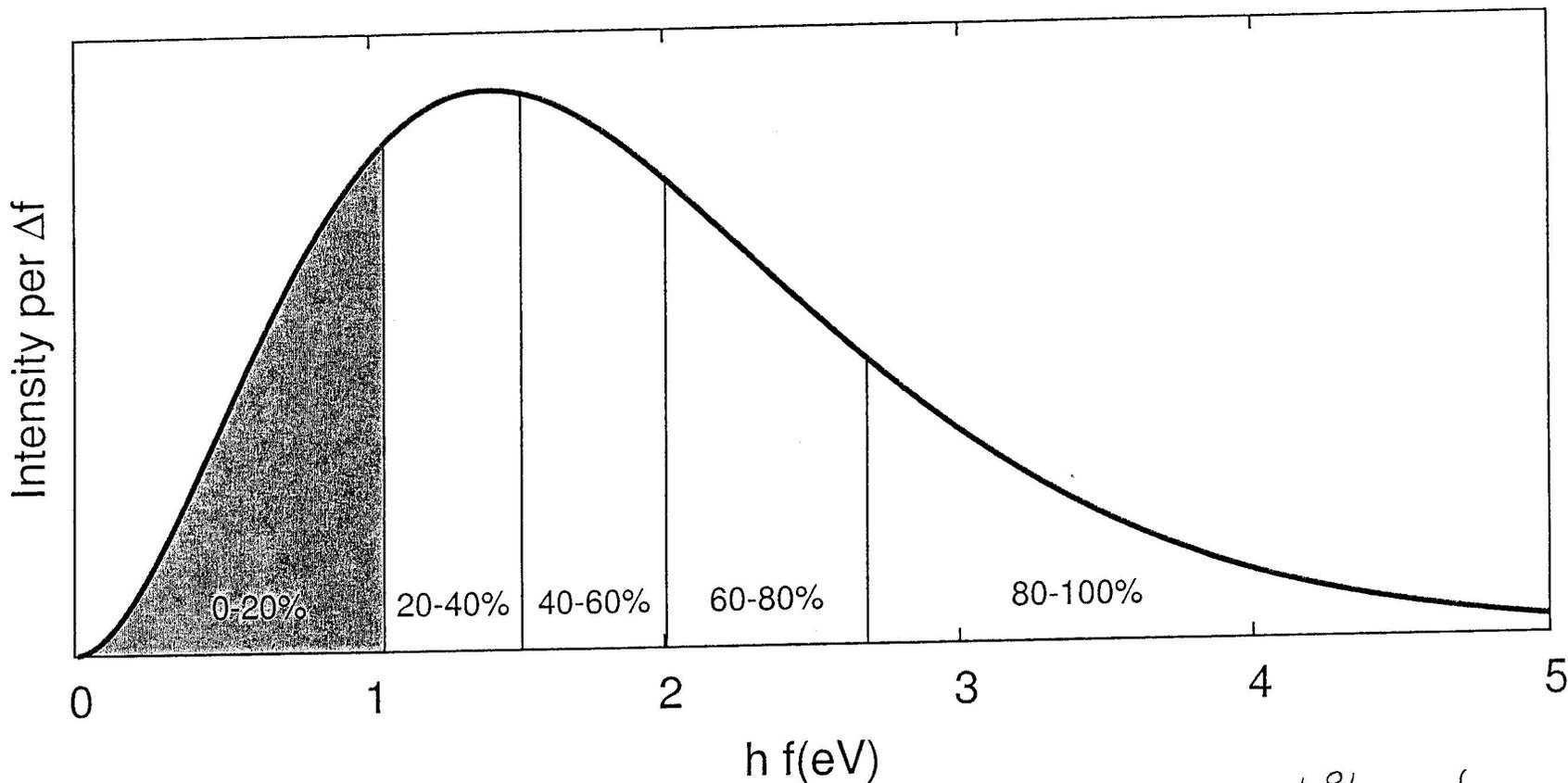
Black Body Radiation



1. The height increases rapidly with with temperature (as T^3).
2. The typical frequency (the base) increases linearly T .

Problem on the frequency-dep of BB radiation

Black Body Radiation at 6000°K

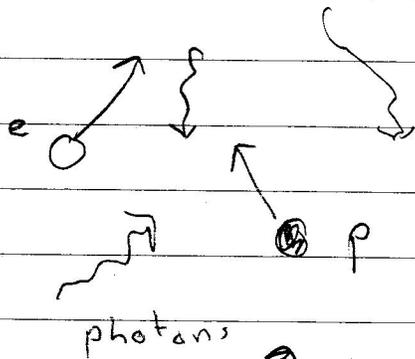


Problem: Photons with energy $> E_*$ carry only 10% of the energy. Estimate E_* .

High Temperature Plasma & Recombination

$$k_B T \sim 10 \text{ eV}$$

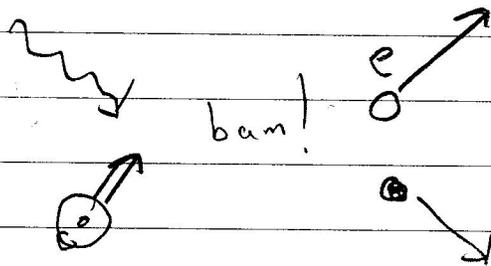
High
Temperature



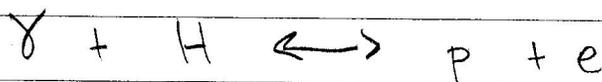
The number of electrons and protons are equal

Very occasionally make Hydrogens

It takes a certain amount of energy to rip an electron from hydrogen $\Delta E = 13.6 \text{ eV}$



This is known as "photo-dissociation."

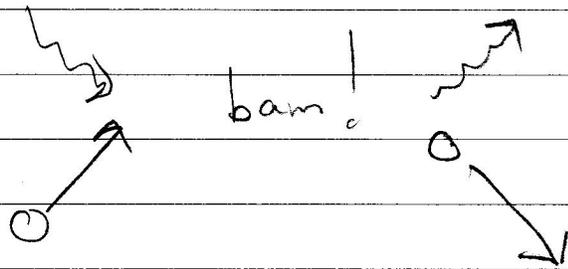


In high temperature plasma there are plenty of photons @ energy greater than 13.6 eV

Recombination Continued

known as Compton scattering

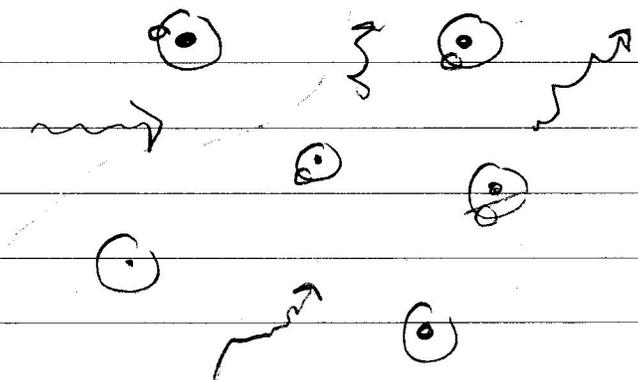
- There is also photon-electron scattering which happens a lot, keeping the plasma in equilibrium.



As the plasma cools to $k_B T \sim 0.2 \text{ eV}$
there are almost no photons (ω) energy $> 13.6 \text{ eV}$

→ less than one in a million have $E > 13.6 \text{ eV}$

- As a result the photons can no longer ionize the hydrogen gas



• This is known
as recombination

Recombination Continued

The free electron fraction



free electrons



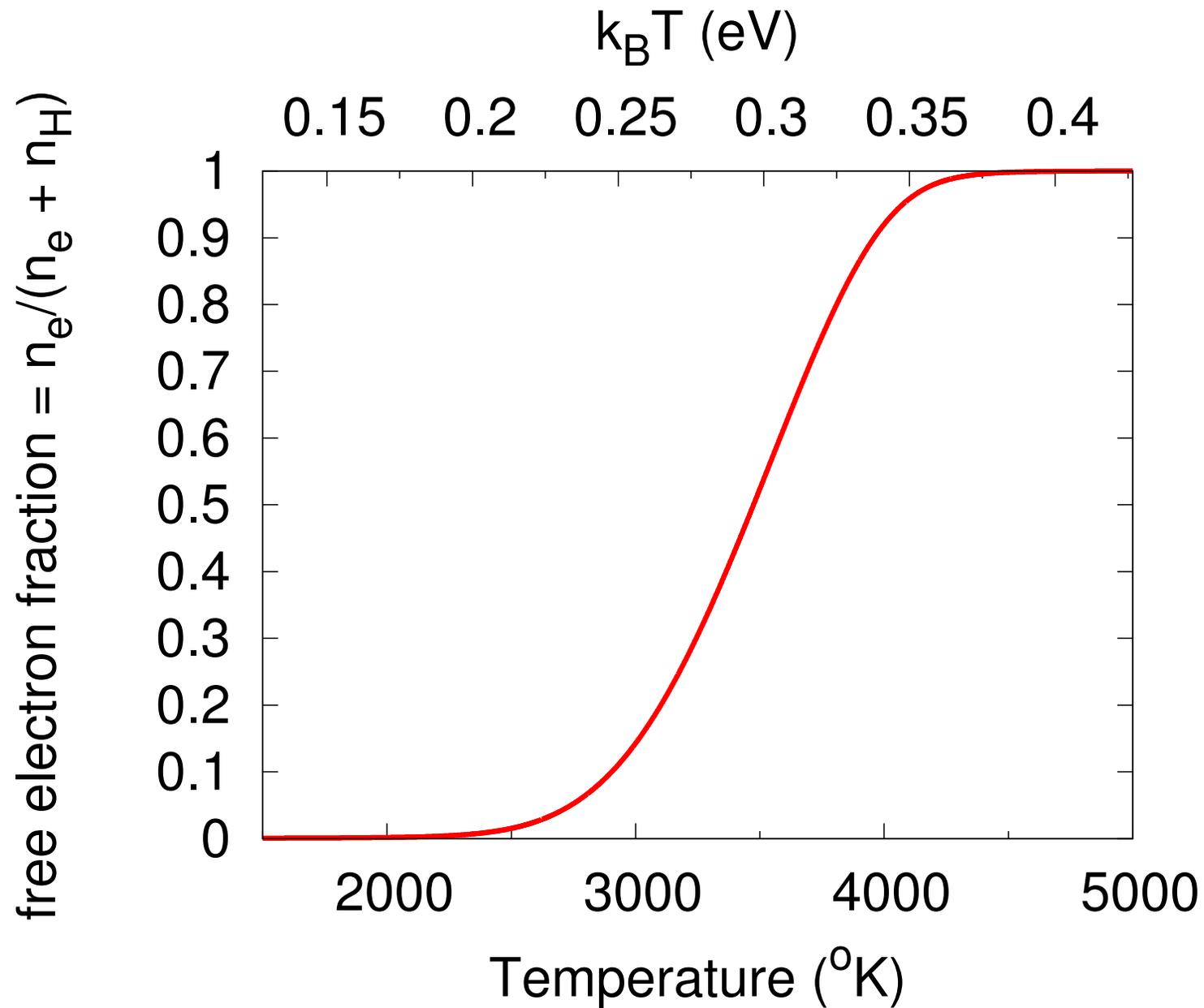
Bound electron



$$X_e = \text{fraction of electrons free} = \frac{n_e}{n_e + n_H}$$

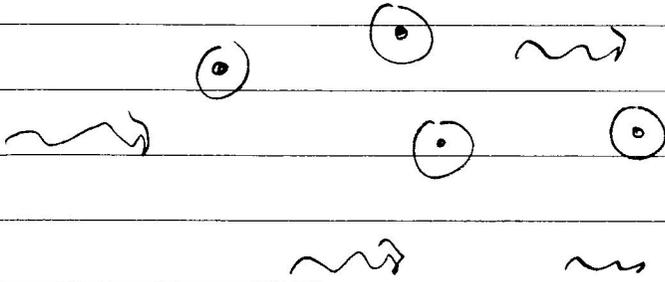
See slide for plot of X_e vs Temperature

The free electron fraction in high temperature plasma:



Recombination Continued

- At this point the photons are effectively free, since electron-photon (or Compton) scattering has stopped.



- The light scatters only very weakly off neutral hydrogen (at least if the energy is much less than 13.6 eV as is the case here)
- Then the light flies freely until today!!
- What is the spectrum of light that is seen today?

Computing The Spectrum Today

- Definitions:

T_r = the temperature at recombination

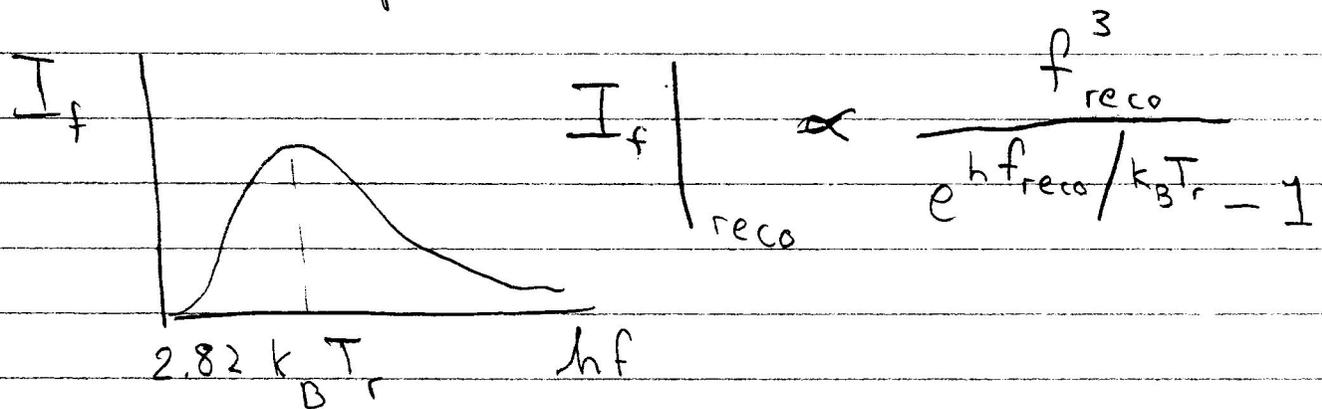
a_r = the scale factor at reco.

t_r = the time at reco

f_r and λ_r = the frequency and wavelength at reco

Computing The Spectrum of Temperature Today

At recombination, the spectrum (i.e. the probability to find a photon with frequency f) is given by a black body radiation spectrum @ $T \approx 3000^{\circ}K$, T_r is the temperature at recombination.



• Now the universe continues to expand
The wave lengths all increase by a

uniform factor

$$\lambda_0 = \frac{\lambda_{reco}}{a_{reco}}$$

λ_0 ← wavelength today
 λ_{reco} ← wavelength at recombination
 a_{reco}

$$f_0 = a_{reco} f_{reco}$$

f_0 ← frequency today
 f_{reco}

Computing the spectrum today continued:

$$I_f \Big|_{\text{today}} \propto \frac{(f_0/a)^3}{e^{hf_0/k_B a T_r} - 1}$$

$$\propto \frac{f_0^3}{e^{hf_0/k_B T_0} - 1}$$

Where

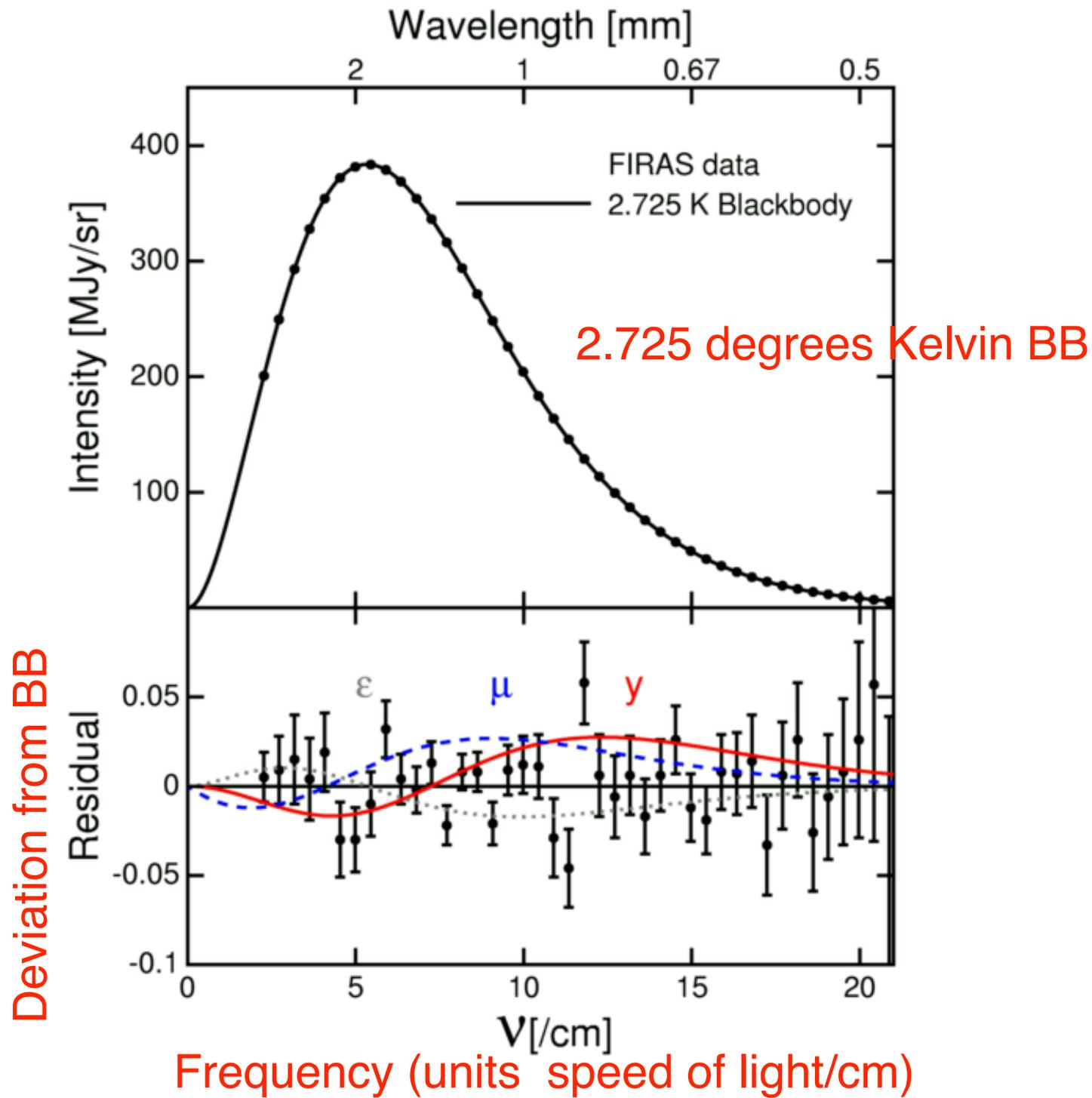
$$T_0 = a_r T_r$$

Thus we expect a blackbody spectrum today with a lower temperature

$$T_0 = a_r T_r$$

- The temperature today is observed to be

$$T_0 \approx 2.728 \text{ K}$$



History:

- The basic description of recombination was Gamov, Alpher, and Herman \sim 1940's

→ They estimated that the scale factor at recombination was

$$1/a_r \sim 100 \leftrightarrow 1000$$

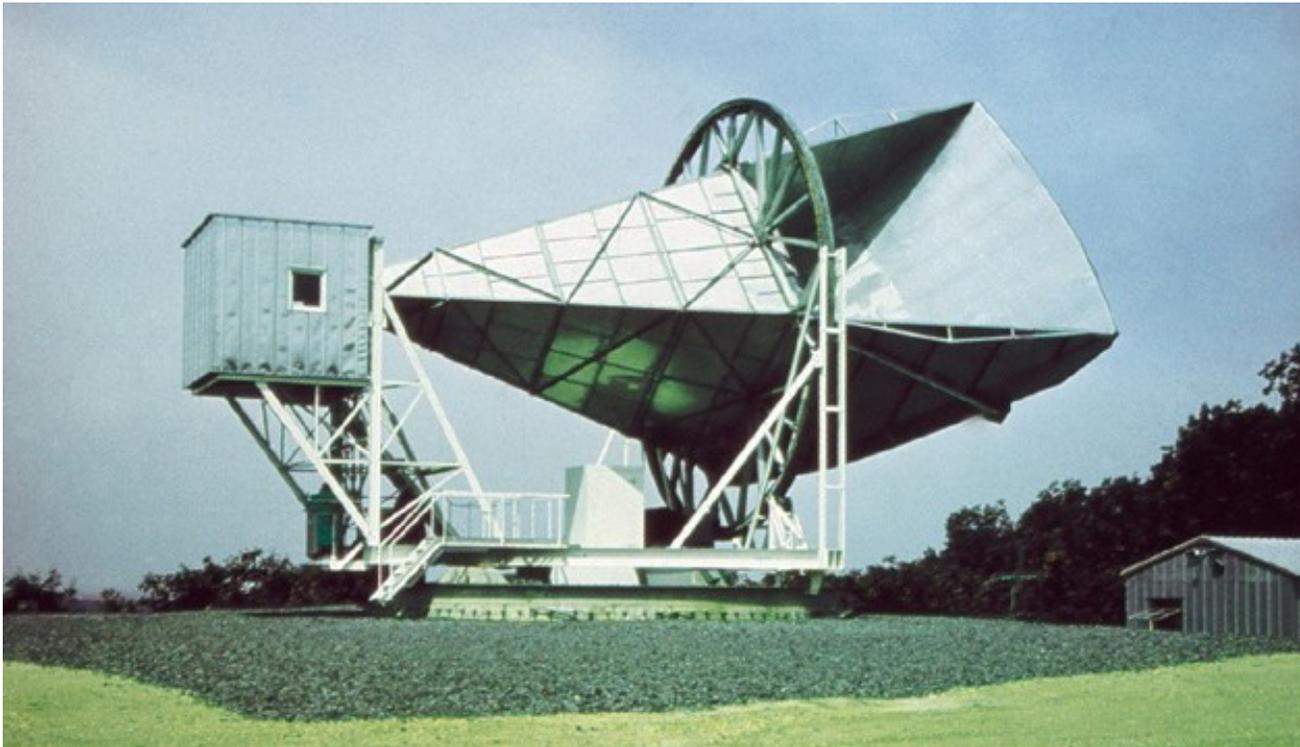
This leads to a prediction that with $T_r \approx 3000^\circ\text{K}$, the temperature today would be

$$T_0 \sim 3^\circ \leftrightarrow 30^\circ\text{K}$$

History Continued

- The original prediction of the microwave background lay dormant for 20 yrs
 - An important paper ~ 1960 by Jim Peebles revived interest.
 - Dicke, Peebles, + Wilkenson set out to measure the micro waves (Princeton)
 - At the same time, "Penzias and Wilson discovered" the microwave background by accident 1965
- Penzias and Wilson were awarded the Nobel Prize.

Penzias and Wilson discover the Black Body Radiation in Crawford Hill, NJ. 1965



These bell Labs astronomers found a mysterious microwave signal
Princeton astronomers Dicke, Peebles, and Wilkenson describe the origin of the signal

The time and Scale factor at recombination:

Then, we also can estimate the scale factor at recombination:

$$T_0 = a_r T_r$$

↑
2.72°K

↑
3400°K

So $1/a_r = (1+z) = \frac{T_r}{T_0} = 1250$
(See next slide)

Then we can also estimate the time at recombination from $a(t)$

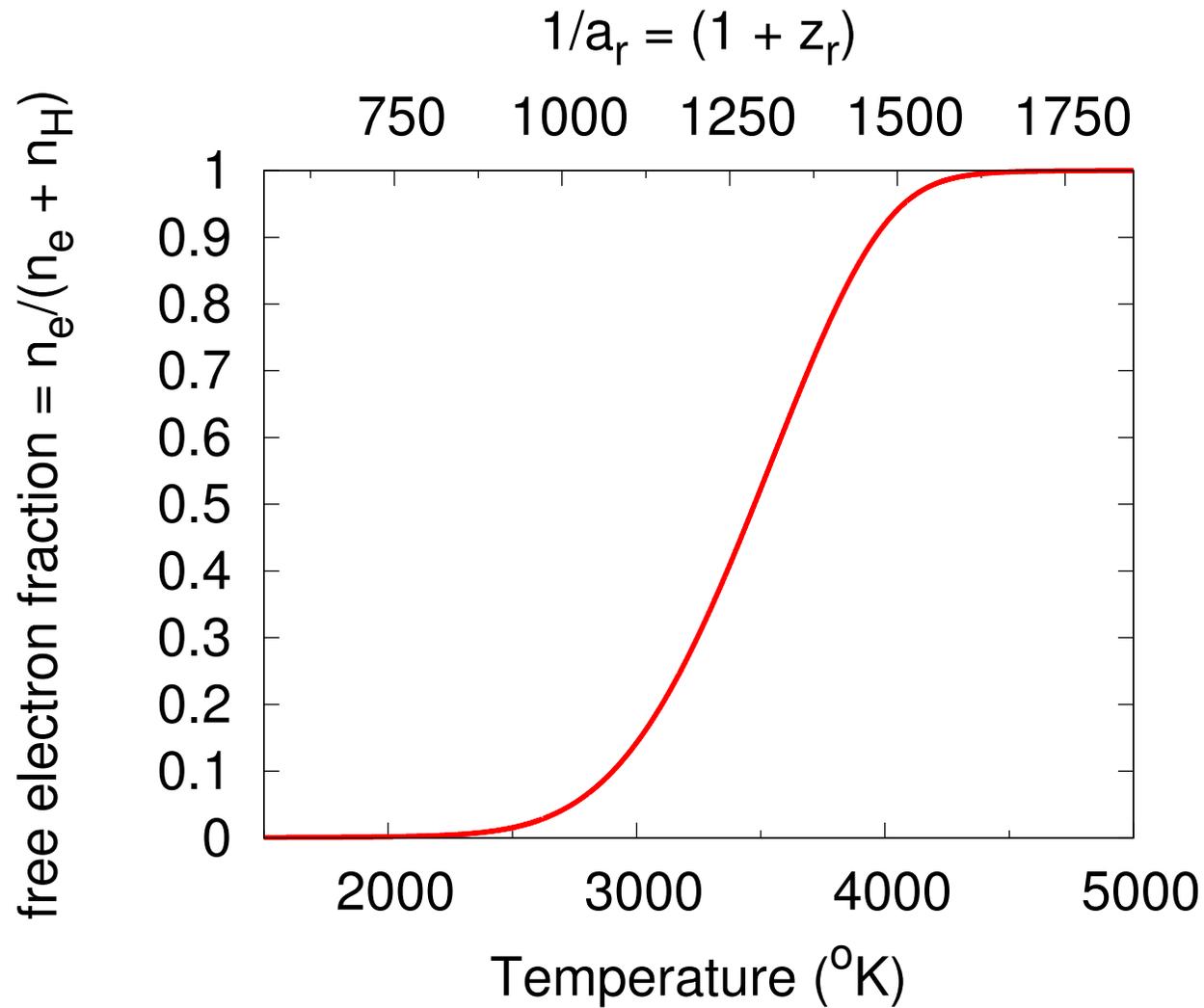
$$a_r(t_r) = \frac{1}{1250}$$

So from the graph of $a(t)$ find

$$t_r \approx 400,000 \text{ y}$$

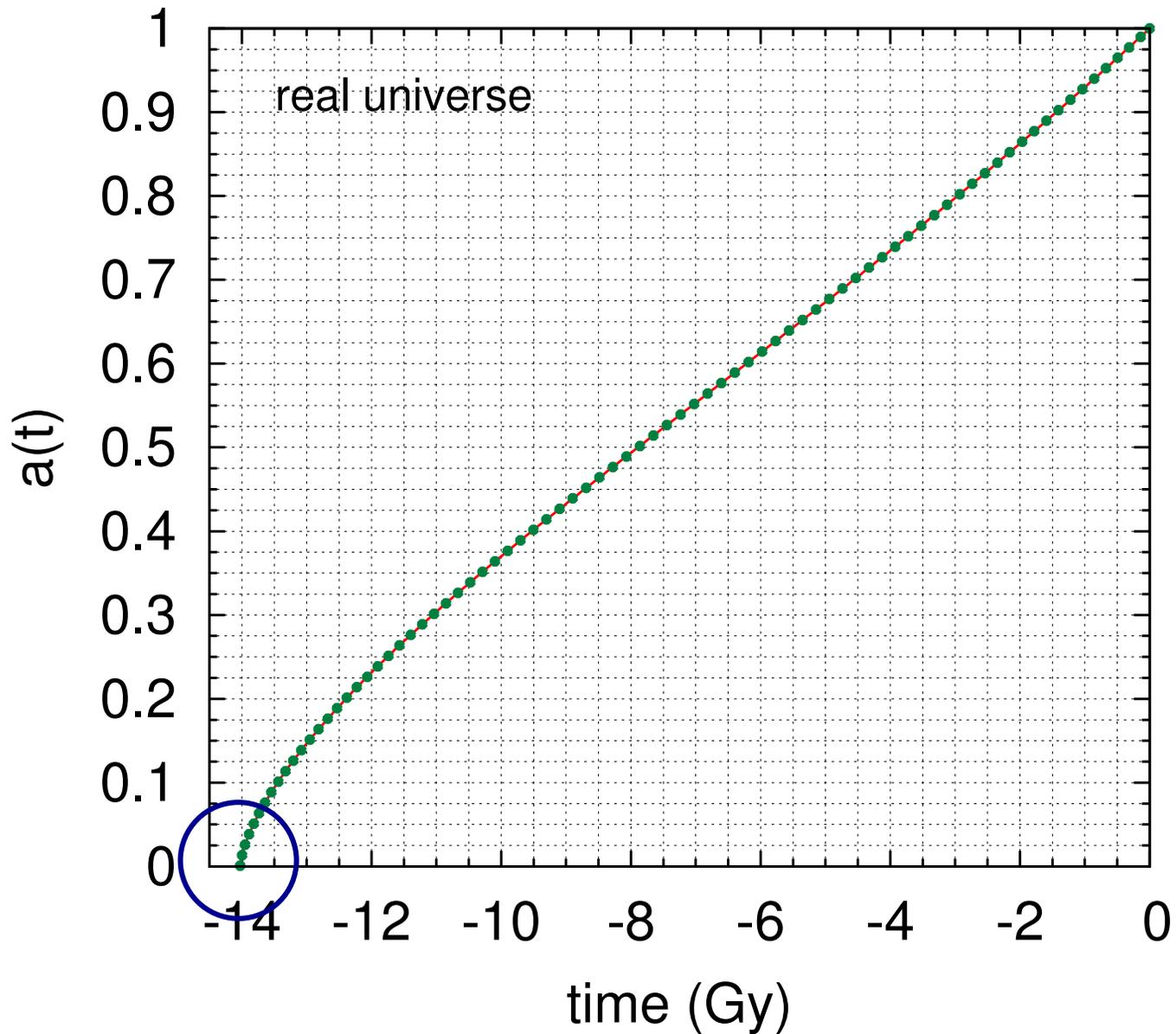
(See slides, 2+3 from now)

The red shift (scale factor) at recombination:



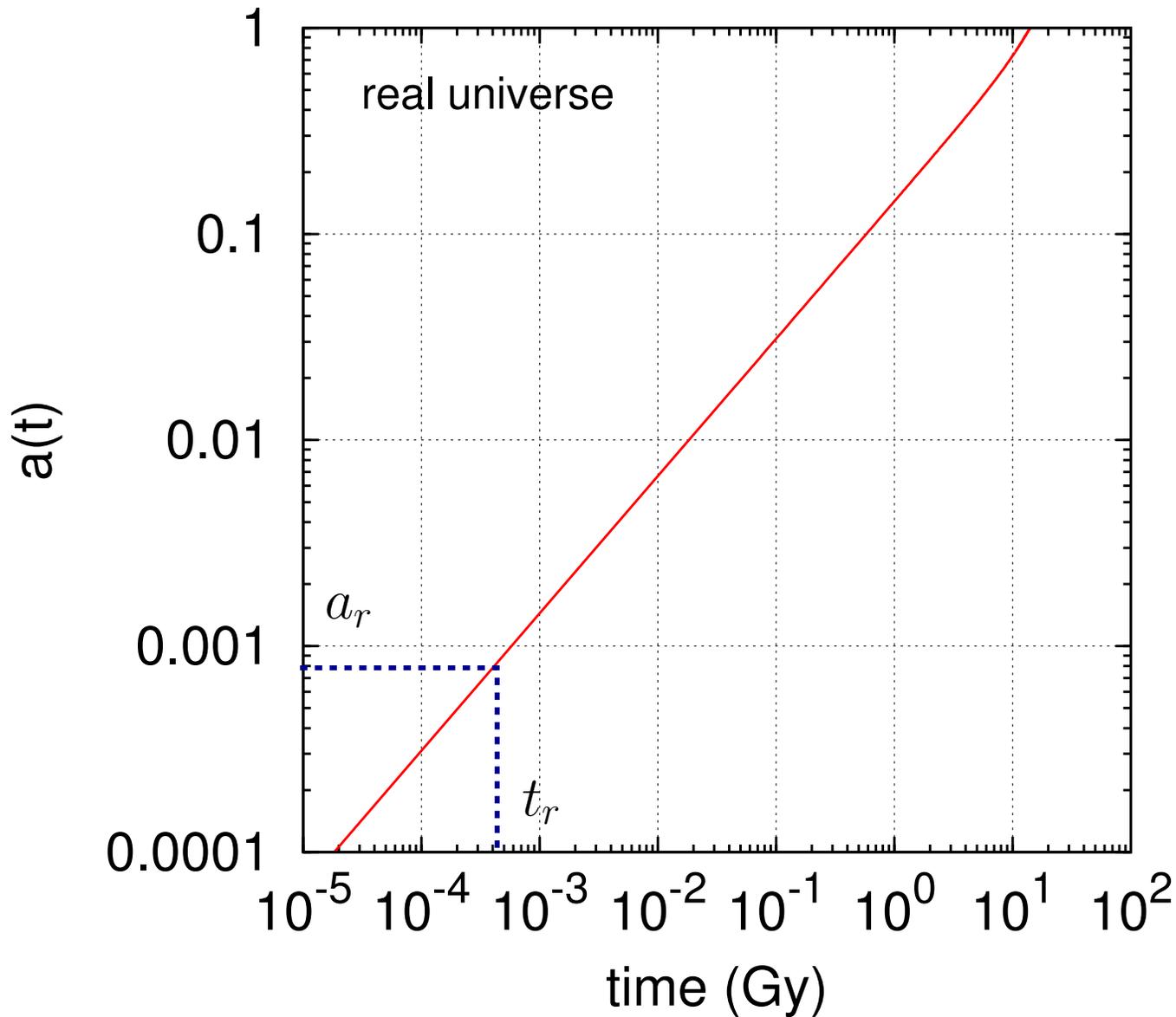
$$\frac{1}{a_r} = (1 + z_r) = \frac{T_r}{T_o} = \frac{3400^{\circ}\text{K}}{2.72^{\circ}\text{K}} \simeq 1250$$

Given $a(t)$ can estimate the time (after the big bang) of recombination



We are looking for $a(t) \simeq 1/1250$. This graph wont work

Given $a(t)$ (on a log scale) can estimate the time after the big bang of recombination



For $a_r \simeq 1/1250$, the time is $t_r \simeq 4 \times 10^{-4}$ Gy $\simeq 400,000$ years

The uniformity of the Cosmic Microwave Background

- The microwave background is remarkably uniform throughout the sky

$$\bar{T} = 2.728 \text{ K} \quad (\text{see slide})$$

In You subtract the average, and look at the difference:

$$T - \bar{T} \approx 3.353 \text{ mK}$$



Last digit of 2,728 K

$$\pm 0.00353 \text{ mK}$$

Then you begin to see structure to CMB:

This ^{uniformity} is another example of the cosmic principle:

"On the largest distance scales, everything is the same. All observers measure the same thing"

Two Examples:

① Cosmic Microwave Background

② The "end of greatness" in the distribution of galaxies

