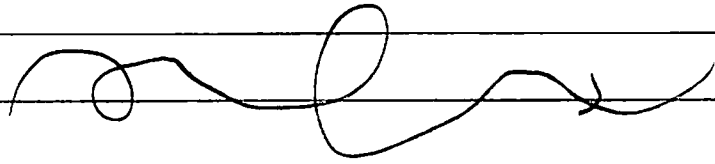


Last Time

Started From wave equation

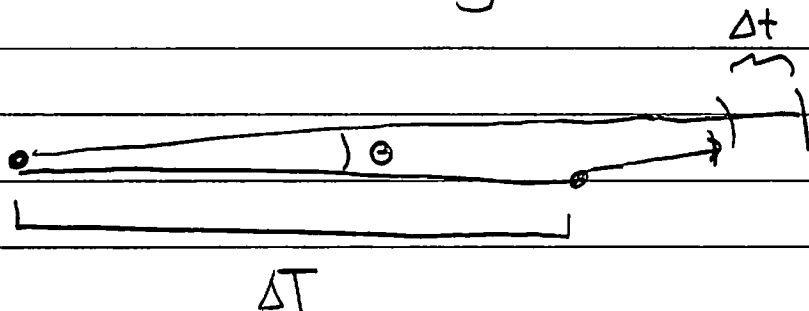


Derived the radiation for arbitrary relativistic motion

$$A_{\text{rad}} = \frac{1}{4\pi r} \frac{q \vec{v}(T)/c}{(1 - n \cdot \beta)}$$

$$A_{\text{rad}} = \frac{q}{4\pi} \frac{\vec{v}(T)/c}{dt} \quad T = t - \frac{r}{c} + \frac{\vec{n} \cdot \vec{r}_0(T)}{c}$$

Again dT/dt is very important



if a waveform at infinity is observed to have a characteristic time scale Δt , then it was formed over a time-scale $\Delta T = \frac{\Delta t}{(1 - n \cdot \beta)}$

Last Time Continued

Once we know A_{rad} we can find the electric field

$$\vec{E}_{\text{rad}} = \frac{e}{4\pi\epsilon_0} \frac{n \times (n - \vec{\beta}) \times \vec{a}(\tau)}{(1 - n \cdot \vec{\beta})^3} = n \times n \times \frac{1}{c} \frac{\partial \vec{A}}{\partial t}$$

And we can find the radiated power

$$\frac{dW}{dt d\Omega} = \frac{dP(t)}{d\Omega} = c |E(t)|^2$$

Often care about the energy lost during a period of acceleration

$$\frac{dW}{dT d\Omega} = \frac{dP(\tau)}{d\Omega}$$

$$= c |E|^2 \frac{dt}{dT}$$

$$= c |E|^2 (1 - n \cdot \vec{\beta}(\tau))$$

Problem - SHO

- ① A relativistic harmonic oscillator with charge q , moves on the z -axis

$$z = H \cos \omega_0 t$$

Show that the power emitted is

$$\frac{dW}{dT d\Omega} = \frac{dP(T)}{d\Omega} = \frac{e^2}{16\pi^2} \frac{c\beta^4}{H^2} \frac{\cos^2 \omega_0 t \sin^2 \theta}{(1 + \beta \cos \theta \sin \omega_0 t)^5}$$

where $\beta \equiv \frac{\omega_0 H}{c}$, $\gamma \equiv \frac{1}{\sqrt{1 - \beta^2}}$

- ② Most of the energy in the $\gamma \rightarrow \infty$ limit is radiated at particular points during its motion. Determine these points. The radiation is pulsed. Explain why.

- ③ Show that for $\gamma \rightarrow \infty$ for $T = \frac{3\pi}{2\omega_0} + \Delta T$, have: small

$$\frac{dP}{d\Omega}(T) = \frac{2e^2}{\pi^2} \frac{c\beta^4}{H^4} \gamma^6 \left[\frac{(\gamma\theta)^2 (\gamma\omega_0 \Delta T)^2}{(1 + (\gamma\theta)^2 + (\gamma\omega_0 \Delta T)^2)^5} \right]$$

- ④ Use the integral

$$\int_{-\infty}^{\infty} \frac{x^2 dx}{(1 + y^2 + x^2)^5} = \frac{5\pi}{128} \frac{1}{(1 + y^2)^{7/2}}$$

Problem - SHO continued

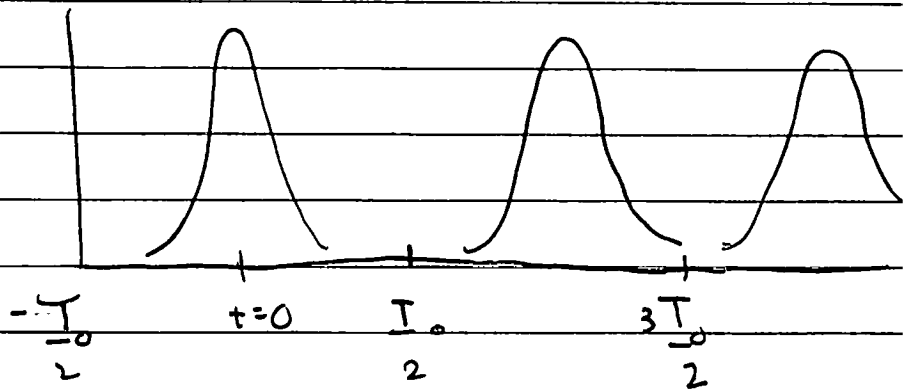
to show that the time averaged power is

$$\frac{d\overline{P}}{d\Omega} = \frac{e^2}{128\pi^2} \frac{c\beta^4}{H^2} \gamma^5 \left[\frac{5(\gamma\theta)^2}{(1+(\gamma\theta)^2)^{7/2}} \right]$$

A repeating sequence of pulses:

$E_1(t) \equiv$ The electric field in a single period,

zero elsewhere.

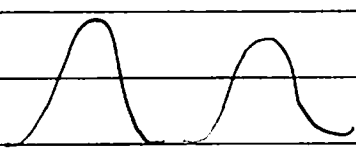


$E_2(\omega) =$ Fourier transform $E_1(t)$

Problem 1

Suppose that the waveform repeats once.

Show that the Fourier transform is



$$E_2(\omega) \equiv E_1(\omega) (1 + e^{i\omega T_0})$$

And the power spectrum is

$$|E_2(\omega)|^2 = |E_1(\omega)|^2 (2 + 2\cos(\omega T_0))$$

Problem 2

Show that the Electric field after a sum of n -pulses

is A repeating Sequence of pulses pg. 2

$$E_N(\omega) = E_1(\omega) e^{-i\omega(n-1)T_0/2} \frac{\sin(\omega n T_0/2)}{\sin \omega T_0/2}$$

So that the power is

$$|E_N(\omega)|^2 = |E_1(\omega)|^2 \left(\frac{\sin(\omega n T_0/2)}{\sin \omega T_0/2} \right)^2$$

Problem 3

Show that the electric field approaches

$$E(\omega) \xrightarrow{N \rightarrow \infty} \sum_m E_1(\omega_m) \frac{2\pi}{T_0} \delta(\omega - \omega_m)$$

and

$$\frac{|E(\omega)|^2}{nT_0} \xrightarrow{N \rightarrow \infty} \sum_m |E_1(\omega_m)|^2 \frac{2\pi}{T_0} \delta(\omega - \omega_m)$$

Use

$$\int_{-\infty}^{\infty} \frac{\sin(x/2)}{x/2} dx = 2\pi$$

and

$$\int_{-\infty}^{\infty} \left(\frac{\sin(x/2)}{x/2} \right)^2 dx = 2\pi$$

Problem 4

Use the results of this problem to claim

$\Delta_T(t) \equiv \sum_n \delta(t - nT_0)$, the int's fourier transform is

$$\Delta_T(\omega) = \sum_n e^{i\omega n T_0} = \frac{2\pi}{T_0} \sum_m \delta(\omega - \omega_m) \quad \omega_m = \frac{2\pi}{T_0}$$

And

$$\frac{dP_m}{d\Omega} = \frac{1}{(\text{Time})} \int_{\omega_m - \text{bit}}^{\omega_m + \text{bit}} 2\pi \frac{dW}{d\omega} \frac{d\omega}{2\pi}$$

$$\frac{dP_m}{d\Omega} = \frac{C}{T_0^2} |E_1(\omega_m)|^2$$