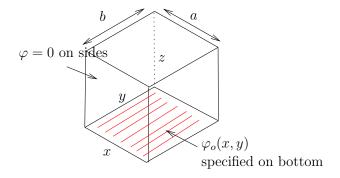
D Separation of Variables

D.1 Cartesian coordinates



(a) Laplacian

$$\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2} + \frac{\partial^2}{\partial z^2}\right)\Phi = 0 \tag{D.1}$$

(b) Eigen functions along boundary vanishing at x = 0 and x = a and y = 0 and y = b

$$\psi_{nm}(x,y) = \sin\left(\frac{n\pi x}{a}\right)\sin\left(\frac{m\pi y}{b}\right) \qquad n = 1\dots\infty \qquad m = 1\dots\infty$$

(c) Orthogonality

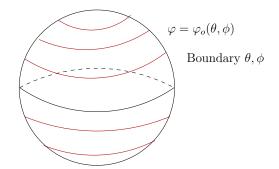
$$\int_0^a dx \int_0^b dy \,\psi_{nm} \,\psi_{n'm'} = \left(\frac{a}{2}\right) \left(\frac{b}{2}\right) \delta_{nn'} \delta_{mm'}$$

(d) Solution

$$\sum_{n=1}^{\infty} \sum_{m=1}^{\infty} \left[A_{nm} e^{-\gamma_{nm}z} + B_{nm} e^{+\gamma_{nm}z} \right] \psi_{nm}(x,y) \tag{D.2}$$

where $\gamma_{nm} = \sqrt{(n\pi/a)^2 + (m\pi/b)^2}$

D.2 Spherical coordinates



(a) Laplacian

$$\left[\frac{1}{r^2}\frac{\partial}{\partial r}r^2\frac{\partial}{\partial r} + \frac{1}{r^2\sin\theta}\frac{\partial}{\partial\theta}\sin\theta\frac{\partial}{\partial\theta} + \frac{1}{r^2\sin^2\theta}\frac{\partial^2}{\partial\phi^2}\right]\Phi = 0$$
(D.3)

(b) Eigen functions along boundary θ, ϕ , regular at $\theta = 0$ and $\pi, 2\pi$ periodic in ϕ

$$\psi_{\ell m}(\theta,\phi) = Y_{\ell m}(\theta,\phi) \qquad \ell = 0\dots\infty \quad m = -\ell\dots\ell$$

(c) Orthogonality:

$$\int d\Omega Y_{\ell m}^*(\theta,\phi) Y_{\ell'm'}(\theta,\phi) = \delta_{\ell\ell'} \delta_{mm'}$$

(d) Solution

$$\Phi = \sum_{\ell=0}^{\infty} \sum_{m=-\ell}^{\ell} \left[A_{\ell m} r^{\ell} + \frac{B_{\ell m}}{r^{\ell+1}} \right] Y_{\ell m}$$
(D.4)

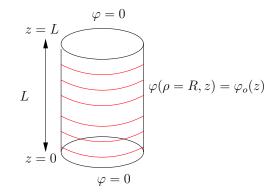
(e) When there is no azimuthal dependence things simplify to

$$\Phi = \sum_{\ell=0}^{\infty} \left[A_{\ell} r^{\ell} + \frac{B_{\ell}}{r^{\ell+1}} \right] P_{\ell}(\cos \theta)$$
(D.5)

where $P_{\ell}(\cos \theta)$ is the legendre polynomial, which up to a normalization if $Y_{\ell 0}(\theta, \phi)$, satisfying the orthogonality

$$\int_{-1}^{1} d(\cos\theta) P_{\ell}(\cos\theta) P_{\ell'}(\cos\theta) = \frac{2}{2\ell+1} \delta_{\ell\ell'}$$

D.3 Cylindrical Boundary: z, ϕ are the boundary.



(a) Laplacian:

$$\left[\frac{1}{\rho}\frac{\partial}{\partial\rho}\rho\frac{\partial}{\partial\rho} + \frac{1}{\rho^2}\frac{\partial^2}{\partial\phi^2} + \frac{\partial^2}{\partial z^2}\right]\Phi = 0 \tag{D.6}$$

(b) Eigenfunctions along boundary z, ϕ vanishing at z = 0 and z = L and 2π periodic in ϕ

$$\psi_{nm}(z,\phi) = \sin(k_n z) e^{im\phi} \qquad k_n \equiv \frac{n\pi}{L} \qquad n = 1\dots\infty \quad m = -\infty\dots\infty$$

(c) Orthogonality:

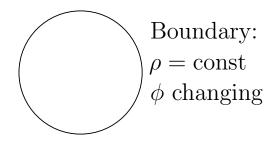
$$\int_{0}^{L} dz \int_{0}^{2\pi} \psi_{nm}(z,\phi) \,\psi_{nm}(z,\phi) = \frac{L}{2} \,(2\pi) \delta_{nn'} \delta_{mm'}$$

(d) Solution:

$$\Phi = \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} \left[A_{nm} I_m(k_n \rho) + B_{nm} K_m(k_n \rho) \right] \psi_{nm}(z, \phi)$$
(D.7)

Here $I_{\nu}(x)$ and $K_{\nu}(x)$ is the modified bessel function of the first and second kinds. Note that $K_{-m}(x) = K_m(x)$ and $I_{-m}(x)$

D.4 2D cylindrical coordinates



(a) Laplacian:

$$\left[\frac{1}{\rho}\frac{\partial}{\partial\rho}\rho\frac{\partial}{\partial\rho} + \frac{1}{\rho^2}\frac{\partial^2}{\partial\phi^2}\right]\Phi = 0 \tag{D.8}$$

(b) Eigenfunctions along boundary $\phi : \, 2\pi$ periodic in ϕ

 $\psi_m(\phi) = e^{im\phi} \qquad m = -\infty \dots \infty$

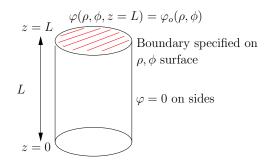
(c) Orthogonality

$$\int_{0}^{2\pi} \psi_{m}^{*}(\phi)\psi_{m'}(\phi) = 2\pi\delta_{mm'}$$
(D.9)

(d) Solution

$$\Phi = A_0 + B_0 \ln \rho + \sum_{m=-\infty}^{\infty} \left(A_m \rho^{|m|} + \frac{B_m}{\rho^{-|m|}} \right) \psi_m$$

D.5 Cylindrical Boundary: ρ, ϕ are the boundary



(a) Laplacian:

$$\left[\frac{1}{\rho}\frac{\partial}{\partial\rho}\rho\frac{\partial}{\partial\rho} + \frac{1}{\rho^2}\frac{\partial^2}{\partial\phi^2} + \frac{\partial^2}{\partial z^2}\right]\Phi = 0 \tag{D.10}$$

(b) Eigenfunctions along boundary ρ, ϕ vanishing at $\rho = R$ and regular at $\rho = 0, 2\pi$ periodic in ϕ :

$$\psi_{mn}(\rho,\phi) = J_m(k_{mn}\rho) e^{im\phi} \qquad n = 1\dots\infty \quad m = -\infty\dots\infty$$

Here:

$$k_{mn} = \frac{x_{mn}}{R} \tag{D.11}$$

where x_{mn} is the *n*-th zero of the *m*-th Bessel function, *e.g.* the zeros of $J_0(x)$ are

 $(x_{01}, x_{02}, x_{03}) = 2.40483, 5.52008, 8.65373$ (D.12)

These are given by $x_{mn} = \text{BesselZeroJ}[m, n]$ in Mathematica. Note also that $J_{-m}(x) = J_m(x)$

(c) Orthogonality:

$$\int_{0}^{R} \rho d\rho \int_{0}^{2\pi} \psi_{mn}(\rho,\phi) \psi_{mn}(\rho,\phi) = \left(\frac{R^2}{2} \left[J_{m+1}(k_{mn}R)\right]^2\right) (2\pi) \,\delta_{nn'} \delta_{mm'}$$

(d) Solution:

$$\Phi = \sum_{n=1}^{\infty} \sum_{m=-\infty}^{\infty} \left[A_{mn} e^{-k_{mn}z} + B_{nm} e^{k_{mn}z} \right] \psi_{mn}(\rho,\phi)$$
(D.13)