

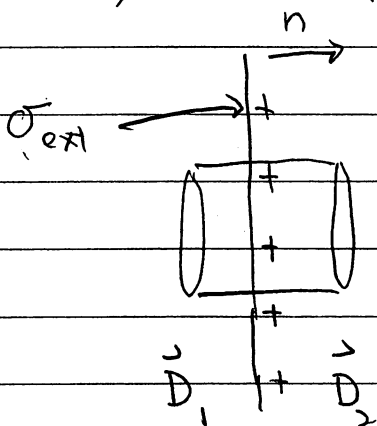
Boundary Conditions

Now we want to solve

$$\begin{aligned}\nabla \cdot \mathbf{D} &= \rho_{\text{ext}} \\ \nabla \times \mathbf{E} &= 0\end{aligned}$$

Before we can do so we need boundary conditions

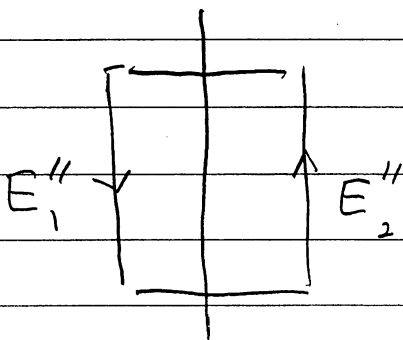
From $\nabla \cdot \mathbf{D} = \rho_{\text{ext}}$, we apply Gauss law to this and find



$$\vec{n} \cdot (\vec{D}_2 - \vec{D}_1) = \sigma_{\text{ext}}$$

i.e. the normal components of \mathbf{D} jump across the surface

From $\nabla \times \mathbf{E} = 0$ or $\oint \vec{E} \cdot d\vec{l} = (E_2'' - E_1'')l = 0$

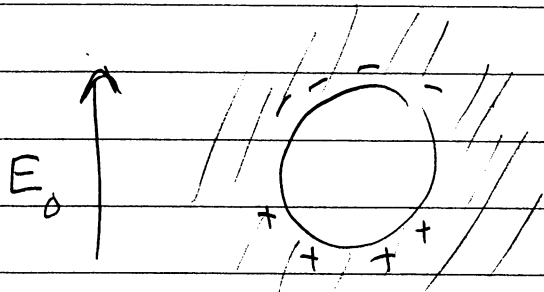


we have the continuity of parallel electric field component

$$E_2'' = E_1''$$

A model problem: A spherical cavity (Jackson 4.4)

Consider a cavity in a dielectric of radius \vec{a} in an external field.



The polarized charge we will find is sketched. (Why does it look like this?)

Solution

$$\nabla \cdot \mathbf{D} = \rho \quad \text{and for constant } \epsilon \quad \nabla \cdot (\epsilon \nabla \phi) = \rho$$

$$\nabla \times \mathbf{E} = 0 \quad \Rightarrow \quad \mathbf{E} = -\nabla \phi$$

So outside the sphere $\rho = 0$ and $-\nabla^2 \phi = 0$

$$\phi^{\text{out}} = \sum_{\ell} \left(A_{\ell} r^{\ell} + \frac{B_{\ell}}{r^{\ell+1}} \right) P_{\ell}(\cos \theta)$$

while inside

$$\phi^{\text{in}} = \sum_{\ell} \left(C_{\ell} r^{\ell} + \frac{D_{\ell}}{r^{\ell}} \right) P_{\ell}(\cos \theta)$$

The only non-vanishing A_{ℓ} , is A_1 .

$$\phi^{\text{out}} = -E_0 r \cos \theta + \sum_{\ell} \frac{B_{\ell}}{r^{\ell+1}} \cos \theta$$

Now on surface $r=a$ we have:

$$\vec{n} \cdot \vec{D}^{\text{out}} - n \cdot \vec{D}^{\text{in}} = 0 \Rightarrow \left. \begin{aligned} \varepsilon \frac{\partial \varphi^{\text{out}}}{\partial r} &= \frac{\partial \varphi^{\text{in}}}{\partial r} & (1) \\ E_2'' - E_1'' &= 0 & \Rightarrow \\ \frac{\partial \varphi^{\text{out}}}{\partial \theta} &= \frac{\partial \varphi^{\text{in}}}{\partial \theta} & (2) \end{aligned} \right\}$$

In general these equations are easily satisfied for $l \neq 1$. Just set $B_l = C_l = 0$. The $l=1$ equation gives a non-trivial condition

From (1) and (2) we find:

$$\left(-\varepsilon E_0 - 2\varepsilon \frac{B_1}{a^3} \right) = C_1$$

$$\left(-E_0 a + \frac{B_1}{a^2} \right) = C_1 a$$

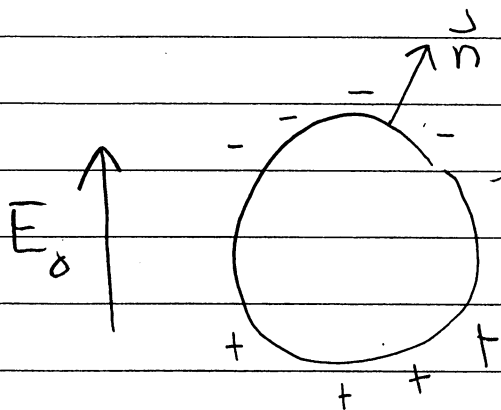
Solving for C_1 and B_1 we have: $(1-\varepsilon)$ is negative

$$\varphi = \begin{cases} -E_0 r \cos\theta + \frac{E_0 a^3 (1-\varepsilon)}{r^2 (1+2\varepsilon)} \cos\theta & r > a \\ -E_0 r \cos\theta \left(\frac{3\varepsilon}{1+2\varepsilon} \right) & r < a \end{cases}$$

We make the following remarks:

① The electric field is constant in the sphere

② The induced charge is:



$$\sigma_p = -\vec{n} \cdot (\vec{P}_2 - \vec{P}_1)$$
$$= -\chi \vec{n} \cdot \vec{E}$$

$$= -(\epsilon - 1) \left(-\frac{\partial \phi^{\text{out}}}{\partial r} \right) \Big|_{r=a}$$

This yields after differentiating ϕ^{out}

$$\sigma_p = -\frac{3(\epsilon - 1) E_0 \cos \theta}{1 + 2\epsilon}$$

thus we find the appropriate sign.