

## Dimensional Analysis of Maxwell Eqs

Examining the Maxwell Eqs (In heavyside-lorentz units)

$$\nabla \cdot E = \rho$$

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- Note every time derivative comes with  $1/c$ , i.e.  $\frac{1}{c} \frac{\partial}{\partial t}$

$$\nabla \times B = \frac{j}{c} + \frac{1}{c} \frac{\partial E}{\partial t}$$

↓

- Every velocity is measured in units of  $c$

$$\nabla \cdot B = 0$$

↓

$$-\nabla \times E = \frac{1}{c} \frac{\partial B}{\partial t}$$

$j = \frac{\text{charge}}{\text{volume}} \times \frac{\text{velocity}}{c}$

We see that if the system has characteristic length  $L$ , and characteristic time scale  $T$ , then the solutions will change rather dramatically from when

$$\frac{L}{c} \ll T \quad \text{vs.} \quad T \ll \frac{L}{c}$$

$\underbrace{\hspace{10em}}$  this is the regime of electrostatics, magnetostatics, and quasi-statics. The fields very rapidly adjust (with speed  $c$ ) to changes of charges.

$\underbrace{\hspace{10em}}$  This regime is radiation dominated. The system evolves significantly over the time it takes for the fields to propagate (at  $c$ ) across the system.

Consider the quasi-static regime:

$$\frac{L}{T} \ll c, \quad \frac{L}{cT} \sim 10^{-8} \text{ say.}$$

Then the different terms in the Maxwell equations have very different magnitudes.

For instance:

$$\nabla \cdot \bar{E} \sim \frac{\bar{E}}{L}$$

while

$$\frac{1}{c} \frac{\partial \bar{E}}{\partial t} \sim \frac{\bar{E}}{cT} \sim \frac{L}{cT} \left( \frac{\bar{E}}{L} \right) \ll \nabla \cdot \bar{E}$$

Thus we should set up a series solution in powers of  $\frac{1}{c}$ .

Most undergraduate courses stay entirely in this approximation scheme (without telling you).

Set up a series in  $1/c$

$$\mathbf{E} = \mathbf{E}^{(0)} + \mathbf{E}^{(1)} + \mathbf{E}^{(2)} + \dots$$

$$\mathbf{B} = \mathbf{B}^{(0)} + \mathbf{B}^{(1)} + \mathbf{B}^{(2)}$$

where  $\mathbf{E}^{(1)} \sim 10^{-8} \mathbf{E}^{(0)}$  and  $\mathbf{E}^{(2)} \sim 10^{-16} \mathbf{E}^{(0)}$  etc.

Then substituting this series into the Maxwell equations we find to zeroth order:

$$\begin{aligned}\nabla \cdot \mathbf{E}^{(0)} &= \rho & \uparrow & \text{This is electrostatics, } \mathbf{B}^{(0)} = 0 \\ \nabla \times \mathbf{B}^{(0)} &= 0 & \text{and} \\ \nabla \cdot \mathbf{B}^{(0)} &= 0 & & \nabla \cdot \mathbf{E} = \rho \\ -\nabla \times \mathbf{E}^{(0)} &= 0 & \downarrow & \nabla \times \mathbf{E} = 0\end{aligned}$$

At first order:

$$\begin{aligned}\nabla \cdot \mathbf{E}^{(1)} &= 0 & \text{determined} & \uparrow \text{This is magneto-} \\ \nabla \times \mathbf{B}^{(1)} &= \vec{j} + \frac{1}{c} \frac{\partial \mathbf{E}^{(0)}}{\partial t} & \downarrow \text{from electro} \\ \nabla \cdot \mathbf{B}^{(1)} &= 0 & \text{Statics} & \text{statics. } \mathbf{E}^{(1)} = 0 \\ \nabla \times \mathbf{E}^{(1)} &= \frac{1}{c} \frac{\partial \mathbf{B}^{(0)}}{\partial t} & & \text{and } \vec{B} \text{ solves:} \\ & & & \nabla \times \mathbf{B} = \vec{j}/c + \frac{1}{c} \frac{\partial \mathbf{E}^{(0)}}{\partial t} \\ & & & \nabla \cdot \mathbf{B} = 0\end{aligned}$$

One can continue this way and find corrections to electrostatics and magneto-statics. We will do this later in the course.