

Phase Space

- It is informative to study the motion of the system in the $x-p$ plane. Take the simple pendulum:



$$\theta \approx \pi$$

$$p_\theta = \frac{\partial L}{\partial \dot{\theta}} = ml^2 \dot{\theta}$$

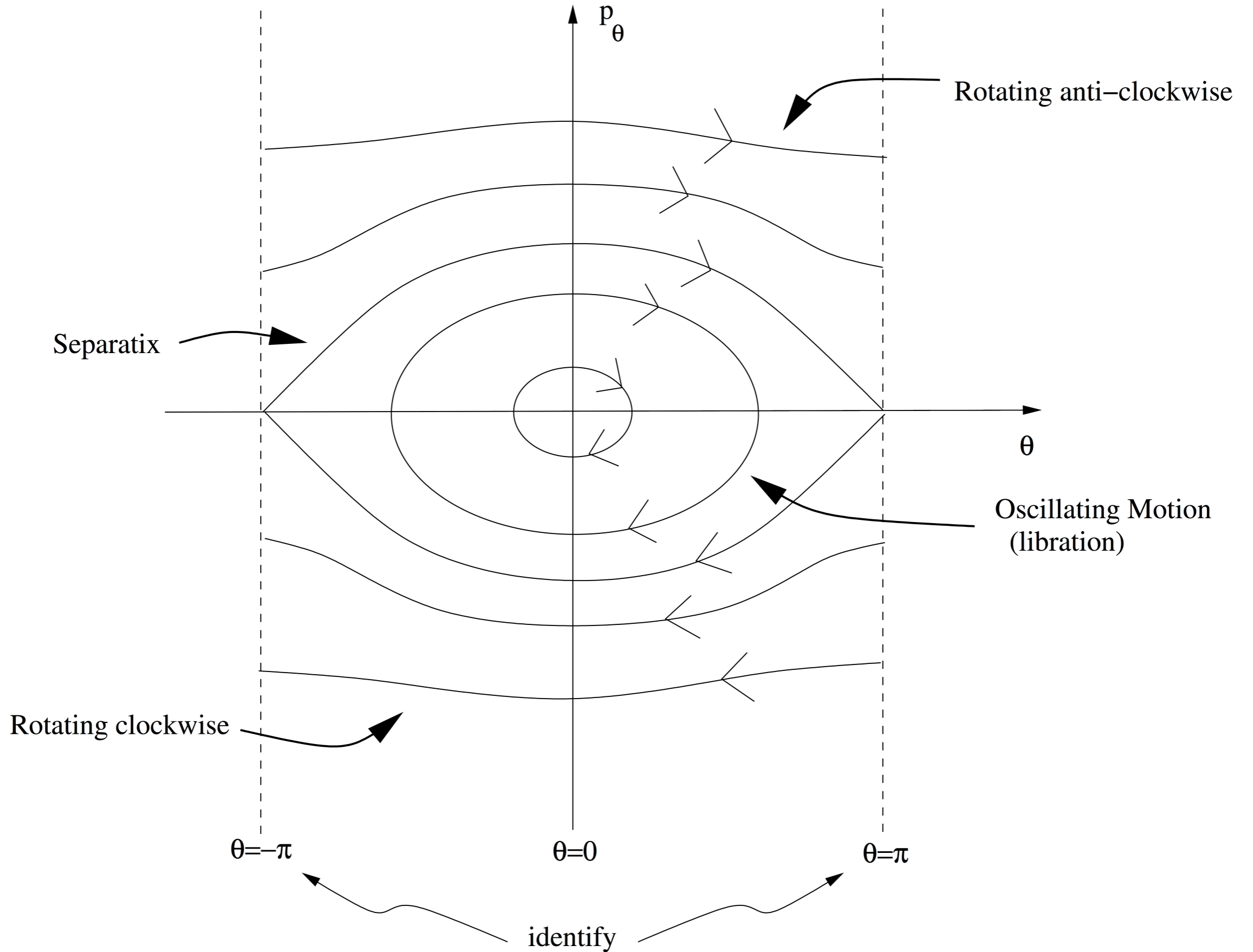
- For small oscillations, a trajectory's energy is:

$$E = \frac{p_\theta^2}{2ml^2} - mgl \cos \theta$$

$$\approx \frac{p_\theta^2}{2ml^2} + mgl \frac{\theta^2}{2} + \text{const}$$

- Thus in the phase space (θ, p_θ) the trajectory is a circle, after choosing appropriate units $m=g=l=1$. This is "Libration".
- For larger values of E , the circles are deformed, and ultimately changed completely when $E > E_{\text{separatrix}} = 2mgl$. Then the pendulum goes around in a circle, clockwise or counter-clockwise.
- We see how informative the phase space can be!

Phase portrait of a pendulum



Liouville Theorem

(next page)

- Now consider another phase-space plot. For an ensemble of particles undergoing constant acceleration ($x(t) = x_0 + v_0 t + 1/2 a t^2$)

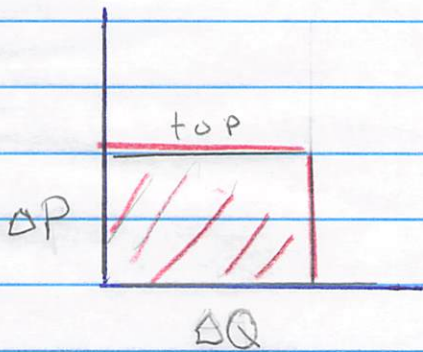
$$\left. \begin{aligned} Q(t) &= \frac{1}{2} a t^2 + q + \frac{p}{m} t \\ P(t) &= m a t + p \end{aligned} \right\} \begin{array}{l} \text{A map from:} \\ (q, p) \rightarrow (Q, P) \\ \uparrow \qquad \qquad \qquad \nwarrow \\ \text{initial state} \qquad \qquad \text{final state} \end{array}$$

- The next figure shows how an ensemble of particles starting in the square in the lower left evolves in time (see next page)

Liouville Theorem

- Quite generally the "phase-space" volume covered by the ensemble, does not grow in magnitude in time, but stretches and changes shape:

Simple Proof



$$q \rightarrow Q(t) = q + \dot{Q}(q, p) t$$

$$p \rightarrow P(t) = p + \dot{P}(q, p) t$$

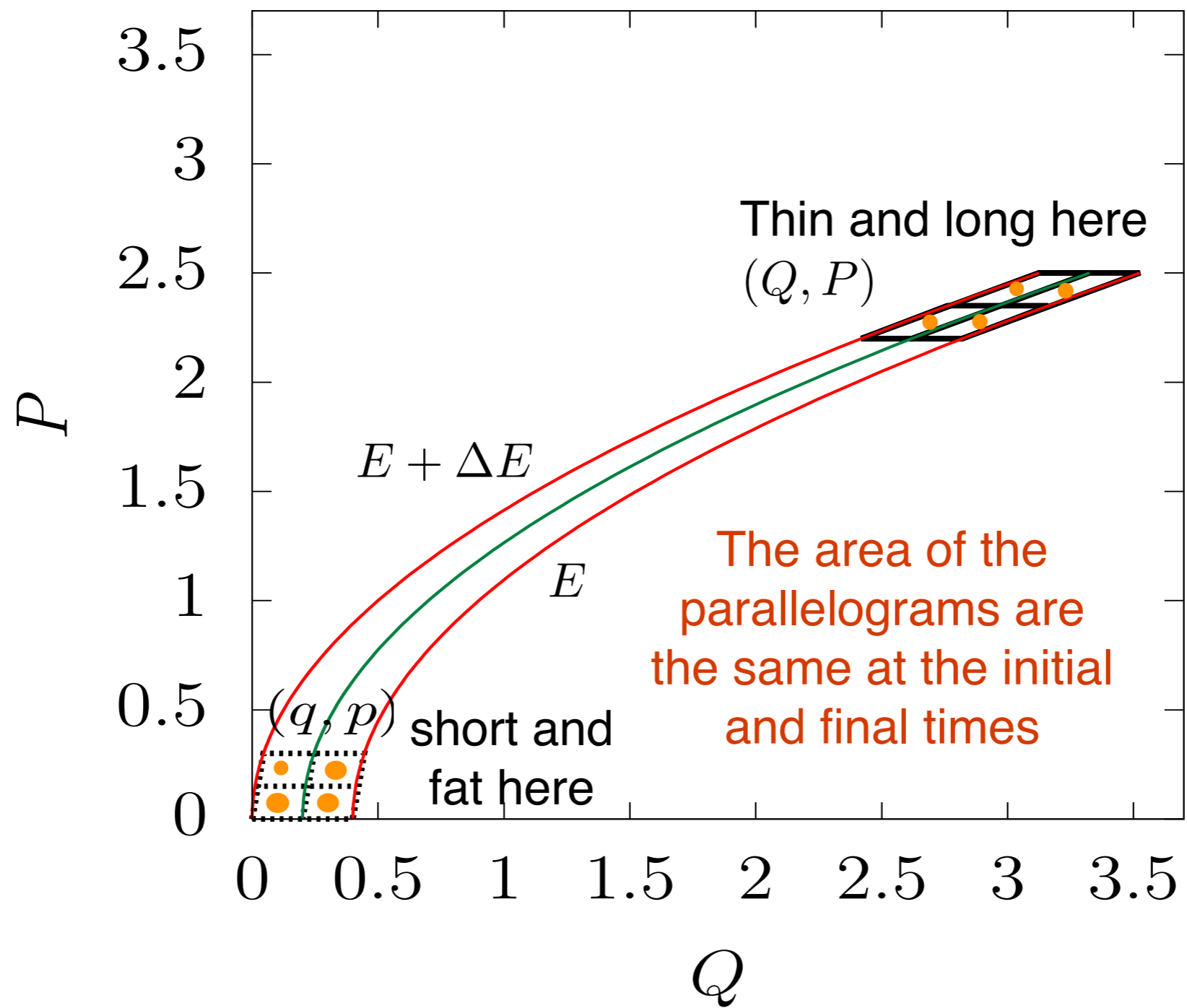
t is small

$$\dot{Q}(q, p) = \frac{\partial H(q, p)}{\partial p} \quad \dot{P} = \frac{\partial H(q, p)}{\partial q}$$

- How does the volume change in time?

$$d(\Delta P) = \left(\dot{P} \Big|_{\text{top}} - \dot{P} \Big|_{\text{bottom}} \right) dt = \frac{\partial \dot{P}}{\partial p} \Delta p dt$$

Phase space volume conserved



Similarly

$$d(\Delta Q) = \left(\dot{Q} \Big|_{\text{right}} - \dot{Q} \Big|_{\text{left}} \right) dt$$
$$= \frac{\partial \dot{Q}}{\partial q} \Delta q dt$$

• So

$$d(\Delta P \Delta Q) = \Delta P d(\Delta Q) + d(\Delta P) \Delta Q$$
$$= \left(\frac{\partial \dot{Q}}{\partial q} + \frac{\partial \dot{P}}{\partial p} \right) dt \Delta P \Delta q$$

So

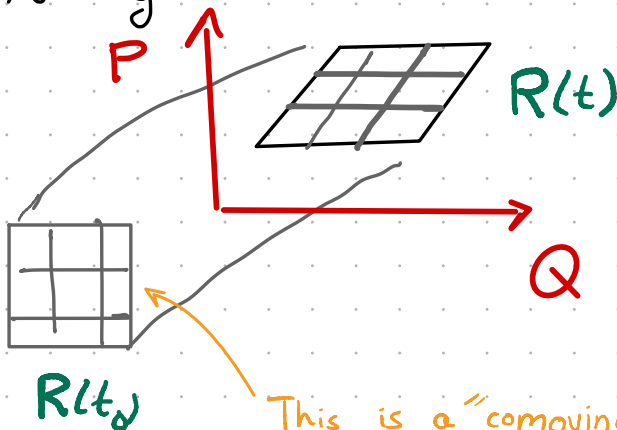
$$d(\Delta P \Delta Q) = \left(\frac{\partial^2 H}{\partial q \partial p} - \frac{\partial^2 H}{\partial p \partial q} \right) dt \Delta P \Delta q = 0$$



Change in phase-space volume of square

Liouville Theorem: the comoving phase-space volume is constant in time

- A more formal procedure/proof is the following:



This is a "comoving" volume meaning the boundary follows the EOM

- Then the area at time t is

$$A(t) = \int_{R(t)} dQdP = \int_{R(t_0)} \left| \frac{\partial(Q,P)}{\partial(q,p)} \right| dqdp$$

we will show now that the determinant is unity. Thus the areas at the final and initial times are equal

$$A(t) = \int_{R(t_0)} dqdp = A(t_0)$$

Proof:

$$Q(t) = q + \frac{\partial H(q,p)}{\partial p} t$$

$$P(t) = p - \frac{\partial H(q,p)}{\partial q} t$$

$$\begin{aligned}
 \text{So } \left| \frac{\partial(Q, P)}{\partial(q, p)} \right| &= \begin{vmatrix} \partial Q / \partial q & \partial Q / \partial p \\ \partial P / \partial q & \partial P / \partial p \end{vmatrix} \\
 &= \begin{vmatrix} 1 + t \frac{\partial^2 H}{\partial q \partial p} & t \frac{\partial^2 H}{\partial p^2} \\ t \frac{\partial^2 H}{\partial p^2} & 1 - t \frac{\partial^2 H}{\partial p \partial q} \end{vmatrix} \\
 &= 1 + t \left(\frac{\partial^2 H}{\partial q \partial p} - \frac{\partial^2 H}{\partial p \partial q} \right) + \underbrace{O(t^2)}_{\text{neglect}} \\
 &= 1 + O(t^2)
 \end{aligned}$$

So the volume is preserved by the map.

More Coordinates

In general

$$\begin{aligned}
 A(t) &= \int_{R(t)} dQ^1 dP_1 dQ^2 dP_2 \dots \\
 &= \int_{R(t_0)} \left| \frac{\partial(Q, P)}{\partial(q, p)} \right| dq^1 dp_1 dq^2 dp_2 \dots \\
 &= A(t_0), \quad \text{provided this determinant is unity (it is!)}
 \end{aligned}$$

Proof:

$$q^i \rightarrow Q^i = q^i + \frac{\partial H}{\partial p_i} t$$

$$p_i \rightarrow P_i = p_i - \frac{\partial H}{\partial q_i} t$$

The Jacobian of the map is

$$M = \det \begin{pmatrix} \frac{\partial Q^i}{\partial q^j} & \frac{\partial Q^i}{\partial P_j} \\ \frac{\partial P_i}{\partial q^j} & \frac{\partial P_i}{\partial P_j} \end{pmatrix}$$

$$= \det \begin{pmatrix} \delta_j^i + \frac{\partial^2 H}{\partial p_i \partial q^j} t & \frac{\partial^2 H}{\partial p_i \partial p_j} t \\ t \frac{\partial^2 H}{\partial q^i \partial q^j} & \delta_j^i - t \frac{\partial^2 H}{\partial p_j \partial q^i} \end{pmatrix}$$

• Then note for a matrix M

$$\det M = \exp \text{Tr} \log M$$

↖ just prove me by working in the eigenbasis of M

- Now in this case the matrix M is close to the identity

$$M = \mathbb{I} + \overset{\text{small time}}{t} M^{(1)}$$

$$\log M \simeq t M^{(1)}, \quad \text{Tr} \log M \simeq t \text{Tr} M^{(1)}$$

$$\det M \simeq e^{t \text{Tr} M^{(1)}} \simeq 1 + t \text{Tr} M^{(1)}$$

- But in this case

$$\text{Tr} M^{(1)} = \frac{\partial^2 H}{\partial p_i \partial q^i} - \frac{\partial^2 H}{\partial p_i \partial q^i} = 0$$

So the volume in phase space of any comoving region is constant in time!

(since the det of the map is unity)