

Answer Note that the result for the moment of inertia of a cylinder does not depend on L , the length of the cylinder. In other words, it applies equally well to a long cylinder and


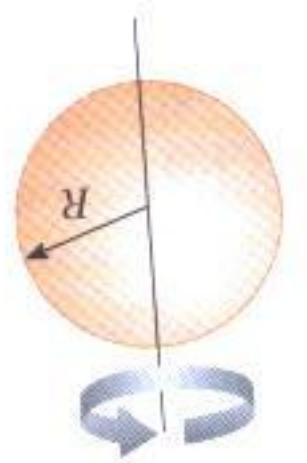


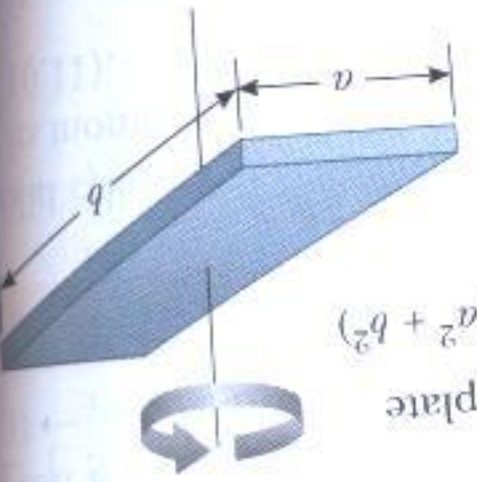
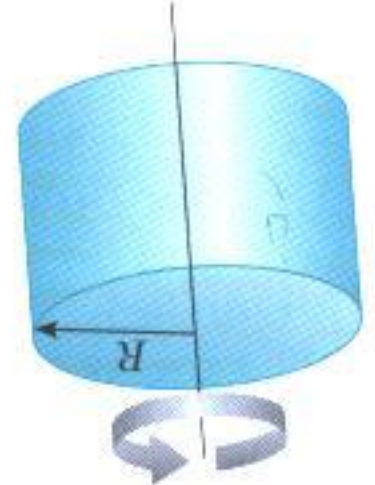
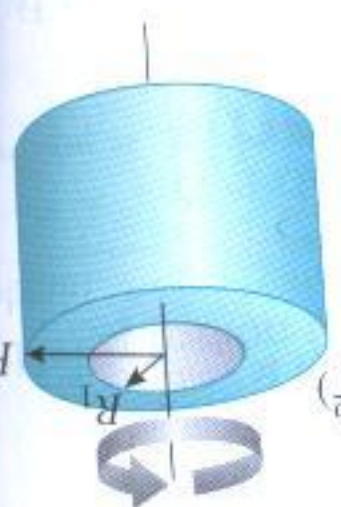
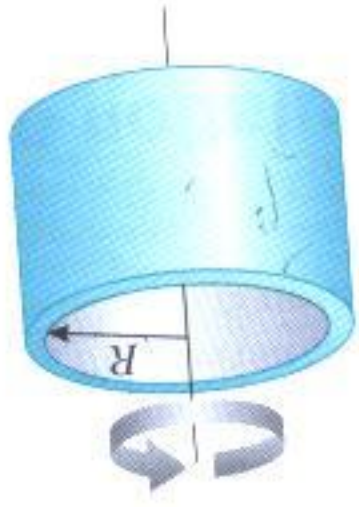
Parallel-axis theorem

Table 10.2 gives the moments of inertia for a number of objects about specific axes. The moments of inertia of rigid objects with simple geometry (high symmetry) are relatively easy to calculate provided the rotation axis coincides with an axis of symmetry. The calculation of moments of inertia about an arbitrary axis can be cumbersome, however, even for a highly symmetric object. Fortunately, use of an important theorem called the **parallel-axis theorem**, often simplifies the calculation. Suppose the moment of inertia about an axis through the center of mass of an object is I_{CM} . The parallel-axis theorem states that the moment of inertia about any axis parallel to the parallel-axis theorem states that the moment of inertia about any axis parallel to a distance D away from this axis is

$$I = I_{CM} + MD^2$$

To prove the parallel-axis theorem, suppose that an object rotates in the xy plane about the z axis, as shown in Figure 10.12, and that the coordinates of the center of mass are x_{CM} , y_{CM} . Let the mass element dm have coordinates x , y . Because

Table 10.2
Moments of Inertia of Homogeneous Rigid Objects with Different Geometries

	Thin spherical shell $I_{CM} = \frac{2}{3} MR^2$		Solid sphere $I_{CM} = \frac{2}{5} MR^2$
	Long thin rod with rotation axis through end $I = \frac{1}{3} ML^2$		Long thin rod with rotation axis through center $I_{CM} = \frac{1}{12} ML^2$
	Rectangular plate $I_{CM} = \frac{1}{12} M(a^2 + b^2)$		Solid cylinder or disk $I_{CM} = \frac{1}{2} MR^2$
	Hollow cylinder $I_{CM} = \frac{1}{2} M(R_1^2 + R_2^2)$		Hoop or thin cylindrical shell $I_{CM} = MR^2$