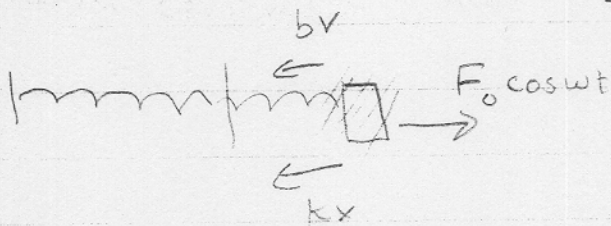


Last Time



- Drive the spring with an external force $F_0 \cos \omega t$
- The motion reaches a steady state
- The amplitude can become large if tuned @ ω_0
→ Remember the bridge



Newton

$$m a = -kx - bv + F_0 \cos \omega t$$

$$\begin{matrix} \textcircled{1} & \textcircled{2} & \textcircled{3} \\ \frac{d^2 x}{dt^2} + \frac{b}{m} \frac{dx}{dt} + \frac{k}{m} x = \frac{F_0}{m} \cos \omega t \end{matrix}$$

Solution:

$$x = A_{\omega} \cos(\omega t - \phi_{\omega})$$

Ansatz

$$A = \frac{F_0/m}{\left[(\omega^2 - \omega_0^2)^2 + \left(\frac{b}{m} \omega \right)^2 \right]^{1/2}}$$

$$\phi = \tan^{-1} \left[\frac{b/m \omega}{(\omega_0^2 - \omega^2)} \right]$$

Derivation - Skipped in class

$$\ddot{x} = -A \omega^2 \cos(\omega t - \phi)$$

$$\dot{x} = -A \omega (\sin \omega t - \phi)$$

$$x = A \cos(\omega t - \phi)$$

①

②

k/m

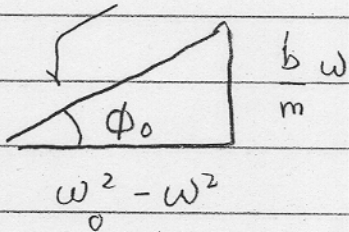
///

③

$$-A \omega^2 \cos(\omega t - \phi) - \frac{b}{m} \omega A \sin \omega t - \phi + \omega_0^2 A \cos(\omega t - \phi) = \frac{F_0}{m} \cos \omega t$$

$$A \left[(\omega_0^2 - \omega^2) \cos(\omega t - \phi) - \frac{b}{m} \omega \sin(\omega t - \phi) \right] = \frac{F_0}{m} \cos \omega t$$

Same trick



$$\sqrt{\quad} = \left((\omega_0^2 - \omega^2)^2 + \left(\frac{b}{m} \omega \right)^2 \right)^{\frac{1}{2}}$$

$$A \sqrt{\quad} \left[\frac{\omega_0^2 - \omega^2}{\sqrt{\quad}} \cos(\omega t - \phi) - \frac{(b/m)\omega}{\sqrt{\quad}} \sin(\omega t - \phi) \right] = \frac{F_0}{m} \cos \omega t$$

$$A \sqrt{\quad} \left[\cos \phi_0 \cos(\omega t - \phi) - \sin \phi_0 \sin(\omega t - \phi) \right] = \frac{F_0}{m} \cos \omega t$$

$$A \sqrt{\quad} \left[\cos(\omega t - \phi + \phi_0) \right] = \frac{F_0}{m} \cos \omega t$$

Need to choose (to have cos match)

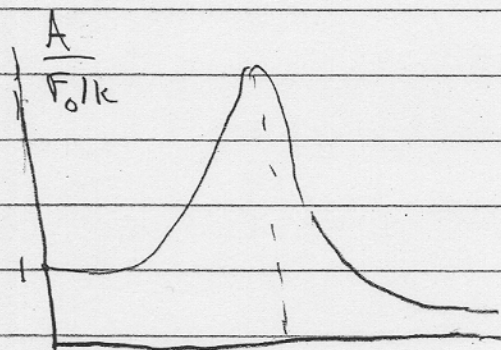
$$\phi = \phi_0 = \tan^{-1} \left(\frac{b/m\omega}{\omega_0^2 - \omega^2} \right) = \phi$$

Then

$$A\Gamma = \frac{F_0}{m}$$

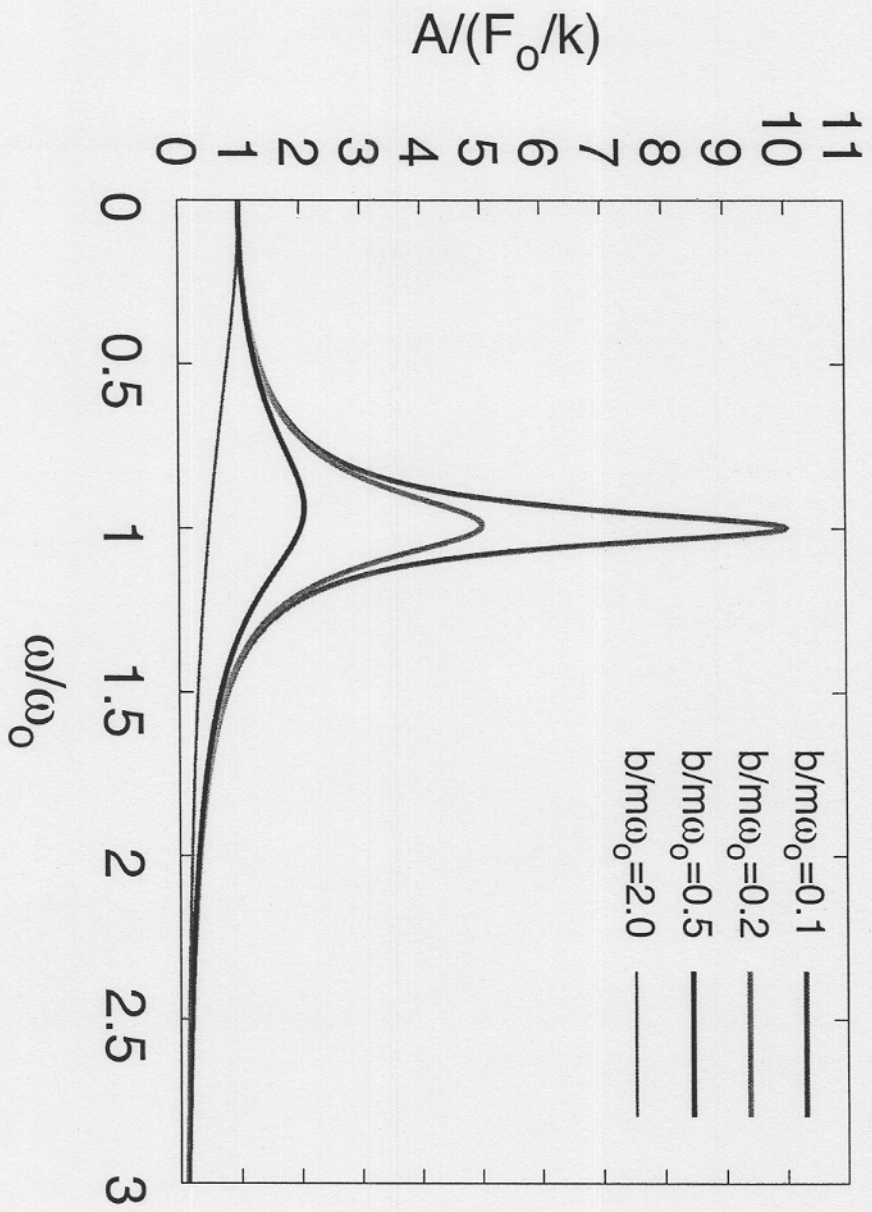
$$\text{Or } A = \frac{F_0/m}{\left[(\omega_0^2 - \omega^2)^2 + \left(\frac{b}{m} \omega \right)^2 \right]^{1/2}}$$

Usually we are interested in $\frac{b}{m\omega_0} \ll 1$ (Continuous)



$$\omega \approx \omega_0$$

- Many interesting features
- See slides



At Resonance $\omega \approx \omega_0$

$$A = \frac{F_0/m}{\max \left(\omega^2 + \left(\frac{b}{m}\omega_0 \right)^2 \right)^{1/2}}$$

$$A = \frac{F_0/m}{\max \left(\frac{b}{m}\omega_0 \right)} = \frac{F_0}{m\omega_0^2} \cdot \frac{1}{\left(\frac{b}{m\omega_0} \right)} = \boxed{\frac{F_0}{k} \left(\frac{m\omega_0}{b} \right)}$$

Usually we define

$$Q \equiv \frac{m\omega_0}{b} \sim \text{ " } \underbrace{\hspace{10em}} \text{ "}$$

damping in dimensionless units

So

$$A_{\max} = \frac{F_0}{k} \cdot Q$$

For a violin $Q \sim 30-50$

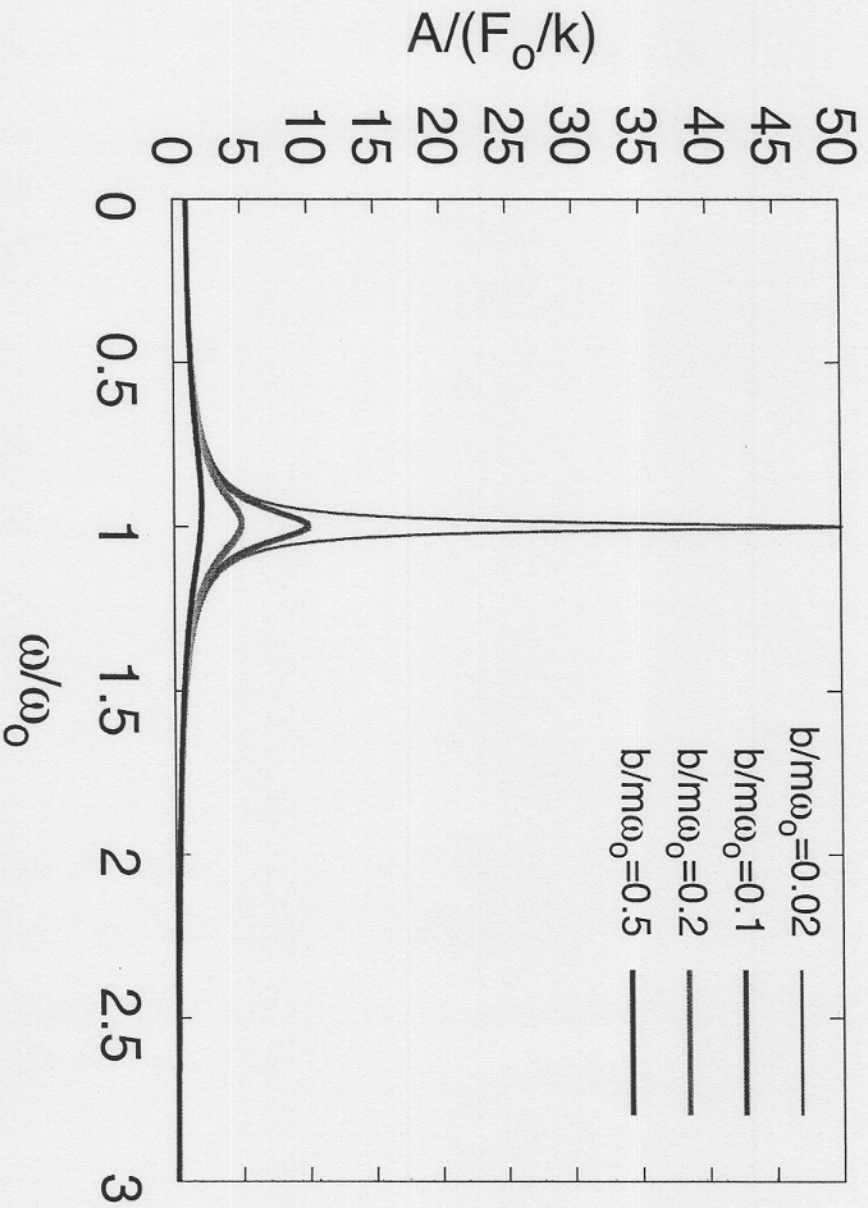
Notice:

\rightarrow As the Q becomes large the "spike" becomes increasingly sharp

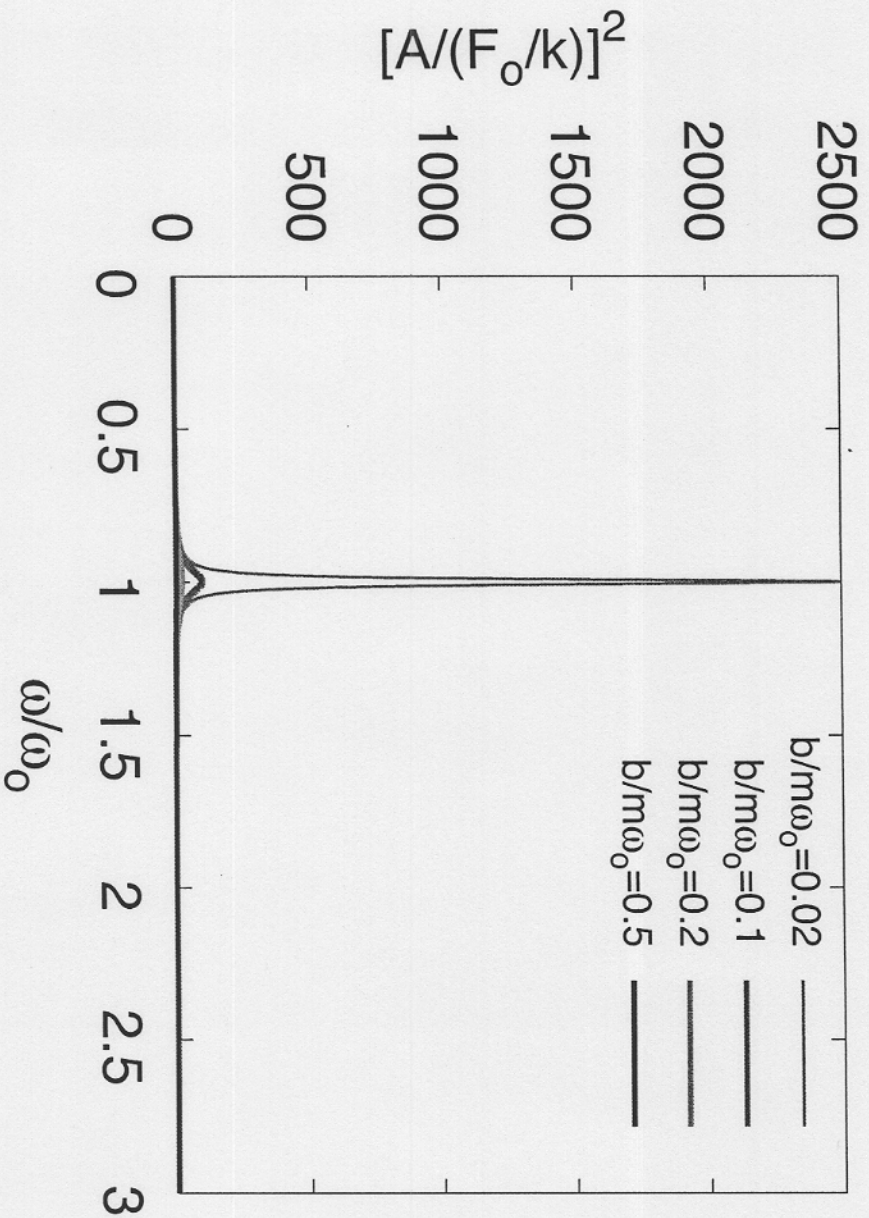
\swarrow see figure

$$\text{Small} \rightarrow \frac{1}{Q} \approx \frac{\Delta\omega}{\omega} \equiv \text{Full width at half max in } A^2_{\text{plot}}$$

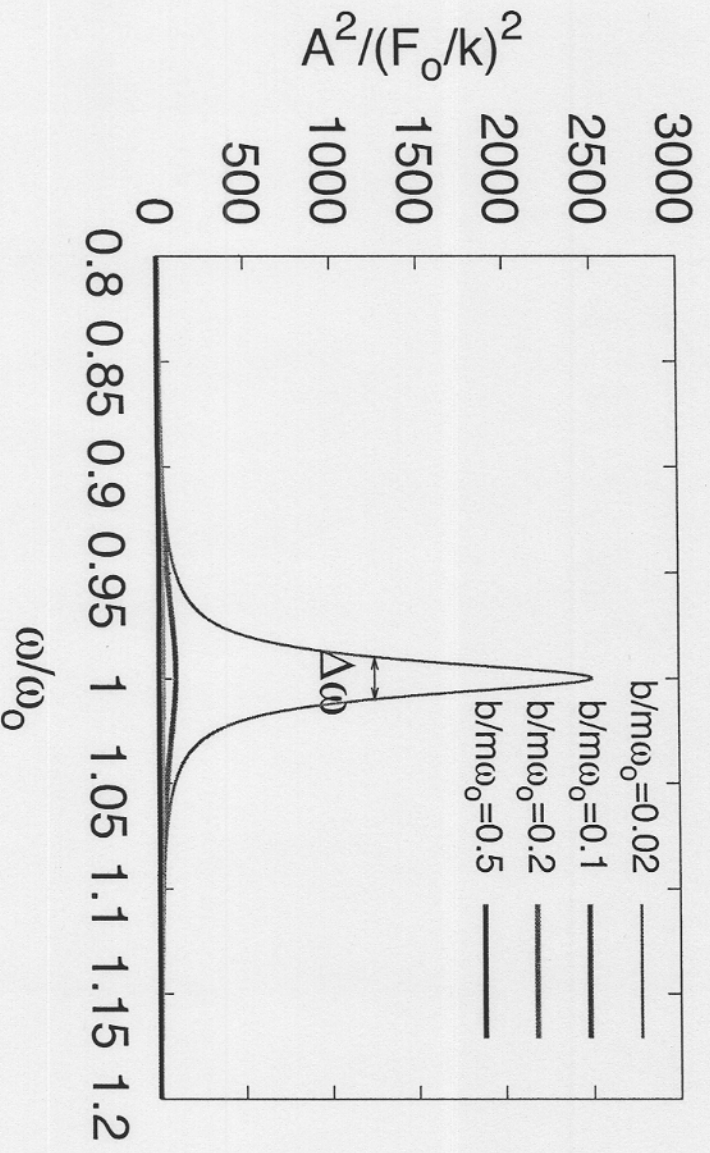
Realistic $Q \approx 50$ for a violin



Realistic $Q \approx 50$ for a violin, we also care about energy $\propto A^2$



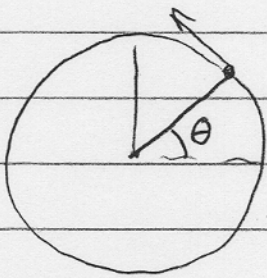
A^2 proportional to energy in the oscillator



$$\frac{1}{Q} \approx \frac{\Delta\omega}{\omega} = \frac{\text{Full width at half maximum in } A^2 \text{ plot}}{\omega}$$

Rotational Motion

Angular Quantities -- Always use radians internally



$$\theta = \frac{l}{R}$$

$$\omega = \frac{d\theta}{dt} = \frac{dl/dt}{R} = \frac{v}{R}$$

$$\boxed{\omega = \frac{v}{R} \quad \text{or} \quad v = \omega R}$$

Units of ω

• rad/s

or

• rev/s, 1 rev = 2π rad, rpm = rotations per minute

Also use the frequency

$$f = \frac{\omega}{2\pi} = \frac{\# \text{ of complete cycles}}{\text{sec}} = \text{Hz}$$

If the motion is uniform define

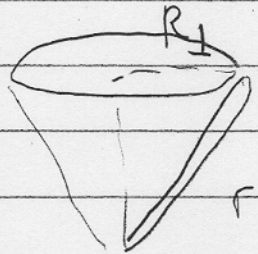
$$T = \frac{1}{f}$$

What else needs to be said?

- ① Implicit in ω is an agreed upon center point
- a fixed point
 - Or the CM if the system is moving

- ② ω is independent of where you are on the wheel, a property of body & center point

③

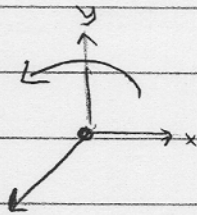


$$\omega = \frac{d\theta}{dt} = \frac{v}{R_{\perp}}$$

$$v = \omega R_{\perp} = \omega r \sin\theta$$

Angular Velocity is actually a vector

- ① We choose the angular velocity $\vec{\omega}$ to be the direction of rotating axis according to the "right hand rule"



- ② So a counter clockwise rotation in x-y plane we would say is

$$\vec{\omega} = \omega \hat{k} \leftarrow \text{out of the page}$$

↖ magnitude of ω , pos number

While a clockwise rotation is into the page

$$\vec{\omega} = -\omega \hat{k}$$

↖ pos number

- ③ If we only ever talk about rotation in the x-y plane we can use the sign to indicate the direction of rotation

