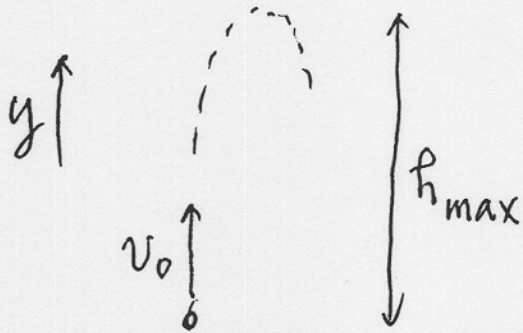


Lecture

d-1

throwing ball upward: with friction



if no air friction

$$h_{\max} = \frac{v_0^2}{2g}$$

suppose $f = \text{friction force} = \alpha v^2$
 \uparrow cst

no friction
 $v_y = v_0 - gt \Rightarrow dt = \frac{v_0}{g} - \frac{v}{g}$
 $h = v_0 t - \frac{gt^2}{2} = \frac{v_0^2}{2g}$

find $h_{\max} = ?$

Newton:
 laws of physics are diff. equs!

$$F = ma$$

$$-mg - \alpha v^2 = ma = m \frac{dv}{dt}$$

$$k \equiv \frac{\alpha}{m}$$

$$-g - kv^2 = \frac{dv}{dt} \quad \dots \quad (1)$$

$$v = \frac{dy}{dt} = \frac{dy}{dv} \frac{dv}{dt}$$

? how to solve ?
 either $dt = \frac{-dv}{g + kv^2} \rightarrow t = -\int_{v_0}^v \frac{dv'}{g + kv'^2}$
 but $v = \frac{dy}{dt}$; so $dt = \frac{dy}{v}$

$$\rightarrow dy = \frac{v dv}{\frac{dv}{dt}} \quad \text{using (1)} \rightarrow dy = \frac{v dv}{-g - kv^2} \quad (2)$$

(2) is easy:

$$\int_0^{h_{\max}} dy = \int_{v_0}^0 \frac{v dv}{-g - kv^2} = -\frac{1}{2k} \int_{v_0}^0 \frac{d(kv^2)}{g + kv^2}$$

$$= -\frac{1}{2k} \ln(g + kv^2) \Big|_{v_0}^0$$

$$\therefore h_{\max} = -\frac{1}{2k} \left[\ln g - \ln (g + kV_0^2) \right]$$

$$= -\frac{1}{2k} \ln \frac{g + kV_0^2}{g} \quad (4)$$

naively $\ln \frac{g}{g} = 0$
 but $\frac{1}{k}$ so looks $\frac{0}{0}$

when no fric. force $k \rightarrow 0$

above should reduce to $\frac{V_0^2}{2g}$?

$$(4) \rightarrow h_{\max} = -\frac{1}{2k} \ln \left(1 + \frac{kV_0^2}{g} \right)$$

(using $\ln(1-x) = -(x + \frac{x^2}{2} + \frac{x^3}{3} + \dots)$)

$$\xrightarrow{k \rightarrow 0} -\frac{1}{2k} \left(-\frac{kV_0^2}{g} \right) = \frac{V_0^2}{2g} \quad | \quad \text{no, } x = -\frac{kV}{g}$$

let me keep the next term

$$\frac{1}{2k} \left(\frac{kV_0^2}{g} + \frac{1}{2} \left(\frac{kV_0^2}{g} \right)^2 \right) = + \frac{V_0^2}{g} + \frac{kV_0^4}{4g^2} + \dots$$

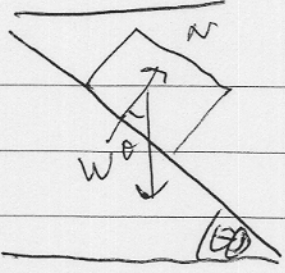
now we get correct

Lecture

~~Friction~~

Frictionless

Motion With Inclined Planes



$$\begin{cases} -W + N \cos \theta = m a_y \\ N \sin \theta = m a_x \end{cases}$$

3 unknown
+ const.

constraint

$$y = -x \tan \theta ; \quad \ddot{y} = -\ddot{x} \tan \theta$$

$$a_y = -a_x \tan \theta$$

$$a_x = N \frac{\sin \theta}{m} ; \quad a_y = -N \frac{\sin \theta}{m} \frac{\sin \theta}{\cos \theta} = -N \frac{\sin^2 \theta}{\cos \theta}$$

$$-mg + N \cos \theta = -N \frac{\sin^2 \theta}{\cos \theta}$$

~~N cos theta~~

$$N \left(\frac{\cos^2 \theta}{\cos \theta} + \frac{\sin^2 \theta}{\cos \theta} \right) = mg$$

$$N = mg \cos \theta$$

less than weight
zero if $\theta = \pi/2$

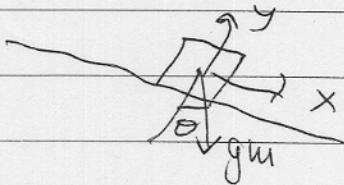
$$a_x = g \sin \theta \cos \theta$$

zero if $\theta = 0$ or 90°

$$a_y = -g \sin^2 \theta$$

at $\theta = 0$ zero
at $\theta = 90^\circ$ it is $-g$
yes

The "turning head" solution



$$gm \sin \theta = m a_x$$

$$gm \cos \theta - N = m a_y$$

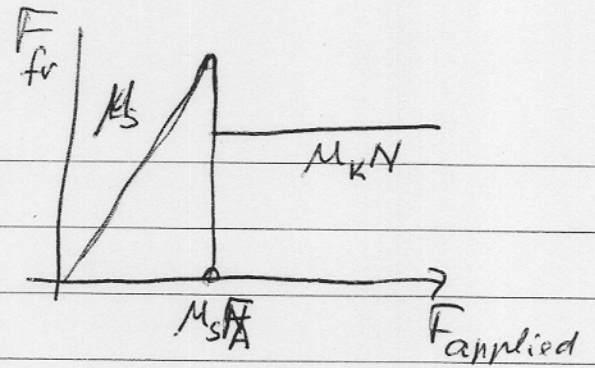
now $y = \text{const}$
 $\ddot{y} = 0$

thus they decouple

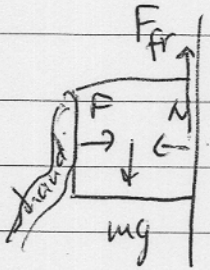
$$N = gm \cos \theta$$

$$a_x = g \sin \theta$$

Now with friction



A box against a wall



$$\mu_s F = F_{Fr} > mg$$

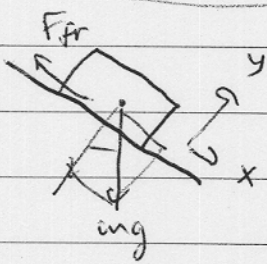
condition of no slipping

if it is less, acceleration

$$\mu_k F - mg = ma$$

$$y = \frac{at^2}{2} = \left(-g + \frac{\mu_k F}{m}\right) \frac{t^2}{2} < 0$$

↑ pushing makes falling slower



$$N = mg \cos \theta$$

$$-\mu_k N + mg \sin \theta = ma_x$$

$$-\mu_k \cos \theta + \sin \theta = \frac{a_x}{g}$$

should be > 0 !

critical angle $\mu_k = \tan \theta$

if friction is larger, no sliding

more accurately,

$$\tan \theta = \mu_s$$