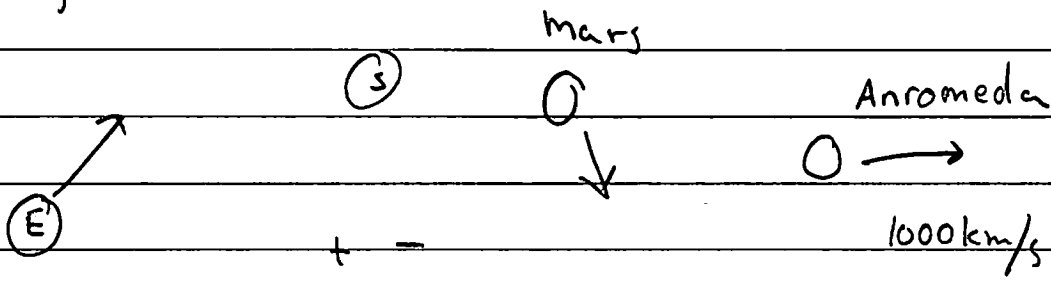


Relativity



- The speed of light is constant for all inertial observers. Nothing travels faster than c .
- The laws of physics are the same for all observers.

Earth

$$\frac{d\mathbf{p}}{dt} = q\vec{E} + q\frac{\vec{v}}{c} \times \vec{B}$$

$$\nabla \cdot \vec{E} = \rho$$

$$-\frac{1}{c} \frac{\partial \vec{E}}{\partial t} + \nabla \times \vec{B} = \mathbf{j}$$

$$\nabla \cdot \vec{B} = 0$$

$$\frac{1}{c} \frac{\partial \vec{B}}{\partial t} + \nabla \times \vec{E} = 0$$

the only thing relativity changes is the relation between \vec{v} and \mathbf{j}

$$\vec{v} = \frac{\mathbf{j}}{E/c^2}$$

Andromeda measures the same equations:

eg.

$$\frac{d\mathbf{p}}{dt} = q\vec{E} + q\frac{\vec{v}}{c} \times \vec{B}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot \mathbf{j} = 0$$

Relativity relates $(t, \vec{x}, \vec{E}, \vec{B}, \rho, \mathbf{j})$ to $(\underline{t}, \underline{x}, \underline{E}, \underline{B}, \underline{\rho}, \underline{\mathbf{j}})$

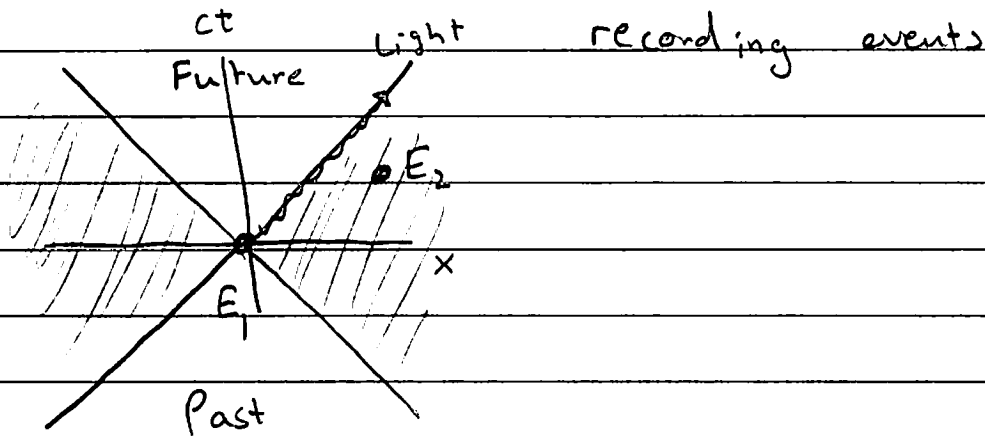
- Our goal is to introduce a notation to make this manifest.

Transformation of Coordinates

We describe physics as a sequence of events labelled by space-time coords $x^m = (x^0, x^1, x^2, x^3)$

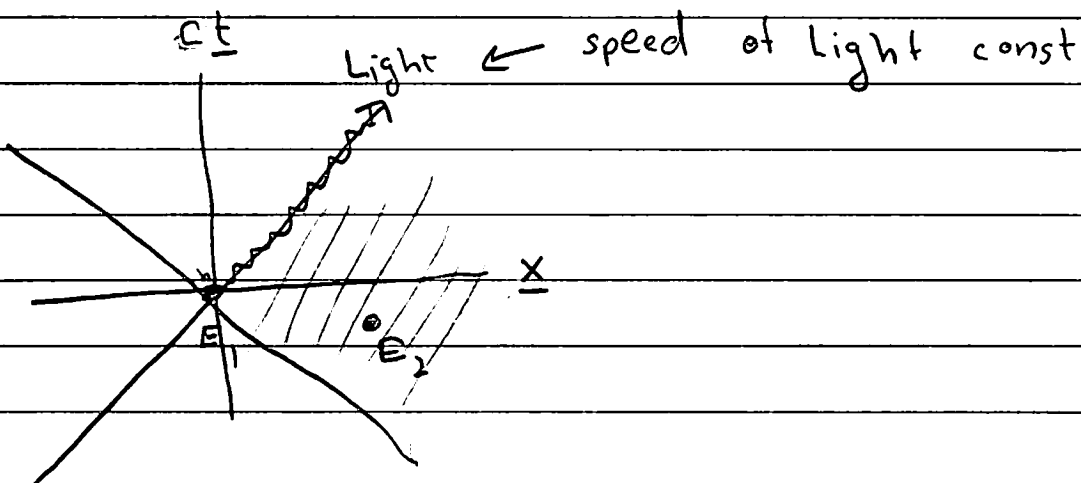
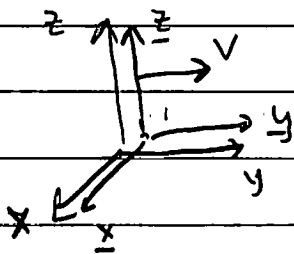
$$\equiv (ct, \vec{x})$$

Each observer set up his own coordinate system



An observer moving with speed v relative to the first makes measurement of the

same events. The order of events is sacred only if causally connected



We look for a linear change of coordinates which leaves the trajectory of light at fixed speed, $c = x/t$

$$\star \quad -c^2 t^2 + \underline{x}^2 = -c^2 \underline{t}^2 + \underline{\underline{x}}^2$$

So $x \rightarrow \underline{x}^m = L^m_{\nu}(v) x^\nu$

$$\begin{pmatrix} \underline{x}^0 \\ \underline{x}^1 \\ \underline{x}^2 \\ \underline{x}^3 \end{pmatrix} = \begin{pmatrix} L(v) \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix} =$$

how space time coords of events in one frame K transform, to another frame \underline{K} moving @ velocity v

Properties

$$L(-\underline{v}) L(\underline{v}) = 1$$

$$L(-v) = L^{-1}(v)$$

$\star\star$

$$L(v_2) L(v_1) = L(v_3)$$

This is a group of transformations, The Lorentz Group. With \star and $\star\star$ find

$$\begin{pmatrix} \underline{x}^0 \\ \underline{x}^1 \\ \underline{x}^2 \\ \underline{x}^3 \end{pmatrix} = \begin{pmatrix} \gamma & -\gamma\beta & & \\ -\gamma\beta & \gamma & & \\ & & 1 & \\ & & & 1 \end{pmatrix} \begin{pmatrix} x^0 \\ x^1 \\ x^2 \\ x^3 \end{pmatrix}$$

$$\gamma \equiv \frac{1}{\sqrt{1-\beta^2}} \quad \beta \equiv \frac{v}{c}$$

Parametrize the frame velocity $\frac{V}{c}$ with rapidity

$$\frac{V}{c} = \tanh \mathcal{J}$$

So

$$\left. \begin{aligned} \gamma &= \cosh \mathcal{J} \\ \gamma \beta &= \sinh \mathcal{J} \end{aligned} \right\}$$

$$\underline{x}^0 = x^0 \cosh \mathcal{J} - x^1 \sinh \mathcal{J}$$

$$x^1 = -x^1 \sinh \mathcal{J} - x^0 \cosh \mathcal{J}$$

Four vectors

• We denote $x_\mu = (x_0, x_1, x_2, x_3) = (-ct, \vec{x})$

Then

$$\underline{x}_\mu \underline{x}^\mu = x_0 x^0 + x_i x^i = -(ct)^2 + \vec{x} \cdot \vec{x}$$

is invariant. Since under a change of frame

$$\underline{x}^\mu = L^\mu_\nu x^\nu$$

We find

$$\underline{x}_\mu = x_\nu (L^{-1})^\nu_\mu \quad \text{or}$$

$$\begin{pmatrix} x_0 \\ x_1 \\ x_2 \\ x_3 \end{pmatrix} = \begin{pmatrix} x_0 & x_1 & x_2 & x_3 \end{pmatrix} \begin{pmatrix} L^{-1} \end{pmatrix}$$

Then

$$\begin{aligned} \underline{x}_\mu \underline{x}^\mu &= \underbrace{(x_0 \dots x_3)}_{\text{Lower components}} \begin{pmatrix} x^0 \\ \vdots \\ x^3 \end{pmatrix} = \underbrace{(x_0 \dots)}_{\text{upper components}} (L^{-1}) (L) \begin{pmatrix} x^0 \\ \vdots \end{pmatrix} \\ &= (x_0 \dots) \begin{pmatrix} x^0 \\ \vdots \end{pmatrix} \\ &= \underline{x}_\mu \underline{x}^\mu \end{aligned}$$

In general define a four vector $A^\mu = (a^0, \vec{a})$ as a quantity that transforms like

$$\underline{A}^\mu = L^\mu{}_\nu A^\nu \quad (\text{contravariant components})$$

The covariant components $A_\mu = (a_0, \vec{a}) = (-a^0, \vec{a})$

$$\underline{A}_\mu = A_\nu (L^{-1})^\mu{}_\nu \quad \leftarrow \text{transformation rule for lower indices}$$

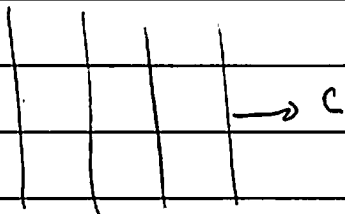
Then

$$A_\mu A^\mu = -(a^0)^2 + \vec{a}^2 \quad \text{is invariant}$$

$$\text{Similarly } A_\mu B^\mu = a_0 b^0 + \vec{a} \cdot \vec{b}$$

$$= -a^0 b^0 + \vec{a} \cdot \vec{b} \quad \text{is invariant}$$

Example - Relativistic Doppler Shift



$$A e^{-i\omega t + i\vec{k}\cdot\vec{x}}$$

a wave with $\omega = ck$

Since the speed of light is constant

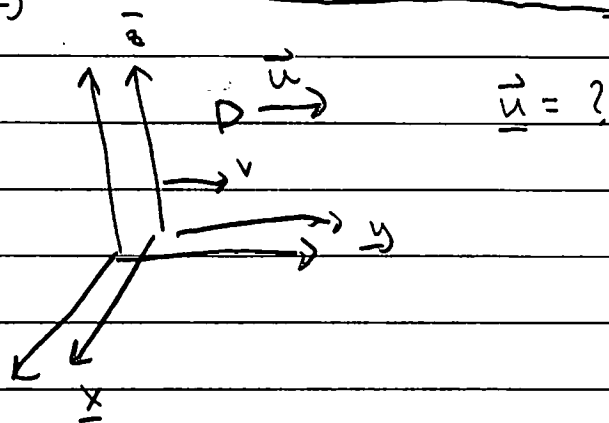
$$\phi = -\omega t + \vec{k}\cdot\vec{x} = -\underline{\omega}t + \underline{\vec{k}}\cdot\underline{\vec{x}}$$

must be constant. This shows that $(\frac{\omega}{c}, \vec{k})$ form a four vector

$$K^m = (\frac{\omega}{c}, \vec{k})$$

So $K_\mu X^\mu = \underline{K}_\mu \underline{X}^\mu = -\omega t + \vec{k}\cdot\vec{x}$ is const.

Velocity transformation rules & Four velocity



$$\underline{u} = \frac{d\underline{x}}{dt}$$

$$\underline{u}' = \frac{d\underline{x}'}{dt'}$$

Use: $\underline{x}' = L^{\mu}_{\nu} x^{\nu}$

$$\underline{u}' = \frac{u' - v}{1 + \vec{u}\vec{v}/c^2} \quad \text{etc'}$$

$$\underline{u}'_{\perp} = \frac{u_{\perp}}{\gamma(1 + \vec{u}\vec{v}/c^2)}$$

Complicated, because ^{both} the numerator and denominator transform. Consider a particle moving, $d\underline{x}' = (dx^0, d\underline{x}') = dx^0 (1, \vec{\beta}'_u)$

$$ds^2 = -(dx^0)^2 + d\underline{x}'^2 = d\underline{X}'_{\mu} d\underline{X}'^{\mu}$$

$$ds^2 = -(dx^0)^2 (1 - \vec{\beta}'_u{}^2)$$

\uparrow ds^2 is invariant. In ^{the} instantaneous rest frame of particle $ds^2 = -c^2 d\tau^2 + d\underline{x}'^2$

proper time of particle

So one finds that $d\tau$ is a Lorentz invariant

$$d\tau = dt \sqrt{1 - \beta^2} = \frac{dt}{\gamma_u} = d\tau$$

So define the four velocity: (distance per proper time)

$$U^\mu = \frac{dX^\mu}{d\tau} = \left(\gamma_u \frac{dx^0}{dt}, \gamma_u \frac{d\vec{x}}{dt} \right)$$

$$U^\mu = (\gamma_u c, \gamma_u \vec{u})$$

and

$$U^\mu = L^\mu_{\nu} U^\nu$$

As we will see the energy and momentum are proportional to the four velocity.

Energy momentum and velocity

• For a particle at rest

$$E = mc^2$$

• For a particle moving slowly

$$U = \frac{c^2 \vec{p}}{E} \approx \frac{\vec{p}}{m}$$