6.2 Magnetic Matter

Basic equations

(a) We are considering materials in the presence of a magnetic field. We write j_{mat} (the medium (material) currents) as an expansion in terms of the derivatives in the magnetic field. For weak fields, and an isotropic medium, the lowest term in the derivative expansion, for a parity and time-reversal invariant material is

$$\frac{\boldsymbol{j}_{\text{mat}}}{c} = \chi_m^B \nabla \times \boldsymbol{B} \tag{6.33}$$

where we have inserted a factor of c for later convenience.

(b) The current takes the form

$$\frac{\boldsymbol{j}_{\text{mat}}}{c} = \nabla \times \boldsymbol{M} \tag{6.34}$$

- i) M is known as the magnetization, and can be interpreted as the magnetic dipole moment per volume.
- ii) We have worked with linear response for an isotropic medium where

$$\boldsymbol{M} = \chi_m^B \boldsymbol{B} \tag{6.35}$$

This is most often what we will assume.

iii) Usually people work with H (see the next items (c), (d) for the definition of H) not B^{-1}

$$\boldsymbol{M} = \chi_m \boldsymbol{H} \tag{6.36}$$

iv) For not-that soft ferromagnets M(B) can be a very non-linear function of B. This will need to be specified (usually by experiment) before any problems can be solved. Usually this is expressed as the magnetic field as a function of H

$$\boldsymbol{B}(\boldsymbol{H}) \tag{6.37}$$

where H is small (of order gauss) and B is large (of order Tesla)

(c) After specifying the currents in matter, Maxwell equations take the form

$$\nabla \times \boldsymbol{B} = \nabla \times \boldsymbol{M} + \frac{\boldsymbol{j}_{ext}}{c} \tag{6.38}$$

$$\nabla \cdot \boldsymbol{B} = 0 \tag{6.39}$$

or

$$\nabla \times \boldsymbol{H} = \frac{\boldsymbol{j}_{ext}}{2} \tag{6.40}$$

$$\nabla \cdot \boldsymbol{B} = 0 \tag{6.41}$$

where 2

$$\boldsymbol{H} = \boldsymbol{B} - \boldsymbol{M} \tag{6.43}$$

² In the MKS system one has $H_{MKS} = \frac{1}{\mu_o} B_{MKS} - M_{MKS}$ so that B and H have different units. In a system of units where $\varepsilon_o = 1$ (so $1/\mu_o = c^2$) we have $H_{HL} = H_{MKS}/c$, $M_{HL} = M_{MKS}/c$ or since $1/c = \sqrt{\mu_o}$:

$$\boldsymbol{H}_{HL} = \sqrt{\mu_o} \boldsymbol{H}_{MKS} \qquad \boldsymbol{M}_{HL} = \sqrt{\mu_o} \boldsymbol{M}_{MKS} \tag{6.42}$$

¹There are a couple of reasons for this. One reason is because the parallel components of H are continuous across the sample. But, ultimately it is B which is the curl A, and it is ultimately the average current which responds to the gauge potential, through a retarded medium current-current correlation function that we wish to categorize.

(d) For linear materials :

$$\boldsymbol{B} = \mu \boldsymbol{H} = \frac{1}{1 - \chi_m^B} \boldsymbol{H} = (1 + \chi_m) \boldsymbol{H}$$
(6.44)

Implying the definitions

$$\mu \equiv \frac{1}{1 - \chi_m^B} \equiv (1 + \chi_m) \tag{6.45}$$

Solving magnetostatic problems with linear magnetic media:

All of the methods described in Sect. (6.1) will work with minor modifications due to the boundary conditions described below

(a) For linear materials in the coulomb gauge we get

$$\nabla \times \boldsymbol{H} = \mu \frac{\boldsymbol{j}_{ext}}{c} \tag{6.46}$$

$$\nabla \cdot \boldsymbol{B} = 0 \tag{6.47}$$

and with $\boldsymbol{B} = \nabla \times \boldsymbol{A}$ and constant μ we find

$$-\nabla^2 \mathbf{A} = \mu \frac{j_{ext}}{c} \tag{6.48}$$

which can be solved using the methods of magnetostatics.

(b) To solve magneto static equations we have boundary conditions:

$$\boldsymbol{n} \times (\boldsymbol{H}_2 - \boldsymbol{H}_1) = \frac{\boldsymbol{K}_{\text{ext}}}{c} \tag{6.49}$$

$$\boldsymbol{n} \cdot (\boldsymbol{B}_2 - \boldsymbol{B}_1) = 0 \tag{6.50}$$

i.e. if there are no external currents then the parallel components of H are continuous and the perpendicular components of B are continuous.

(c) At an interface the there are bound currents which are generated

$$\boldsymbol{n} \times (\boldsymbol{M}_2 - \boldsymbol{M}_1) = \frac{\boldsymbol{K}_{\text{mat}}}{c} \tag{6.51}$$