

## Maxwell Equations + Induction + Energy in Mag fields

$$\nabla \cdot \vec{E} = \rho_{\text{mat}} + \rho_{\text{ext}}$$

$$\nabla \times \vec{B} = \frac{\vec{j}_{\text{mat}}}{c} + \frac{\vec{j}_{\text{ext}}}{c} + \frac{1}{c} \partial_t \vec{E}$$

$$\nabla \cdot \vec{B} = 0$$

$$-\nabla \times \vec{E} = \frac{1}{c} \partial_t \vec{B}$$

Then for the material current we write

$$\frac{\vec{j}_{\text{mat}}}{c} = \overset{0 \text{ for insulator}}{\sigma} \vec{E} + \frac{1}{c} \partial_t \vec{P} + \nabla \times \vec{m}$$

Then with continuity,  $\rho_{\text{mat}} = -\nabla \cdot \vec{P}$ ,  $\vec{D} \equiv \vec{E} + \vec{P}$ ,  $\vec{H} \equiv \vec{B} - \vec{m}$ ,  
find

$$\nabla \cdot \vec{D} = \rho_{\text{ext}}$$

$$\nabla \times \vec{H} = \frac{\vec{j}_{\text{ext}}}{c} + \frac{1}{c} \partial_t \vec{D}$$

$$\nabla \cdot \vec{B} = 0$$

$$-\nabla \times \vec{E} = \frac{1}{c} \partial_t \vec{B}$$

← maxwell eqs  
in simple matter

Last Time pg. 2 + Induction pg. 1

Then we expand in powers of  $c$

Electrostatics

$$\nabla \cdot D^{(0)} = \rho_{\text{ext}}$$

$$\nabla \times E^{(0)} = 0$$

Magnetostatics:

$$\nabla \times H^{(1)} = j_{\text{ext}}/c + \frac{1}{c} \partial_t D^{(0)}$$

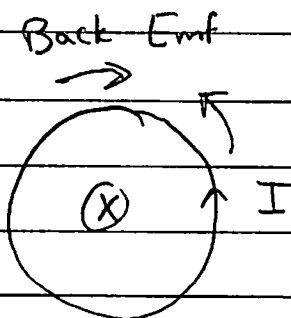
$$\nabla \cdot B^{(1)} = 0$$

Induced Electric fields / Back Emf

$$\nabla \cdot D^{(2)} = 0$$

$$-\nabla \times E^{(2)} = \frac{1}{c} \partial_t B^{(1)} \quad \text{could call } E^{(2)} = E^{\text{ind}}$$

Want to compute the energy stored in magnetic field



Imagine slowly increasing the current. Changing the current makes a changing magnetic field inducing a Back Emf. The work the Battery does to increase the current is the energy stored in the fields.

# Induction pg. 2

almost  
Take  $\nabla$  magneto statics, i.e.  $D^{(0)} = 0$

$$\begin{aligned}\nabla \times \mathbf{H} &= \mathbf{j} \\ \nabla \cdot \mathbf{B} &= 0\end{aligned}$$

And

$$-\nabla \times \mathbf{E}^{\text{ind}} = \frac{1}{c} \partial_t \mathbf{B}$$

Then the work by battery is

$$\frac{\delta U}{\delta t} = \frac{\delta W_{\text{batt}}}{\delta t} = - \int \mathbf{j} \cdot \delta \mathbf{E}^{\text{ind}} / \delta t$$

$\nwarrow -\vec{F} \cdot \vec{v}$

$$= - \int (\nabla \times \mathbf{H}) \cdot c \delta \mathbf{E}^{\text{ind}}$$

$$= - \int_V \vec{H} \cdot c \nabla \times \delta \mathbf{E}^{\text{ind}}$$

$\left. \begin{aligned} \nabla \cdot (\mathbf{H} \times \mathbf{E}) &= (\nabla \times \vec{H}) \cdot \mathbf{E} \\ &- \mathbf{H} \cdot \nabla \times \vec{E} \end{aligned} \right\}$

$$\frac{\delta U}{\delta t} = \int_V \vec{H} \cdot \frac{\delta \vec{B}}{\delta t}$$

$$\delta U = \int_V \vec{H} \cdot \delta \vec{B}$$

# Induction pg. 3

Then for linear media  $\delta B = \mu \delta H$

$$U = \frac{1}{2} \int_V \frac{H^2}{\mu} = \boxed{\frac{1}{2} \int_V \vec{H} \cdot \vec{B} \, d^3x = U}$$

These equations are often expressed in terms of  $\vec{J}$  and  $\vec{A}$  rather than  $\vec{B}$

Indeed,

$$\delta U = \int_V \vec{H} \cdot \delta \vec{B}$$

$$\delta E^{\text{ind}} = -\frac{1}{c} \partial_t \delta \vec{A}$$

$$\delta U = \int_V \vec{H} \cdot \nabla \times \delta \vec{A}$$

$$\delta U = \int_V \underbrace{\nabla \times \vec{H}}_{\frac{\vec{J}}{c}} \cdot \delta \vec{A}$$

By parts (no minus) because cross prod

$$\boxed{\delta U_B = \int_V \frac{\vec{J}}{c} \cdot \delta \vec{A}}$$

For linear media  $\delta \vec{A} \propto \mu \delta \vec{J}$

$$\boxed{U = \frac{1}{2} \int_V \frac{\vec{J}}{c} \cdot \vec{A}}$$

Wires pg. 1

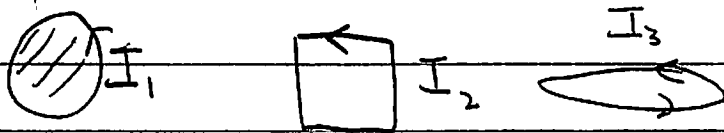
## Inductance in Wires

$$\textcircled{1} \quad U_B = \frac{1}{2} \int_V \vec{j} \cdot \vec{A} \, d^3x = \frac{1}{2} \int_V \vec{H} \cdot \vec{B}$$

•  $U_B$  is a property of state

$$\textcircled{2} \quad \delta U_B = \int_V \frac{\vec{j}}{c} \cdot \delta \vec{A}$$

For a set of wires:  $\vec{j} \, d^3x = I \, d\vec{l}$



Then find  $\leftarrow$  summed over  $a = \text{loops}$

$$\textcircled{1} \quad U = \frac{1}{2} \frac{I_a}{c} \Phi_a \quad \Phi_a = \oint_{\substack{a\text{-th} \\ \text{loop}}} \vec{A} \cdot d\vec{l} = \int \vec{B} \cdot d\vec{A}$$

$$\textcircled{2} \quad \delta U = \frac{I_a}{c} \delta \Phi_a = \text{flux through } a\text{-th loop}$$

Note that,  $\vec{A}(x) = \mu \int \frac{\vec{j}(x_0)}{4\pi |x-x_0|}$

$$U_B = \frac{\mu}{2} \int d^3x \, d^3x_0 \frac{\vec{j}(x) \cdot \vec{j}(x_0)}{4\pi |x-x_0|}$$

Wires pg. 2

So for a set of wires

$$U = \frac{1}{2} \bar{I}_a M_{ab} \bar{I}_b$$

↖ inductance matrix

$M_{11}$  is the self inductance of the first loop

$M_{12}$  is the mutual inductance between the 1st + 2nd

Then since  $U_B = \frac{1}{2} \frac{\bar{I}_a \bar{\Phi}_a}{C}$

$$\frac{\bar{\Phi}_a}{C} = M_{ab} \bar{I}_b$$

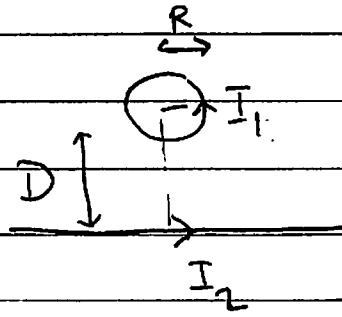
↖ back emf

And for any circuit

$$\mathcal{E}_a = -\frac{1}{C} \frac{d\bar{\Phi}_B}{dt} = -M_{ab} \frac{d\bar{I}_b}{dt}$$

## Problem on Mutual Inductance & Force

- Compute the mutual inductance of a ring and a long straight wire



Solution

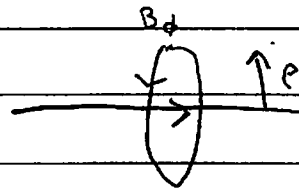
current in wire one

field from wire 2 at wire 1

$$U_{12} = \int_C \vec{j}_1 \cdot \vec{A}_2$$
$$= \frac{I_1}{c} \int_C \vec{A}_2 \cdot d\vec{l}_1 = \frac{I_1}{c} \int \vec{B}_2 \cdot d\vec{a}_1$$

Then the field from the wire is

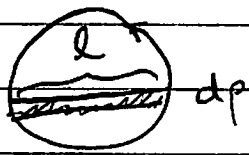
$$\vec{B}_2 = \frac{I_2}{2\pi\rho} \hat{\phi}$$



So we need to integrate this field from the wire over the area of the ring

# Mutual Inductance & Force pg. 2

points out given by circulation of  $I_1$   $l$  (see picture)



$$\vec{B}_2 \cdot \vec{n} da = \frac{I_2 l}{2\pi \rho} \cdot 2(R^2 - (\rho - D)^2)^{1/2} d\rho$$

points out

We have

$$U_{12} = \int_{D-R}^{D+R} d\rho \frac{I_1 I_2}{c^2} \frac{2(R^2 - (\rho - D)^2)^{1/2}}{2\pi \rho}$$

$$U_{12} = \frac{I_1 I_2}{c^2} [D - \sqrt{D^2 - R^2}]$$

$$\text{So } M_{12} = \frac{1}{c^2} (D - \sqrt{D^2 - R^2})$$

Then we might want to compute the force between the ring and the wire. To do this we ask about the change in  $U_B$ , as the distance between the ring and the wire is changed:

$$\delta U_B = \underbrace{I_a \delta \Phi_a}_{\substack{\text{change in energy} \\ \text{Stored in fields}}} + \delta W_{\text{mech}} \quad \begin{array}{l} \leftarrow \text{work done on system} \\ \text{by mechanical forces} \end{array}$$

$$= -\vec{F}_{\text{ring}} \cdot \vec{\delta D}$$

$\vec{F}_{\text{ring}} = -\vec{F}_{\text{applied}}$   
 force on ring      applied mechanically

$\delta W_{\text{batt}}$   
 work done by battery to keep current fixed



# Mutual Inductance and Force pg. 3

$$\delta U = \frac{1}{2} I_a \delta M_{ab} I_b$$

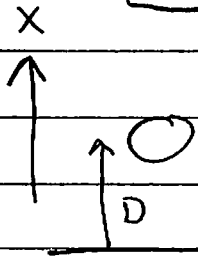
$$I_a \delta \Phi_a = I_a \delta M_{ab} I_b$$

So

$$\delta U_B - I_a \delta \Phi_a = -\frac{1}{2} I_a \delta M_{ab} I_b = -\vec{F} \delta \vec{D}$$

So

$$F^x = + \frac{\delta M_{ab}}{\delta D} \frac{I_a I_b}{2} = \frac{I_1 I_2}{c^2} \left( 1 - \frac{D}{\sqrt{D^2 - R^2}} \right)$$



$$= -\frac{I_1 I_2}{c^2} \left( \frac{D}{\sqrt{D^2 - R^2}} - 1 \right)$$

indicates an attractive force  
i.e. force in negative x-direction

