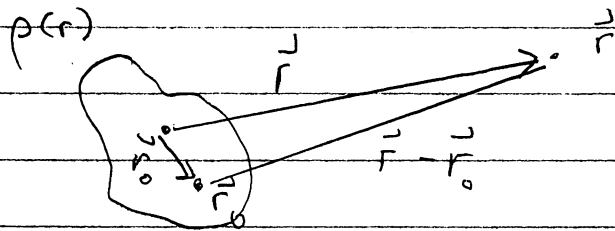


Multipole Expansion with Spherical Harmonics



Let us redo
the multipole
expansion

Then

$$\phi(\vec{r}) = \int d^3 \vec{r}_0 \frac{\rho(\vec{r}_0)}{4\pi |\vec{r} - \vec{r}_0|}$$

For $r \gg r_0$ we can expand $r_{>} = r$ and $r_{<} = r_0$ and we have the expansion

$$\frac{1}{4\pi |\vec{r} - \vec{r}_0|} = \sum_{lm} \frac{r_0^l}{r^{l+1}} \frac{1}{(2l+1)} Y_{lm}(\theta, \phi) Y_{lm}^*(\theta_0, \phi_0)$$

This leads to

$$\phi(r) = \sum_{lm} \frac{q_{lm}}{(2l+1)} \frac{Y_{lm}(\theta, \phi)}{r^{l+1}} = \frac{q_{00}}{r} Y_{00} + \frac{1}{3} \frac{q_{1m}}{r^2} Y_{1m} + \mathcal{O}\left(\frac{1}{r^3}\right)$$

where

$$q_{lm} = \int d^3 r_0 \rho(\vec{r}_0) r_0^l Y_{lm}^*(\theta_0, \phi_0)$$

↑
spherical multipole moment

This multipole expansion is entirely equivalent to the expansion we had previously

$$\varphi(r) = \frac{Q_{\text{Tot}}}{4\pi r} + \frac{\vec{p} \cdot \hat{r}}{4\pi r^2} + \frac{Q_{ij}(\hat{r}_i \hat{r}_j - \frac{1}{3} \delta_{ij})}{4\pi r^3} + \dots$$

To see this one needs to understand what $Y_{\ell m}(\theta, \phi)$ are. $Y_{\ell m}$ are linearly combos of the components of a symmetric traceless ℓ -th rank tensor constructed out of \hat{r}

Cartesian	Spherical	Rank
1	Y_{00}	0
\hat{r}_i	Y_{1m}	1
$\hat{r}_i \hat{r}_j - \frac{1}{3} \delta_{ij}$	Y_{2m}	2
$\frac{\hat{r}_i \hat{r}_j \hat{r}_k}{5} (\hat{r}_i \delta_{jk} + \hat{r}_j \delta_{ki} + \hat{r}_k \delta_{ij})$	Y_{3m}	3

And so on.

To understand my meaning, take the dipole term:

$$Y_{11} \propto (\hat{x} + i\hat{y})$$

$$Y_{1-1} \propto (\hat{x} - i\hat{y})$$

$$Y_{10} \propto \hat{z}$$

We see that Y_{1m} is a linear combo of \hat{r}^i

Similarly q_{1m} is a linear combo of \vec{p} , e.g.:

$$q_{11} = \int d^3r_0 \underbrace{r_0 Y_{1m}^*}_{\propto (x-iy)\rho(r_0)} \rho(r_0) \propto \underbrace{p^x - ip^y}_{\text{The } x, y \text{ components of } \vec{p}}$$

since

$$p^x \equiv \int d^3\vec{r}_0 x \rho(\vec{r}_0) \text{ etc.}$$

The relation between p^i and q_{1m} is the same as the relation (i.e. linear-combo) between \hat{r}^i and Y_{1m}^*

The relations and normalizations are chosen so that the series agree, e.g.

$$\frac{\vec{p} \cdot \hat{r}}{4\pi r^2} = \sum_m \frac{1}{3} q_{1m} \frac{Y_{1m}}{r^2} \Rightarrow \vec{p} \cdot \hat{r} = \frac{4\pi}{3} \sum_m q_{1m} Y_{1m}$$

↖ $2l+1$ for $l=1$

This is the statement that

$$\vec{p} \cdot \hat{r} = \frac{(p_x - ip_y)}{\sqrt{2}} \left(\frac{\hat{r}_x + i\hat{r}_y}{\sqrt{2}} \right) + \frac{(p_x + ip_y)}{\sqrt{2}} \left(\frac{\hat{r}_x - i\hat{r}_y}{\sqrt{2}} \right) + p_z \hat{r}_z$$

$$= \frac{4\pi}{3} (q_{11}^* Y_{11} + q_{1-1}^* Y_{1-1} + q_{10} Y_{10})$$

Similarly Y_{2m} is a linear combo of $\hat{r}^i \hat{r}^j - \frac{1}{3} \delta^{ij}$

(There are five components of $\hat{r}^i \hat{r}^j - \frac{1}{3} \delta^{ij}$, and five $l=2$ spherical harmonics). And, q_{2m} is a linear combo of the quadrupole tensor Q_{ij} components (The map between q_{2m} and Q_{ij} is the same as between Y_{2m}^* and $\hat{r}^i \hat{r}^j - \frac{1}{3} \delta^{ij}$). Then this map is constructed so that

$$\frac{1}{4\pi r^3} Q_{ij} \left(\hat{r}^i \hat{r}^j - \frac{1}{3} \delta^{ij} \right) = \sum_m \frac{1}{5} q_{2m} Y_{2m} / r^3$$

↖ $2l+1$ with $l=2$