1.1 Linear operators, boundary conditions, and adjoints

Consider the simplest of all differential operator

$$\mathcal{L}_t = \begin{bmatrix} \frac{d}{dt} \end{bmatrix} \tag{1.1}$$

The operator is fully specified by its action and the space of functions on which it acts, *i.e.* the boundary conditions and domain of the functions which we are talking about. For instance lets take the space of functions between t_{\min} and t_{\max} with retarded boundary conditions, *i.e.* v(t) = 0 for $t \leq t_{\min}$.

$$\mathcal{L}_t v(t) = \left[\frac{d}{dt}\right] v(t) \qquad v(t_{\min}) = 0 \tag{1.2}$$

Numerically we discretize using backward differences

$$\frac{dv}{dt} = \frac{v(t) - v(t - \Delta t)}{\Delta t} \tag{1.3}$$

Then the action of the differential operator can be put in matrix form

$$\mathcal{L}_{t}v(t) = \begin{pmatrix} 1/\Delta t & 0 & & \\ -1/\Delta t & 1/\Delta t & & \\ & -1/\Delta t & \ddots & & \\ & & \ddots & 1/\Delta t & 0 \\ & & & -1/\Delta t & 1/\Delta t \end{pmatrix} \begin{pmatrix} v(t_{1}) \\ v(t_{2}) \\ v(t_{3}) \\ \vdots \\ v(t_{N}) \end{pmatrix}$$
(1.4)

where $t_n = t_{\min} + n \Delta$. The last point is at $t_N = t_{\max} - \Delta$

Can you see how the retarded boundary conditions are used when placing the entries into the matrix – basically we said

$$\left. \frac{dv(t)}{dt} \right|_{t_1} = \frac{v(t_1)}{\Delta t} \tag{1.5}$$

which is true if $v(t_{\min}) = 0$.

The adjoint differential operator is a unique operator satisfying

$$\int_{t_{\min}}^{t_{\max}} dt \, V_2(t) \left(\mathcal{L}_t v_1(t) \right) = \int dt \left(\mathcal{L}_t^{\mathrm{adj}} V_2(t) \right) v_1(t) \tag{1.6}$$

Here $v_1(t)$ satisfies retarded boundary conditions (that is $v_1(t_{\min}) = 0$) and $V_2(t)$ satisfies the adjoint boundary conditions, which (as we will show below) are advanced boundary conditions $V_2(t_{\max}) = 0$ for the particular boundary conditions satisfied by \mathcal{L}_t .

As you know the adjoint is found by taking the hermitian transpose

$$\mathcal{L}_{t}^{\mathrm{adj}}v(t) = \begin{pmatrix} 1/\Delta t & -1/\Delta t & 0 & 0 & 0\\ 0 & 1/\Delta t & \ddots & 0 & 0\\ 0 & 0 & \ddots & -1/\Delta t & 0\\ 0 & 0 & 0 & 1/\Delta t & -1/\Delta t\\ 0 & 0 & 0 & 0 & 1/\Delta t \end{pmatrix} \begin{pmatrix} v(t_{0}) \\ v(t_{1}) \\ v(t_{2}) \\ \vdots \\ v(t_{N}) \end{pmatrix}$$
(1.7)

We see that the adjoint differential operator is

$$\mathcal{L}_t^{\mathrm{adj}} = -\frac{d}{dt} \tag{1.8}$$

This only half specifies the operator. The full specification is for the space of functions v(t) between $t_{\min} \dots t_{\max}$ which vanish at the upper end

$$\mathcal{L}_t^{\text{adj}} = -\frac{d}{dt} \qquad v(t_{\text{max}}) = 0 \tag{1.9}$$

This follows because the discretization implies that

$$-\frac{dv(t)}{dt}\Big|_{t_N} = \frac{v(t_N)}{\Delta t}$$
(1.10)

which is true provided $v(t_{\text{max}}) = 0$.

Without appealing to the discretization, the appropriate adjoint operator, and its boundary conditions are found by taking any function V(t) and integrating by parts:

$$\int_{t_{\min}}^{t_{\max}} dt \, V(t)(\frac{d}{dt}v_1(t)) = \underbrace{\left[V(t_{\max})v_1(t_{\max}) - V(t_{\min})v_1(t_{\min})\right]}_{\text{bndry terms}} + \int_{t_{\min}}^{t_{\max}} dt \left[-\frac{d}{dt}V(t)\right] v_1(t) \quad (1.11)$$

Ignoring the boundary terms the adjoint operator is

$$\mathcal{L}_t^{\mathrm{adj}} = \left[-\frac{d}{dt} \right] \tag{1.12}$$

But, the full specification also specifies that that $V(t_{\text{max}}) = 0$, so that the boundary terms vanish:

$$\mathcal{L}_t^{\text{adj}} = \left[-\frac{d}{dt} \right] \qquad V(t_{\text{max}}) = 0.$$
(1.13)

We conclude with examples.

Example operators and their adjoints

(a) a particle experiencing drag with time dependent mass and drag coefficient with retarded boundary conditions

$$\mathcal{L}_t = \left[m(t) \frac{d}{dt} + m(t)\eta(t) \right] \qquad v(t_{\min}) = 0 \tag{1.14}$$

The adjoint is

$$\mathcal{L}_t^{\text{adj}} = \left[-\frac{d}{dt} m(t) + m(t)\eta(t) \right] \qquad v(t_{\text{max}}) = 0 \tag{1.15}$$

(b) The damped harmonic oscillator with time dependent mass and drag coefficient

$$\mathcal{L}_t = \left[\frac{d}{dt}m(t)\frac{d}{dt} + \frac{d}{dt}m(t)\eta(t) + m(t)\omega_o^2\right] \qquad v(t) = 0 \quad t \le t_{\min}$$
(1.16)

The adjoint is

$$\mathcal{L}_t^{\mathrm{adj}} = \left[\frac{d}{dt}m(t)\frac{d}{dt} - m(t)\eta(t)\frac{d}{dt} + m(t)\omega_o^2\right] \qquad v(t) = 0 \ge t_{\max} \tag{1.17}$$

(c) The Sturm-Liouville operator with homogeneous boundary conditions

$$\mathcal{L}_x = \left[-\frac{d}{dx} p(x) \frac{d}{dx} + q(x) \right] \qquad y(x_{\min}) = y(x_{\max}) = 0 \tag{1.18}$$

where p(x) > 0 between x_{\min} and x_{\max} The operator is self adjoint, $\mathcal{L}_x^{\mathrm{adj}} = \mathcal{L}_x$.

Self Test:

(a) Show using any of the methods of this section that, the adjoint of Eq. (1.16) is Eq. (1.17).

1.2 The $G_R(t, t_o)$ is adjoint respect to t_o

We now show that $G_R(t, t_o)$ when considered a function of t_o satisfies the adjoint differential equation and boundary conditions, *i.e.* since

$$\mathcal{L}_t G_R(t, t_o) = \delta(t - t_o) \tag{1.19}$$

we will show that

$$\mathcal{L}_{t_o}^{\mathrm{adj}}G_R(t,t_o) = \delta(t-t_o) \tag{1.20}$$

The proof of this goes like this. We start with

$$\int_{t_{\min}}^{t_{\max}} dt' G_R(t,t') \,\mathcal{L}_{t'} G_R(t',t'') = G_R(t,t'') \,. \tag{1.21}$$

If you think of $\mathcal{L}_{t'}$ as a matrix, then it is clear that

$$\int_{t_{\min}}^{t_{\max}} dt' (\mathcal{L}_{t'}^{\mathrm{adj}} G_R(t, t')) G_R(t', t'') = G_R(t, t'') \,. \tag{1.22}$$

Or

$$\mathcal{L}_{t'}^{\mathrm{adj}}G(t,t') = \delta(t-t') \tag{1.23}$$

More formally, we can integrate by parts picking up boundary terms and the adjoint operator

bndry-terms +
$$\int_{t_{\min}}^{t_{\max}} dt' \left[\mathcal{L}_{t'}^{\mathrm{adj}} G(t,t') \right] G(t',t'') = G(t,t'')$$
 (1.24)

The only way this can be satisfied for all t and t'' is if the boundary terms vanish (i.e. G(t, t') satisfies adjoint b.c. with respect to t'), and

$$\mathcal{L}_{t'}^{\mathrm{adj}}G(t,t') = \delta(t-t'') \tag{1.25}$$

1.3 Green Theorem

Now we can prove Green theorem. For definiteness, take the drag equation with retarded boundary conditions:

$$\underbrace{\left\lfloor \frac{d}{dt} + \eta \right\rfloor}_{\mathcal{L}_t} v(t) = F(t) \tag{1.26}$$

Given an initial condition at t = 0, $v(t_{\min}) = v_o$, and the Green function $G(t, t_o)$ we would like to write down the general (formal) solution to the equation.

The procedure to do this is always the same – act with the adjoint operator on the Green function and integrate by parts

$$v(t) = \int_{t_{\min}}^{t_{\max}} dt_o \, v(t_o) \mathcal{L}_{t_o}^{\mathrm{adj}} G_R(t, t_o) \tag{1.27}$$

$$= \int_{t_{\min}}^{t_{\max}} dt_o v(t_o) \left[-\frac{d}{dt_o} + \eta \right] G_R(t, t_o)$$
(1.28)

Integrating by parts we get

$$v(t) = -v(t_{\max})G_R(t, t_{\max}) + v(t_{\min})G_R(t, t_{\min}) + \int_{t_{\min}}^{t_{\max}} G(t, t_o) \left[\frac{d}{dt_o} + \eta\right] v(t_o)$$
(1.29)

The first term vanishes because $G_R(t, t_o)$ vanishes whenever $t < t_o$. The last term uses the equation of motion, Eq. (1.26), and the again the causality condition $G_R(t, t_o)$, leading to

$$v(t) = G_R(t, t_{\min})v(t_{\min}) + \int_{t_{\min}}^t dt_o G_R(t, t_o)F(t_o)$$
(1.30)

In the absence of the external force we get

$$v(t) = G_R(t, t_{\min})v(t_{\min})$$
(1.31)

Self test:

(a) For a damped simple harmonic oscillator

$$\left[m\frac{d^2}{dt^2} + m\eta\frac{d}{dt} + m\omega_o^2\right]x(t) = 0$$
(1.32)

that the solution to the equation of motion (analogous to Eq. (1.31)) is

$$x(t) = m \left[G_R(t, t_o) \partial_{t_o} x(t_o) - x(t_o) \partial_{t_o} G_R(t, t_o) \right] + m \eta G_R(t, t_o) x(t_o)$$
(1.33)

(b) Check that Eq. (1.33) satisfies the equations of motion and the initial conditions. You will need to establish (by looking at the action of the adjoint) that

$$\lim_{t \to t_o} \left[m \partial_t \partial_{t_o} G_R(t, t_o) - m \eta \partial_t G_R(t, t_o) \right] = 0$$
(1.34)

What are the modifications if the parameters m and η depend on time?