

## Dot Product

$$\begin{aligned}\vec{a} \cdot \vec{b} &= (a^i \vec{e}_i) \cdot (b^j \vec{e}_j) \\ &= a^i b^j (\vec{e}_i \cdot \vec{e}_j) = a^i b^j \delta_{ij} = a_j b^j,\end{aligned}$$

- Contracted indices  $a_j b^j$  indicate dot product.   
  $\swarrow \searrow$    
  $\leftarrow j$ 's are "contracted"

- The dot product is invariant under rotations

$$\underline{a}_j \underline{b}^j = a_i (\mathcal{R}^{-1})^i_j (\mathcal{R})^j_k b^k = a_i \delta^i_k b^k = a_i b^i$$

## Cross Product

- For the cross product need the Levi-Civita tensor

$$\varepsilon^{ijk} = \begin{cases} \pm 1 & \text{if } (i, j, k) \text{ is an even/odd perm} \\ 0 & \text{otherwise} \end{cases} \quad \text{of } (1, 2, 3)$$

e.g.

$$\varepsilon^{123} = -\varepsilon^{213} = \varepsilon^{231} = -\varepsilon^{321} = \varepsilon^{312} \dots$$

$$\varepsilon^{223} = 0$$





b) The determinant is anti-symmetric under interchange of two rows

$$\det(\vec{m}^1, \vec{m}^2, \vec{m}^3) = -\det(\vec{m}^1, \vec{m}^3, \vec{m}^2)$$

Indeed:

$$\begin{aligned} m^1_i m^2_j m^3_k \varepsilon^{ijk} &= -m^1_i m^2_j m^3_k \varepsilon^{ikj} \\ &= -m^1_i m^3_k m^2_j \varepsilon^{ikj} \end{aligned}$$

Thus in general

$$(Eq \star) \quad m^l_i m^m_j m^n_k \varepsilon^{ijk} = \det \mathcal{M} \varepsilon^{lmn}$$

c) Homework, use Eq  $\star$  to show:

$$\det(AB) = \det(A) \det(B)$$

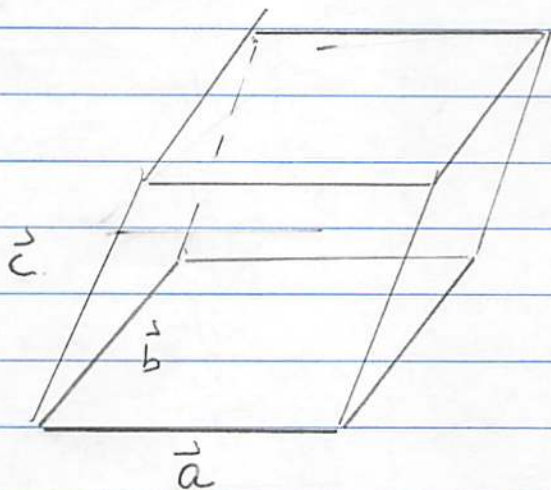
$$\det A^T = \det(A)$$

d) Why is the determinant important for physics? Take three vectors:

$$\vec{a}, \vec{b}, \vec{c}$$

$$\text{Then, } \det(\vec{a}, \vec{b}, \vec{c}) = \begin{vmatrix} a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \\ c_1 & c_2 & c_3 \end{vmatrix} = \text{volume}$$

is the volume of the parallel piped spanned by  $\vec{a}, \vec{b}, \vec{c}$



• (3) The cross product

$$\begin{aligned}\vec{a} \times \vec{b} &= \begin{vmatrix} \vec{e}_1 & \vec{e}_2 & \vec{e}_3 \\ a_1 & a_2 & a_3 \\ b_1 & b_2 & b_3 \end{vmatrix} = \vec{e}_i (\epsilon^{ijk} a_j b_k) \\ &= \vec{e}_i (\vec{a} \times \vec{b})^i\end{aligned}$$

Thus

$$(\vec{a} \times \vec{b})^i = \epsilon^{ijk} a_j b_k = i\text{-th contravariant component of } \vec{a} \times \vec{b}$$

## The $b(ac)$ -abc rule

$$\vec{a} \times (\vec{b} \times \vec{c}) = \vec{b} (\vec{a} \cdot \vec{c}) - (\vec{a} \cdot \vec{b}) \vec{c}$$

Proof

$$\begin{aligned} (\vec{a} \times (\vec{b} \times \vec{c}))^i &= \epsilon^{ijk} a_j (\vec{b} \times \vec{c})_k \\ &= \epsilon^{ijk} a_j \epsilon_{klm} b^l c^m \end{aligned}$$

Now analyze and think

$$\epsilon^{ijk} \epsilon_{klm} = \boxed{\epsilon^{ijk} \epsilon_{lmk} = \delta^i_l \delta^j_m - \delta^i_m \delta^j_l}$$

if  $k=3$ , and  $i, j=1, 2$  and  $l, m$  can be either 1, 2 or 2, 1. So  $l, m$  is  $i, j$  or  $j, i$  which is what this expression says

So

$$\begin{aligned} (\vec{a} \times (\vec{b} \times \vec{c}))^i &= a_j b^l c^m (\delta^i_l \delta^j_m - \delta^i_m \delta^j_l) \\ &= b^i (a_m c^m) - (a_l b^l) c^i \\ &= b^i (\vec{a} \cdot \vec{c}) - (\vec{a} \cdot \vec{b}) c^i \end{aligned}$$



## Derivative Operations:

$$\text{grad} = (\nabla \vec{S})_i = \partial_i S$$

$$\partial_i \equiv \frac{\partial}{\partial x^i}$$

$$\text{curl} = (\nabla \times \vec{V})^i = \epsilon^{ijk} \partial_j V_k$$

$$\text{div} = \nabla \cdot \vec{V} = \partial_i v^i = \partial_x v^x + \partial_y v^y + \partial_z v^z$$

$$\text{laplacian} \quad \nabla \cdot \nabla S = \partial_i \partial^i$$

The  $b(ac) - (ab)c$  rule plays an important role:

$$\nabla \times (\nabla \times \vec{C}) = \nabla (\nabla \cdot \vec{V}) - \nabla^2 \vec{C}$$

- Homework: use the  $b(ac) - (ab)c$  rule to derive the wave equation

## Example of gradient.

• Field of a point charge,  $-\vec{\nabla} \frac{1}{r} = \frac{\vec{r}}{r^2} = \frac{\vec{r}}{r^3}$

Prf

$$-\frac{\partial}{\partial x^i} \left( \frac{1}{r} \right) = -\frac{\partial}{\partial x^i} \frac{1}{(x^l x_l)^{1/2}}$$
$$= +\frac{1}{2} \frac{1}{(x^l x_l)^{3/2}} \frac{\partial (x^l x_l)}{\partial x^i}$$

Use

$$\frac{\partial x^l}{\partial x^i} = \delta_i^l \quad \text{so} \quad \frac{\partial (x^l x_l)}{\partial x^i} = \delta_i^l x_l + x^l \delta_{ij}$$
$$= 2x_i$$

Thus,

$$\left( -\vec{\nabla} \frac{1}{r} \right)_i = +\frac{1}{2} \frac{1}{(x^l x_l)^{3/2}} \cancel{2} x_i$$
$$= \frac{x_i}{r^3}$$

Use

$$r = (x^l x_l)^{1/2}$$

i.e

$$-\vec{\nabla} \left( \frac{1}{r} \right) = \frac{\vec{r}}{r^3}$$

where  $\vec{r} = x_i \vec{e}^i$

$$r = \sqrt{x^l x_l}$$

## Helmholtz Theorems:

① If  $\vec{\nabla} \cdot \vec{C} = 0$ , then there exists  $\vec{D}$  such that:

$$\vec{C} = \vec{\nabla} \times \vec{D}.$$

② If  $\vec{\nabla} \times \vec{C} = 0$ , then there exists a scalar field  $S$  such that:

$$\vec{C} = -\vec{\nabla} S,$$

I won't prove it (but see homework) but I will show the converse, i.e.

$$\text{① } \vec{\nabla} \cdot (\vec{\nabla} \times \vec{D}) = 0 \quad \text{and} \quad \text{② } \vec{\nabla} \times (\vec{\nabla} S) = 0$$

Prf.

$$\text{① } \partial_i C^i = \partial_i \overbrace{\varepsilon^{ijk} \partial_j D_k}^{(\vec{\nabla} \times \vec{C})^i} = \varepsilon^{ijk} \partial_i \partial_j D_k = 0$$

Because  $\varepsilon^{ijk} = -\varepsilon^{jki}$  is antisymmetric while  $\partial_i \partial_j = \partial_j \partial_i$  is symmetric,  $\partial_x \partial_y - \partial_y \partial_x = 0$ .

② Similarly, we show  $\vec{\nabla} \times \vec{\nabla} S = 0$

$$\underbrace{\varepsilon^{ijk} \partial_j C_k}_{(\vec{\nabla} \times \vec{C})^i} = \varepsilon^{ijk} \partial_j \partial_k S = 0$$

These are statements of differential forms  $d \circ d D = 0$