

## Hamilton's Equations from a variational procedure

- So far we derived Hamilton's Equations from the Lagrange equations of motion and the definition of the Hamiltonian. We can also derive the EOM from the action  $H = p\dot{q} - L$  or  $L = p\dot{q} - H$

$$S[q, p] = \int dt L = \int dt p \dot{q} - H(q, p)$$

- Now however we will independently vary  $p$  and  $q$  while keeping  $q$  fixed at ends (previously we kept  $q$ -fixed but  $\dot{q}$  was varied at ends)

$$q \rightarrow q + \delta q(t)$$

$$\delta q(t_1) = \delta q(t_2) = 0$$

$$p \rightarrow p + \delta p(t)$$

$p(t_1) + p(t_2)$  free to change

OK then the action under the variation

$$S[q + \delta q, p + \delta p] = \int dt (p + \delta p) \left( \frac{d}{dt} (q + \delta q) \right)$$

So  $-H(q + \delta q, p + \delta p)$

$$\delta S = S[q + \delta q, p + \delta p] - S[q, p]$$

And

$$\delta S = \int \delta p \frac{dq}{dt} + p \frac{d\delta q}{dt} - \frac{\partial H}{\partial q} \delta q - \frac{\partial H}{\partial p} \delta p = 0$$

by parts

Integrating by parts:

$$\delta S = p \delta q \Big|_{t_1}^{t_2} + \int_{t_1}^{t_2} dt \left( \frac{dq}{dt} - \frac{\partial H}{\partial p} \right) \delta p - \left( \frac{dp}{dt} + \frac{\partial H}{\partial q} \right) \delta q$$

The endpoints vanish since  $\delta q(t_1) = \delta q(t_2) = 0$ ,  
Leading to (since  $\delta p$  and  $\delta q$  are arbitrary)

$$\frac{dq}{dt} = \frac{\partial H}{\partial p}$$

$$\frac{dp}{dt} = -\frac{\partial H}{\partial q}$$

Hamilton  
Equations