

# 1 Basic Mechanics

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## 1.1 Newtonian mechanics a brief review

### Momentum and Center of Mass

- Newton's equations of motion for a system of particles reads

$$\frac{d\mathbf{p}_a}{dt} = \mathbf{F}_a \quad (1.1)$$

where  $a = 1 \dots N$  labels the particles. Here  $\mathbf{p}_a = m_a \mathbf{v}_a$ . We usually divide up the forces on the  $a$ -th particle into external forces acting on the system from outside, and internal forces acting between pairs of particles:

$$\mathbf{F}_a = \underbrace{\mathbf{F}_a^{\text{ext}}}_{\text{external forces}} + \underbrace{\sum_{b \neq a} \mathbf{F}_{ab}}_{\text{internal forces}} \quad (1.2)$$

Here

$$\mathbf{F}_{ab} \equiv \text{Force on particle } a \text{ by } b, \quad (1.3)$$

and of course we have Newton's equal and opposite rule

$$\mathbf{F}_{ab} = -\mathbf{F}_{ba}. \quad (1.4)$$

- Summing over the particles we find (after using Eq. (1.4)) that the internal forces cancel and the total change in momentum per time is the sum of external forces

$$\frac{d\mathbf{P}_{\text{tot}}}{dt} = \mathbf{F}_{\text{tot}}^{\text{ext}} \quad (1.5)$$

where  $\mathbf{P}_{\text{tot}} = \sum_a \mathbf{p}_a$  and  $\mathbf{F}_{\text{tot}}^{\text{ext}} = \sum_a \mathbf{F}_a^{\text{ext}}$ . If there are no external forces then  $\mathbf{P}_{\text{tot}}$  is constant

- The velocity of the center of mass is

$$\mathbf{v}_{\text{cm}} = \frac{\mathbf{P}_{\text{tot}}}{M_{\text{tot}}} = \frac{1}{M_{\text{tot}}} \sum_a m_a \mathbf{v}_a. \quad (1.6)$$

The position of the center of mass (relative to an origin  $O$ ) is

$$\mathbf{R}_{\text{cm}} = \frac{1}{M_{\text{tot}}} \sum_a m_a \mathbf{r}_a. \quad (1.7)$$

### Angular momentum:

- Angular momentum is defined with respect to a specific origin  $O$  (i.e.  $\mathbf{r}_a$  depends on  $O$ ) which is not normally notated

$$\boldsymbol{\ell}_{a,O} \equiv \boldsymbol{\ell}_a \equiv \mathbf{r}_a \times \mathbf{p}_a. \quad (1.8)$$

It evolves as

$$\frac{d\boldsymbol{\ell}}{dt} = \mathbf{r}_a \times \mathbf{F}_a \quad (1.9)$$

- The total angular momentum  $\mathbf{L}_{\text{tot}} = \sum_a \boldsymbol{\ell}_a$  changes due to the total *external* torque

$$\frac{d\mathbf{L}_{\text{tot}}}{dt} = \boldsymbol{\tau}_{\text{tot}}^{\text{ext}}, \quad (1.10)$$

where  $\boldsymbol{\tau}_{\text{tot}}^{\text{ext}} = \sum_a \mathbf{r}_a \times \mathbf{F}_a^{\text{ext}}$  were we have generally assumed that the internal forces are radially directed  $\mathbf{F}_{ab} \propto (\mathbf{r}_a - \mathbf{r}_b)$

- The angular momentum depends on the origin  $O$ . Writing the position of the particle relative to the center of mass as  $\Delta\mathbf{r}_a$ , i.e.

$$\mathbf{r}_a = \mathbf{R}_{\text{cm}} + \Delta\mathbf{r}_a, \quad (1.11)$$

the angular momentum of the system about  $O$  is

$$\mathbf{L}_O = \underbrace{\mathbf{R}_{\text{cm}} \times \mathbf{P}_{\text{tot}}}_{\text{Ang-mom of center of mass about } O} + \underbrace{\sum_a \Delta\mathbf{r}_a \times \mathbf{p}_a}_{\text{Ang-mom about the cm}}. \quad (1.12)$$

## Energy

- Energy conservation is derived by taking the dot product of  $\mathbf{v}$  with  $d\mathbf{p}/dt$ . We find that the change in kinetic energy (on the  $a$ -the particle) equals the work done (on the  $a$ -particle).

$$\left. \frac{1}{2} m_a v_a^2(t) \right|_{t_1}^{t_2} = W_a \quad (1.13)$$

where the work is

$$W_a = \int_{\mathbf{r}_a(t_1)}^{\mathbf{r}_a(t_2)} \mathbf{F}_a \cdot d\mathbf{r}_a \quad (1.14)$$

- *Potential Energy.* For conservative forces the force can be written as (minus) the gradient of a scalar function which we call the potential energy

$$\mathbf{F}_a = -\nabla_{\mathbf{r}_a} U \quad (1.15)$$

Consider the potential energy  $U_{12}$  between particle 1 and 2. Since the force is equal and opposite

$$\mathbf{F}_{12} = -\nabla_{\mathbf{r}_1} U_{12}(\mathbf{r}_1, \mathbf{r}_2) = +\nabla_{\mathbf{r}_2} U_{12}(\mathbf{r}_1, \mathbf{r}_2) = -\mathbf{F}_{21} \quad (1.16)$$

and this is used to conclude that interaction potential between two particles is of the form

$$U_{12}^{\text{int}} = U(|\mathbf{r}_1 - \mathbf{r}_2|) \quad (1.17)$$

Typically we divide up the potential into an external potential and the internal ones

$$U(\mathbf{r}_a) = U^{\text{ext}}(\mathbf{r}_a) + \frac{1}{2} \sum_{ab, a \neq b} U_{ab}^{\text{int}}(\mathbf{r}_a, \mathbf{r}_b) \quad (1.18)$$

The sum over the internal potentials comes with a factor of a half because the energy between particle-1 and particle-2 is counted twice in the sum, e.g. for just two particles

$$U_{12}^{\text{int}}(\mathbf{r}_1, \mathbf{r}_2) = \frac{1}{2} (U(|\mathbf{r}_1 - \mathbf{r}_2|) + U(|\mathbf{r}_2 - \mathbf{r}_1|)) . \quad (1.19)$$

- *Energy.* The total energy is

$$E = \sum_a \frac{1}{2} m_a v_a^2 + U^{\text{ext}}(\mathbf{r}_a) + \frac{1}{2} \sum_{ab, a \neq b} U_{ab}^{\text{int}}(\mathbf{r}_a, \mathbf{r}_b) \quad (1.20)$$

and is constant if there are no non-conservative forces.

If there are non-conservative forces then

$$E(t_2) - E(t_1) = W_{\text{NC}} \quad (1.21)$$

where the work done by the non-conservative forces is  $W_{\text{NC}} = \sum_a \int \mathbf{F}_a^{\text{NC}} \cdot d\mathbf{r}_a$

- It is convenient to measure velocities relative to the center of mass

$$\mathbf{v}_a = \mathbf{v}_{\text{cm}} + \Delta \mathbf{v}_a \quad (1.22)$$

where  $\Delta \mathbf{v}_a = \dot{\Delta} \mathbf{r}_a$ , then the kinetic energy

$$K = \underbrace{\frac{1}{2} M_{\text{tot}} v_{\text{cm}}^2}_{\text{KE of center-mass}} + \underbrace{\sum_a \frac{1}{2} m_a \Delta v_a^2}_{\text{KE relative to center-mass}} \quad (1.23)$$

### Galilean invariance:

- Consider newtons laws then for an isolated system of particles

$$\frac{d\mathbf{p}_a}{dt} = \mathbf{F}_a \quad (1.24)$$

where  $\mathbf{F}_a = -\nabla_{\mathbf{r}_a} U$  with

$$U = \frac{1}{2} \sum_{ab, a \neq b} U_{ab}^{\text{int}}(|\mathbf{r}_a - \mathbf{r}_b|) \quad (1.25)$$

Here the space-time coordinates are measured by an observer  $O$  with origin.

Then consider an observer  $O'$  moving with *constant* velocity  $-\mathbf{u}$  relative to  $O$ . The “new” coordinates (those measured by  $O'$ ) are related to the old coordinates via a Galilean boost

$$\mathbf{r}_a \rightarrow \mathbf{r}'_a = \mathbf{r}_a + \mathbf{u}t \quad (1.26)$$

$$t \rightarrow t' = t \quad (1.27)$$

The potential which only depends on  $\mathbf{r}_a - \mathbf{r}_b$  is independent of the shift. The observer measures

$$\mathbf{v}_a \rightarrow \mathbf{v}'_a = \mathbf{v}_a + \mathbf{u} \quad (1.28)$$

$$\mathbf{p}_a \rightarrow \mathbf{p}'_a = \mathbf{p}_a + m_a \mathbf{u} \quad (1.29)$$

The equations of motion for observer  $O'$  are unchanged

$$\frac{d\mathbf{p}'_a}{dt'} = \mathbf{F}'_a \quad \mathbf{F}' \equiv \nabla_{\mathbf{r}'} U(|\mathbf{r}'_a - \mathbf{r}'_b|) \quad (1.30)$$