

Alternate Proof of Adiabatic Invariance

• Follows Landau

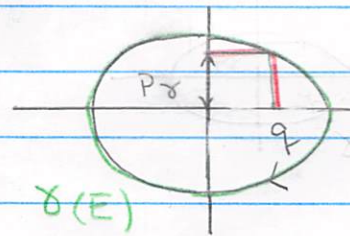
• Consider a 1D system with 1D motion that is periodic and subject to external parameter $\lambda(t)$ which is slow:

$$\text{period } \frac{1}{T} \longrightarrow T \frac{d\lambda}{dt} \ll \lambda \longleftarrow \text{parameter, e.g. } \ell(t) \text{ in example}$$

• The Hamiltonian is $H(q, p, \lambda(t))$. If λ is constant, the energy of the system is $E = H(q, p, \lambda)$. Then,

The system makes a periodic orbit in phase space on curve γ , for constant energy. We can, for given $E = H(q, p, \lambda)$, find p implicitly:

$$p = p_\gamma(q, E, \lambda)$$



• From Hamilton's EOM, $dH/dt = \partial H/\partial t$:

$$\frac{dE(t)}{dt} = \frac{\partial H}{\partial t} = \left(\frac{\partial H}{\partial \lambda} \right)_{p, q} \frac{d\lambda}{dt}$$

← q is always fixed below

- Now we can average over a time scale Δt which is long compared to the period T , but short compared to the adiabatic time scale, T_{\max}

$$\frac{\Delta E}{\Delta t} = \overline{\left(\frac{\partial H}{\partial \lambda}\right)_{p,q} \frac{\Delta \lambda}{\Delta t}}$$

- Where the average means averaged over a cycle

$$\overline{\left(\frac{\partial H}{\partial \lambda}\right)_{p,q}} = \frac{1}{T} \oint_{\gamma} \left(\frac{\partial H}{\partial \lambda}\right)_{p,q} dt$$

$$T \equiv \oint_{\gamma} dt \equiv \text{period}$$

← This is the time averaged force exerted by the string on the ring in our example

- rewriting $dt = dq / \dot{q} = dq / (\partial H / \partial p)_{\lambda}$

$$\frac{\Delta E}{\Delta t} = \frac{\Delta \lambda}{\Delta t} \oint_{\gamma} \left(\frac{\partial H}{\partial \lambda}\right)_{p,q} \frac{dq}{\left(\frac{\partial H}{\partial p}\right)_{\lambda}}$$

$$\oint_{\gamma} \frac{dq}{\left(\frac{\partial H}{\partial p}\right)_{\lambda}}$$

- Now at fixed λ , the energy is constant independent var:

$$\left(\frac{\partial E}{\partial \lambda}\right)_E = 0 = \left(\frac{\partial H}{\partial \lambda}\right)_p \cdot 1 + \left(\frac{\partial H}{\partial p}\right)_{\lambda} \left(\frac{dp}{d\lambda}\right)_E \quad p = p(q, E, \lambda)$$

$$\text{or } \left(\frac{\partial H}{\partial \lambda}\right)_p / \left(\frac{\partial H}{\partial p}\right)_{\lambda} = -\left(\frac{\partial p}{\partial \lambda}\right)_E$$

So

$$\frac{\Delta E}{\Delta t} = \frac{\Delta \lambda}{\Delta t} \frac{\oint - (\partial P / \partial \lambda)_E dq}{\oint (\partial P / \partial E)_q dq}$$

Re-arranging

$$\oint \left(\frac{\partial P}{\partial E} \right) \frac{\Delta E}{\Delta t} + \frac{\partial P}{\partial \lambda} \frac{\Delta \lambda}{\Delta t} dq = 0$$

←—————→

or

$$\Delta P / \Delta t$$

And thus

$$\frac{\Delta I}{\Delta t} = 0 \quad I = \oint \frac{p dq}{2\pi}$$

inserting 2π is
conventional in
the classical
theory

During the evolution of the system with a time dependent parameter λ , the Energy will change in time but the adiabatic invariant $I(E, \lambda) \equiv \oint p dq / 2\pi$ will remain fixed in time.