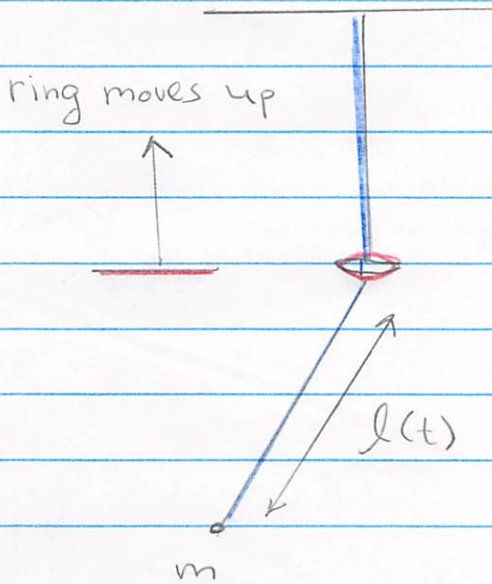


# Adiabatic Invariance: Intro

Example:

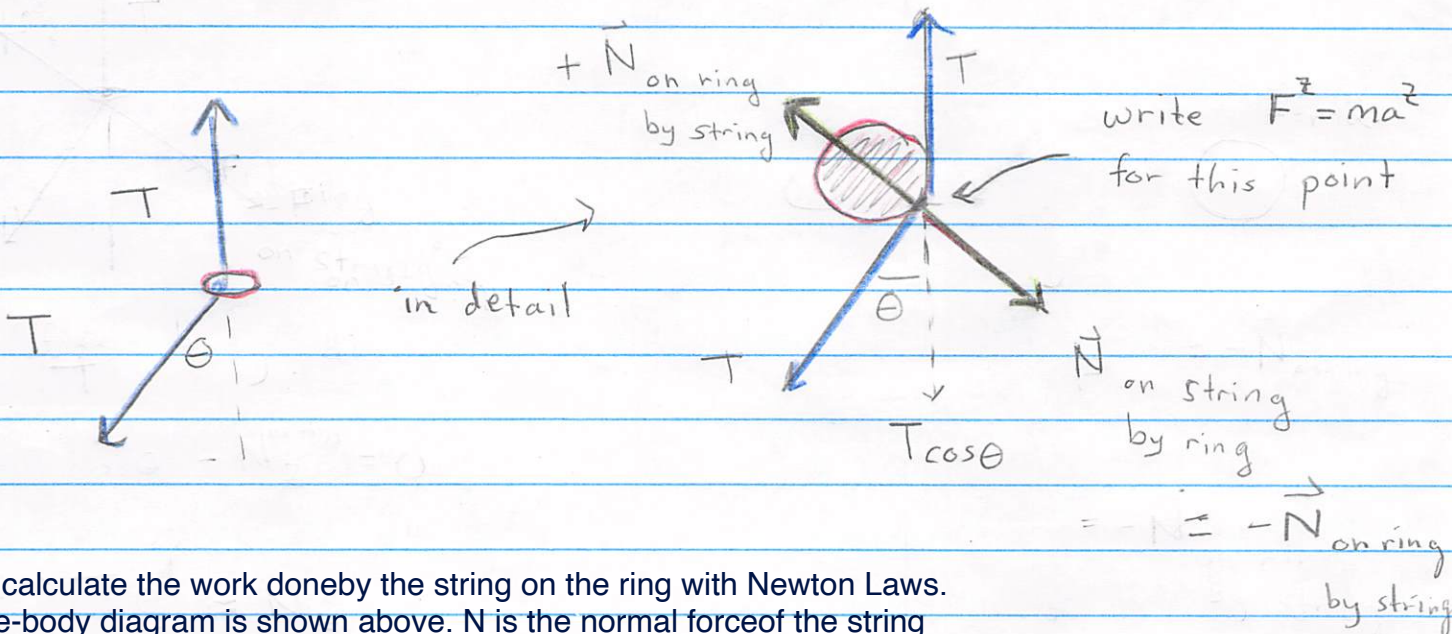


A string pendulum is passed through a ring. The ring is slowly raised.  $l(t)$  gets longer.

The frequency gets slowly lower,  $\omega^2 = g/l(t)$ . The string exerts an upward force on the ring, and as the ring is raised, does work on the ring. The energy available for the oscillations  $E(t)$  gets lower. We will see that

$$\frac{E(t)}{\omega(t)} \approx \text{constant}$$

Simple argument:



We will calculate the work done by the string on the ring with Newton Laws. The free-body diagram is shown above.  $N$  is the normal force of the string on the ring (and vice versa).

- Writing the balance of forces (in  $z$ ) at the pivot:

$$T - T \cos \theta - N_z^{\text{on ring}} = 0$$

←  $z$  - component of force on ring

- So using  $\cos \theta \approx 1 - \theta^2/2$  and  $T \approx mg$  we have:

$$N_z^{\text{on ring}} = mg \theta^2/2$$

- ★ Now let's compute the upward force by averaging over a time,  $\Delta t$ , which is long compared to the period,  $T = 2\pi/\omega(t)$ , but short compared to the total time  $T_{\text{max}}$  over which the length and energy have changed appreciably

(1)

$$2\pi/\omega_0 \ll \Delta t \ll T_{\text{max}}$$

Yielding

$$\overline{N_z} \approx mg \overline{\theta^2} = \frac{\overline{E(t)}}{2l}$$

- Thus the work done,  $\Delta W = -\Delta E = \overline{N_z} \Delta l$ , implies:

$$(2) \quad -\frac{\Delta E}{\Delta t} = \frac{\overline{E}}{2l} \frac{\Delta l}{\Delta t}$$

• This equation is easily solved and has

$$E(t) \sqrt{l(t)} = \text{const}$$

i.e

$$\boxed{\frac{E(t)}{\omega(t)} = \text{const}}$$

$$\omega^2(t) \propto \frac{g}{l(t)}$$

As, we will see in the next section this is a consequence of a general theorem that

$$I = \oint \frac{p dq}{2\pi} \quad \leftarrow \text{adiabatic invariant}$$

is approximately constant under a slow change of a parameter

## Adiabatic Invariants

Consider a 1D Hamiltonian which depends on a parameter  $\lambda(t)$  and assume that  $\lambda$  changes slowly in time compared to other times  $2\pi/\omega_0$  in the problem.

$$E(t) \equiv H(q, p, \lambda(t))$$

Over a time  $\Delta t$  which is long compared to  $2\pi/\omega_0$  but short compared to  $\tau_{\max}$  (the time over which  $\lambda$  and  $E(t)$  change by order one), the energy is approximately constant and the system stays on an iso-energetic contour in phase-space. The area enclosed by this contour (divided by  $2\pi$  by convention) is

$$\underline{I}(t) = \underbrace{\int \frac{dp dq}{2\pi}}_{\text{Area}} = \underbrace{\oint p \frac{dq}{2\pi}}_{\text{1 cycle}} \quad \text{See figure on next page!}$$

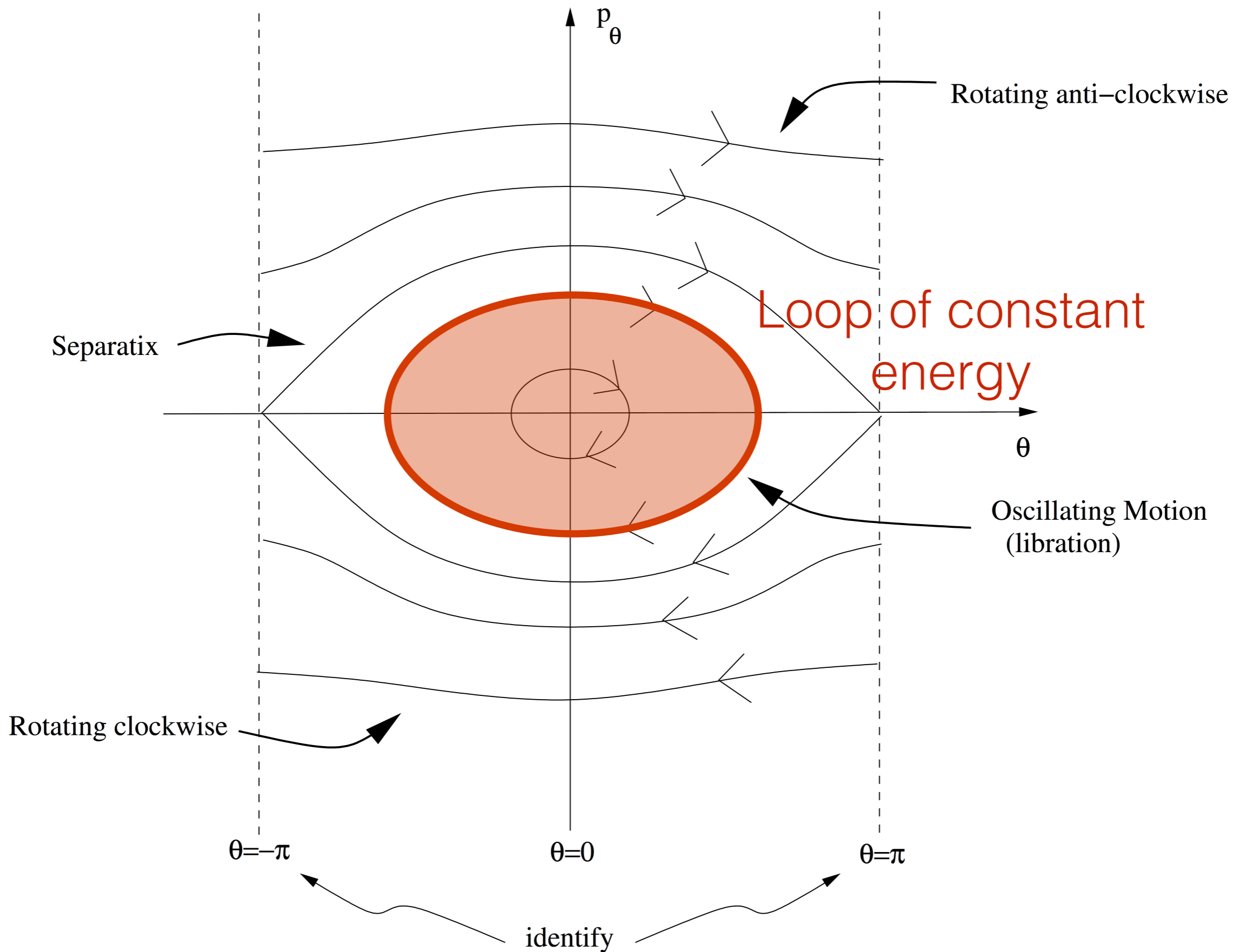
where the integral is done by treating  $E$  as constant for the cycle. See below

We are using the curl theorem in 2D, with the vector field  $\vec{v} = (p, 0)$ . The curl theorem in 2d relates the line integral to the area integral:

$$\oint_{\text{loop}} \vec{v} \cdot d\vec{\ell} = \int_{\text{area}} d\vec{A} \cdot (\nabla \times \vec{v}).$$

We have chosen the loop to circulate in phase-space in a clockwise sense, so that the area points into the page, leading to  $\int_{\text{line}} p dq = \int_{\text{area}} dp dq$ . In terms of forms (if you know it),  $d(pdq) = dp \wedge dq$ .

$$\text{area} = \int dp dq = \oint p dq$$



- Now we have an exact theorem which says that the area enclosed under the full time evolution will not change in time (Liouville theorem). So we must have over

$$I(t) = \oint_{\text{cycle}} \frac{p dq}{2\pi} \approx \text{const} \quad \text{time } T_{\text{max}}$$

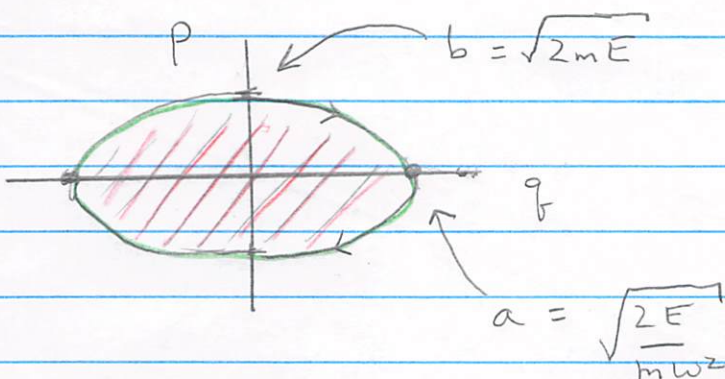
### Example SHO

- The Hamiltonian is

$$H = \frac{p^2}{2m} + \frac{1}{2} m \omega^2(t) q^2$$

with  $\omega$ -changing slowly over time  $T_{\text{max}}$ . For short times  $\Delta t \ll T_{\text{max}}$  the energy is constant, as is  $\omega$

$$E = \frac{p^2}{2m} + \frac{1}{2} m \omega^2 q^2 \quad \leftarrow \text{equation of ellipse: } (x/a)^2 + (y/b)^2 = \text{const}$$



Area  
of ellipse =  $\pi ab$

So

$$I = \oint_{\text{cycle}} \frac{p dq}{2\pi} = \int \frac{dp dq}{2\pi} = \frac{\pi ab}{2\pi}$$

area

$$I = \frac{E}{\omega}$$

Now if  $E$  and  $\omega$  change in time the Liouville theorem guarantees that  $I$  is constant i.e.,

$$I = \frac{E(t)}{\omega(t)} \approx \text{constant}$$