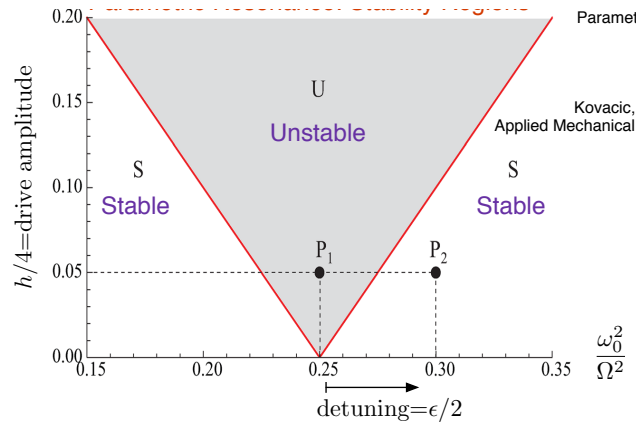


## Problem 1. Parametric resonance with damping

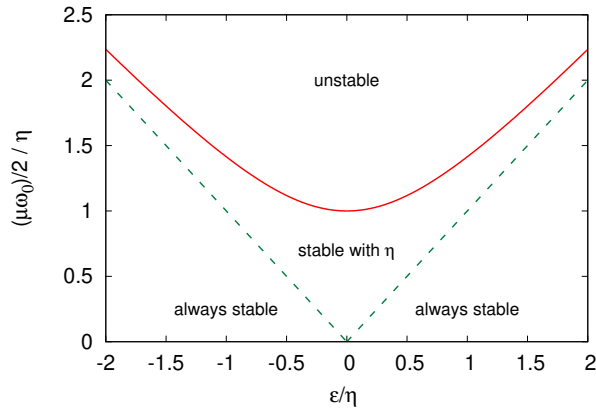
Consider an oscillator with a small damping coefficient  $\eta$ , and a time dependent mass,  $m(t) = m_0(1 + \mu \cos(\Omega t))$ , with  $\mu$  is small. The frequency is  $\Omega \simeq 2\omega_0 + \epsilon$  with  $\epsilon$  also small. Thus the equation of motion is

$$\frac{d}{dt}(m(t)\dot{q}) + m_0\omega_0^2 q(t) + m_0\eta\dot{q} = 0 \quad (1)$$

Determine the regions in the  $\epsilon, \mu$  plane where the oscillations are stable and unstable. How is the plot from class (the first plot below) modified by the non-zero damping coefficient?



You should find the following picture:



## Problem 2. A pendulum in a harmonic electric field

A simple pendulum consists of a particle of mass  $m$  at the end of weightless rod of length  $\ell$ . The particle has a charge  $q$  and sits in an electric field of amplitude  $E_0$ , directed in the horizontal direction, which oscillates rapidly with frequency  $\Omega$ ,  $\Omega \gg \sqrt{g/\ell}$ .

- (a) Determine the Lagrangian for this system, and the equation of motion.
- (b) Above a critical field strength  $E_c$  the position  $\phi = 0$  (the bottom of the pendulum) becomes unstable. Determine  $E_c$  and determine the new point of stability for  $E > E_c$ . Sketch the effective potential for  $E < E_c$  and  $E > E_c$ .
- (c) Analyze the validity of your approximations. Is the critical field large or small compared to  $mg/q$  when your approximation is valid?

### Problem 3. (Laurence Yaffe) A driven set of oscillators

**General Background:** Consider a set of coupled harmonic oscillators interacting with external time dependent forces. The oscillator Lagrangian without the forces reads<sup>1</sup>

$$L_0 = \sum_{ij} \frac{1}{2} M_{ij} \dot{q}^i \dot{q}^j - \frac{1}{2} K_{ij} q^i q^j . \quad (2)$$

The Lagrangian for the forces driving the system is

$$L_{\text{int}} = \sum_i F_i(t) q^i , \quad (3)$$

and the total Lagrangian is  $L = L_0 + L_{\text{int}}$ . As always, switch coordinates to the eigen basis of the generalized eigenvalue problem

$$q^i = \sum_a E_a^i Q^a , \quad (4)$$

where the  $\vec{E}_a$  is the  $a$ -th eigen-vector of the generalized eigenvalue problem,  $K\vec{E}_a = \lambda_a M\vec{E}_a$ . Recall that the natural frequency associated with the  $a$ -th normal mode is  $\lambda_a = \omega_a^2$ , and the eigenvectors are orthonormal with the mass matrix as weight:

$$\sum_{ij} E_a^i M_{ij} E_b^j = \delta_{ab} . \quad (5)$$

- (a) Determine the Lagrangian for the coordinates  $Q^a$ , and show that the resulting equation of motion is

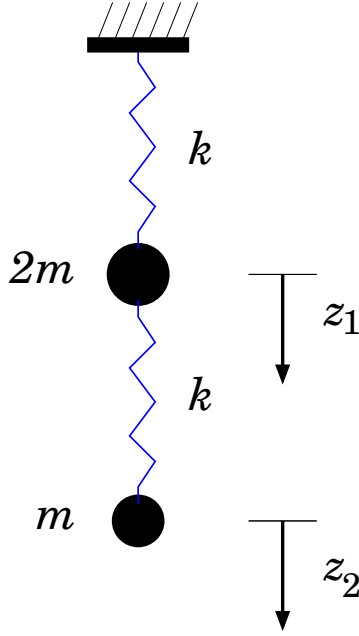
$$\ddot{Q}^a + \omega_a^2 Q^a = F_a , \quad (6)$$

where  $F_a = \sum_i F_i E_a^i$  is the projection of the force vector  $\vec{F} = (F_i)$  onto the  $a$ -th normal mode, i.e.  $F_a = \vec{F} \cdot \vec{E}_a$  .

**Problem:** Now consider two masses,  $m_1 = 2m$  and  $m_2 = m$ , are suspended in a uniform gravitational field  $g$  by identical massless springs with spring constant  $k$ . Assume that only vertical motion occurs, and let  $z_1$  and  $z_2$  denote the vertical displacement of the masses from their equilibrium positions, increasing in the downward direction as shown. An external time-dependent force  $F(t)$  is applied to the lower mass (with  $F > 0$  indicating a downward vertical force). Assume that the external force vanishes as  $t \rightarrow \pm\infty$ , with the system initially at rest in its equilibrium configuration at time  $-\infty$ . Let  $F(\omega)$  denote the Fourier transform of  $F(t)$ .

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<sup>1</sup>For the rest of this problem we will not use the summation convention.



- (b) Construct the Lagrangian for the system without the force and find the normal modes and frequencies. Then include the external force, and find the resulting equations of motion.
- (c) Solve for the motion of both masses (expressed as an integral involving the time-dependent force).
- (d) Find the total work done on the system by the external force,  $W = E(+\infty) - E(-\infty)$ . Show that it can be expressed in the form

$$W = \int_{-\infty}^{\infty} \frac{d\omega}{2\pi} \chi(\omega) |F(\omega)|^2 \quad (7)$$

with  $\chi(\omega)$  real and positive.  $\chi(\omega)$  is known as the response function, and will be proportional to a delta-functions in the absence of damping.

- (e) If a small damping term is added to each equation of motion, so  $m_i \ddot{z}_i \rightarrow m_i \ddot{z}_i + m_i \eta \dot{z}_i$ , make an educated guess how this qualitatively changes the response function  $\chi(\omega)$  and make a sketch of  $\chi(\omega)$ .
- (f) (Optional) With the dissipation described in the previous item, again find the total work done by the force on the system. ( $W$  is not equal to  $E(\infty) - E(-\infty)$ , since the work done is ultimately dissipated away.) Determine  $\chi(\omega)$  in this case.