

Secular Perturbation Theory

- Trying to solve small

$$\ddot{x} + \omega_0^2 x = -\beta x^3 \equiv f_{\text{ind}}/m$$

Now we try

$$x(t) = x^{(0)} + x^{(1)} + x^{(2)}$$

But we try a more general form for the zeroth order solution:

$$\begin{aligned} x^{(0)} &= \text{Re} [A(t) e^{-i\omega_0 t}] \\ &= a(t) \cos(-\omega_0 t + \varphi(t)) \end{aligned}$$

- Now the amplitude and phase are slow functions of time, adjusted to remove secular divergences

$$da/dt \ll a\omega_0 \quad d\varphi/dt \ll \omega_0$$

- Note define $\psi \equiv -\omega_0 t + \varphi(t)$; extra

$$\begin{aligned} \frac{d^2 x^{(0)}}{dt^2} &= -\omega_0^2 x^{(0)} + \overbrace{2\dot{\omega}_0 \dot{\varphi} a \cos \psi + 2\omega_0 \dot{a} \sin \psi}^{\text{extra}} \\ &+ \text{neglectable second derivs of } a \text{ and } \varphi \end{aligned}$$

• As before

$$\frac{f_{\text{ind}}}{m} \approx -\beta (X^{(0)})^3 = -\frac{\beta a^3}{4} \cos(3\varphi) - \frac{3}{4} \beta a^3 \cos(\varphi)$$

• Substituting $X = X^{(0)} + X^{(1)}$ into $\ddot{x} + \omega_0^2 x = f_{\text{ind}}/m$ we find:

$$\frac{d^2 X^{(1)}}{dt^2} + \omega_0^2 X^{(1)} = \underbrace{-\frac{\beta a^3}{4} \cos(3\varphi)}_{f^{(1)}} + \left(-\frac{3}{4} \beta a^3 - 2\omega_0 \dot{\varphi} \right) a \cos \varphi + \underbrace{(-2\omega_0 \dot{a})}_{\text{secular}} \sin \varphi$$

• The $\cos \varphi$ and $\sin \varphi$ terms are secular.
We adjust $\dot{\varphi}$ and \dot{a} so these do not appear

So we find

$$\begin{aligned} -2\omega_0 \dot{a} = 0 &\longrightarrow a = \text{constant} \\ -\frac{3}{4} \beta a^3 - 2\omega_0 \dot{\varphi} = 0 &\longrightarrow \varphi = -\Delta\omega t + \varphi_0 \end{aligned}$$

with $\Delta\omega = \frac{3}{8} \left(\frac{\beta a^2}{\omega_0^2} \right) \omega_0$

• The remaining $f^{(1)}$ term proceeds as before

$$X^{(1)} = -\frac{1}{32} \left(\frac{\beta a^2}{\omega_0^2} \right) a \cos(3\varphi) \quad \varphi = -(\omega_0 + \Delta\omega)t + \varphi_0$$

• So finally

$$x(t) = a \cos(\omega t + \varphi_0) + x^{(1)}$$

Where

$$\omega = \omega_0 + \Delta\omega = \omega_0 \left(1 + \frac{3}{8} \left(\frac{\beta a^2}{\omega_0^2} \right) \omega_0 \right)$$

Comment:

• For a steady state secular perturbation theory reduces to taking $x^{(0)} = a \cos(-\omega t + \varphi)$

$$a = \text{const}$$

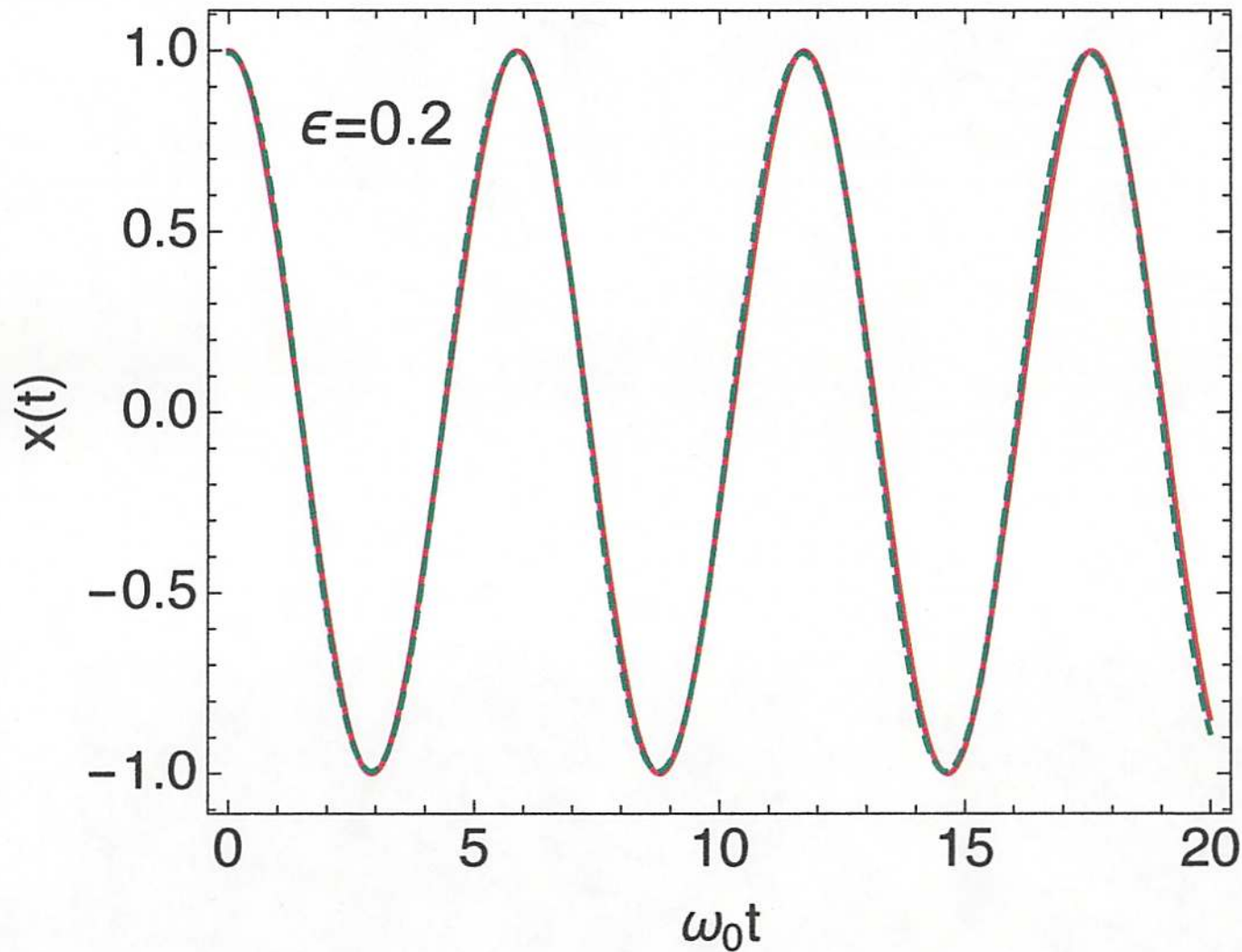
← shift

$$\varphi = -\Delta\omega t + \varphi_0$$

Where a and $\Delta\omega$ are adjusted at each order so that secular terms do not appear.

Non linear oscillator treating secular term as frequency shift

$$\epsilon = \frac{\beta a^2}{\omega_0^2}$$



Solid lines exact solution, dashed lines approximate