

## E Separated coordinates for magnetostatics

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### E.1 Overview

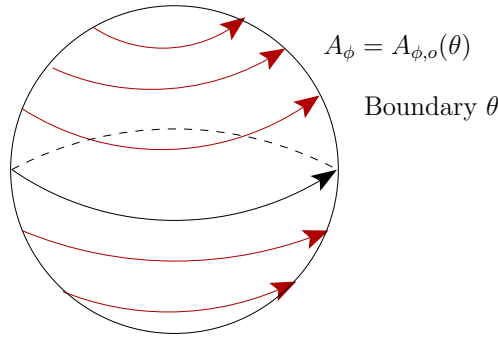
- (a) The magnetostatic equations are complicated, and we refer to [Wikipedia](#) for the form of the vector Laplacian in various coordinate systems.
- (b) For currents running strictly up and down  $A_z(x, y)$  the magnetostatic equations reduce to

$$-\left(\frac{\partial^2}{\partial x^2} + \frac{\partial^2}{\partial y^2}\right)A_z(x, y) = 0 \quad (\text{E.1})$$

which has the same form as 2D electrostatics. The appropriate separated solutions are given in Appendix [D.4](#).

- (c) For currents which are azimuthally symmetric  $\mathbf{j} = j_\phi(r, \theta) \hat{\phi}$  we may either use spherical or cylindrical coordinates. The spherical case is discussed in Appendix [E.2](#).

### E.2 Spherical coordinates for magnetostatics



- (a) The vector Laplacian for azimuthally symmetric currents, and the ansatz  $\mathbf{A} = A_\phi(r, \theta) \hat{\phi}$  reads

$$\left[ -\left( \frac{1}{r^2} \frac{\partial}{\partial r} r^2 \frac{\partial}{\partial r} + \frac{1}{r^2 \sin \theta} \frac{\partial}{\partial \theta} \sin \theta \frac{\partial}{\partial \theta} \right) + \frac{1}{r^2 \sin^2 \theta} \right] A_\phi(r, \theta) = 0 \quad (\text{E.2})$$

This is an appropriate equation only if the current takes a specific symmetric form

$$\mathbf{j} = j_\phi(r, \theta) \hat{\phi}$$

- (b) The eigenfunctions are along the boundary direction  $\theta$ , and are regular at  $\theta = 0$  and  $\pi$ . They are associated Legendre Polynomials with  $m = 1$

$$\psi_\ell(\theta) = P_\ell^1(\cos \theta) \quad \ell = 1 \dots \infty$$

The first few eigenfunctions are given [here](#). Perhaps the most important fact is that they all are proportional to  $\sin(\theta)$  guaranteeing regularity of  $\mathbf{A} = A_\phi \hat{\phi}$  at  $\theta = 0$  and  $\pi$ .

(c) Orthogonality:

$$\int_{-1}^1 d(\cos \theta) P_\ell^1(\cos \theta) P_{\ell'}^1(\cos \theta) = \frac{2}{2\ell + 1} \frac{(\ell + 1)!}{(\ell - 1)!} \delta_{\ell\ell'}$$

(d) Completeness

$$\sum_{\ell=1}^{\infty} \frac{2\ell + 1}{2} \frac{(\ell - 1)!}{(\ell + 1)!} P_\ell^1(x) P_\ell^1(x') = \delta(x - x') \quad (\text{E.3})$$

(e) Solution

$$A_\phi(r, \theta) = \sum_{\ell=1}^{\infty} \left[ A_\ell r^\ell + \frac{B_\ell}{r^{\ell+1}} \right] P_\ell^1(\cos \theta) \quad (\text{E.4})$$